Investigating General-Relativistic Transonic Accretion in Kerr Black Holes

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#### Introduction

- Radial fluid velocity, *n*
- Sound speed,  $a_s$

Mach number

$$M(r) = \frac{u(r)}{a_s(r)}$$

- Accretion flow
  - Subsonic: M < 1
  - Supersonic: M > 1
  - Transonic: crosses M = 1

- Transition
  - Smooth -- Sonic point
  - Discontinuous Shock
- Multi-transonic flow = 3 sonic points

## Motivation

 Black Hole (BH) inner boundary condition : Supersonic flow at Event Horizon (EH)
 Far away from EH – subsonic
 Hence BH accretion is essentially transonic
 Except cases where already supersonic initially
 Shock Formation → Multi-transonic flow

Accretion in Kerr (rotating) Black Holes:

- Transonic
- Consider General-Relativistic effects

#### <u>Notations</u>

Gravitational Radius

 $r_g = \frac{GM_{BH}}{c^2}$ 

Black Hole Spin, Kerr Parameter = a
|a| ∈ (0, 0.9982)
λ = Specific angular momentum of flow, is aligned with a
BH

Co-rotating : (+ve) a
Counter-rotating : (-ve) a

#### Goal

 Study behavior of dynamic and thermodynamic properties of accretion flow very close (≈ 0.01 r<sub>g</sub>) to EH & their dependence on spin of BH

#### <u>What do we do?</u>

 Formulate & solve conservation equations governing general relativistic, multi-transonic, advective accretion flow in Kerr metric

#### Metric

 Kerr metric in equatorial plane of BH (Novikov, I.D. & Thorne, K.S. 1973, in Black Holes, 343)

$$ds^{2} = g_{\mu\nu}dx^{\mu}dx^{\nu} = -\frac{r^{2}\Delta}{A}dt^{2} + \frac{A}{r^{2}}(d\phi - \omega dt)^{2} + \frac{r^{2}}{\Delta}dr^{2} + dz^{2}$$

$$\Delta = r^2 - 2r + a^2$$

$$A = r^4 + r^2 a^2 + 2ra^2 \qquad \qquad \omega = \frac{2ar}{A}$$

$$g_{tt} = \frac{A\omega^2}{r^2} - \frac{r^2\Delta}{A}$$

$$g_{t\phi} = -\frac{A\omega}{r^2}$$

$$g_{\phi\phi} = \frac{A}{r^2}$$



#### **Accretion Disk**



Vertically integrated model (Matsumoto, R. et al. 1984, PASJ, 36, 71)
Flow in hydrostatic equilibrium in transverse direction
Thermodynamic quantities vertically averaged over disk height
Calculate quantities on equatorial plane of BH

Disk height (Abramowicz, M.A. et al. 1997, ApJ, 479, 179)

$$h_{disk}(r) = H(r) = r^{2} \left[ \frac{2na_{s}^{2}}{(n+1)(1-na_{s}^{2})(\lambda^{2}v_{t}^{2}-a^{2}(v_{t}-1))} \right]^{\frac{1}{2}}$$

#### **Conserved Quantities**

Specific energy (including rest mass)

#### Mass accretion rate

$$\dot{M}(r) = 4\pi\Delta^{1/2}H(r)\rho\frac{u}{\sqrt{1-u^2}}$$

#### Entropy accretion rate

$$\dot{\Xi} = K^{1/(\gamma-1)} \dot{M}_{in} = 4\pi \left(\frac{n}{n+1}\right)^n \Delta^{1/2} \left(\frac{a_s^2}{1-na_s^2}\right)^n H(r) \frac{u}{\sqrt{1-u^2}}$$

 $\varepsilon = hv_t = \left| \frac{\gamma - 1}{\gamma - (1 + \theta^2 T)} \right| v_t$ 

## Methodology

Solve conservation eqns. to get du/dr

$$\frac{d\varepsilon}{dr} = 0, \frac{d\dot{\Xi}}{dr} = 0.$$

$$\frac{du}{dr} = \frac{N}{D} = fn(r, u, a_s, \varepsilon, \lambda, \gamma, a)$$

Start from sonic point and integrate to get flow properties at any *r*For a range of [ε-λ-γ-a] → get 3 sonic points on solving eqns.

 $\blacksquare r_{out} > r_{mid} > r_{in}$ 

Result plots for black hole of mass,  $M_{BH}=10M_{sun}$ 

#### Solution topology for multi-transonic accretion in Kerr geometry



#### [ $\epsilon$ - $\lambda$ ] parameter space ( $\gamma$ =4/3, *a*=0.3) for monoand multi-transonic accretion & wind



#### $\epsilon$ - $\lambda$ multi-transonic space for different $\gamma$



Differentiating **Accretion Properties of** Co- & Counter-**Rotation of Central Black Hole** 



Angular Momentum dependence of Multitransonicity for Pro & Retrograde Flow



Angular Momentum dependence of Multitransonicity for Pro & Retrograde Flow

## Pro-vs. Retro-Grade Flows

<u>Prograde</u> Accretion Flow (co-rotating)

<u>Retrograde</u> Accretion Flow (counter-rotating)

 Multi-transonicity much more common

> Covers higher value of angular momentum (λ)

 Multi-transonic regions at lower value of λ

# Terminal Behavior of Accretion Variables

Terminal value of accretion variable A : A<sub>δ</sub> = A (at r<sub>δ</sub> = r<sub>e</sub> + δ)
Integrate flow from r<sub>c</sub> down to r<sub>δ</sub> → get A<sub>δ</sub>
Studied variation of A<sub>δ</sub> with a → BH spin dependence of accretion variables very close to event horizon

■ For  $M_{BH}$ =10 $M_{sun}$ , and  $\delta$ =0.01 → plots next ...





## **Astrophysical Implications**

Preliminary step towards understanding how BH spin affects astrophysical accretion

Shock waves in BH accretion disks must form through multi-transonic flows

Study of the post shock flow helpful in explaining:
 Spectral properties of BH candidates
 Cosmic (galactic & extragalactic) jets powered by accretion (formation & dynamics)
 Origin of Quasi Periodic Oscillations in galactic sources

#### Summary

- 1. Study dependence of multi-transonic accretion properties on BH spin at any radial distance
  - Very close to event horizon
- 2. Found non-trivial difference in,
  - □ Prograde flow (co-rotating BH) : low- $\lambda$
  - $\square \quad Retrograde flow (counter-rotating BH) : multi-transonicity at high \lambda & greater possibility of shock formation$
- 3. Future work:
  - Shock location
  - Analog gravity
     Analog Hawking radiation



#### References

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- Barai, P., Das, T.K. & Wiita, P.J. 2004, ApJ, 613, L49
- Das, T.K., Bilic, N. & Dasgupta, S. 2006, astro-ph/0604477
- Matsumoto, R. et al. 1984, PASJ, 36, 71
- Novikov, I.D. & Thorne, K.S. 1973, in Black Holes, 343

# Thank You All

#### Contents

Introduction & Motivation
What are we doing?
Our system formulations
Results
Conclusions



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## Motivation

- Transonic Flow is relevant in various astrophysical situations:
  - Black Hole (BH) or Neutron Star accretion
  - Collapse / Explosion of stars
  - Solar / Stellar winds
  - Formation of protostars & galaxies
  - Interaction of supersonic galactic (or extragalactic) jets with ambient medium

- BH inner boundary condition : Supersonic flow at Event Horizon (EH)
- Far away from EH subsonic
- Hence BH accretion is essentially transonic
  - Except cases where already supersonic initially
- Multi-transonic flow → Shock Formation

# **Describing Our System**

- No self-gravity, no magnetic field
- Units:  $G = c = M_{BH} = 1$
- Boyer-Lindquist coordinates (-+++)
- Observer frame corotating with accreting fluid
- λ = Specific angular momentum of flow -- aligned with a
  Stationary & axisymmetric flow
- Euler & Continuity eqns :

$$\nabla^{\mu}\mathfrak{I}^{\mu\nu} = 0 \quad , \quad \left(\rho v^{\mu}\right)_{;\mu} = 0$$

Polytropic equation of state
 *K* ~ specific entropy density
 *γ* = adiabatic index, *n* = polytropic index

$$p = K \rho^{\gamma}$$

$$n = \frac{1}{\gamma - 1}$$

# Our System ... Specific proper flow enthalpy, *h*Polytropic sound speed, *a<sub>s</sub>*

$$a_{s} = \left(\frac{\partial p}{\partial \varepsilon}\right)_{s}^{1/2} = \Psi_{1}(p, r, \gamma) = \Psi_{2}(p, \rho, \gamma) ; a_{s}^{2}(r) = \frac{\gamma K_{B}T(r)}{\mu m_{H}}$$

Frame dragging neglected
Weak viscosity limit

- Very large radial velocity close to BH → Timescale (viscous >> infall)
- Effect of Viscosity  $\rightarrow \lambda \downarrow \rightarrow$  Flow behavior as function of  $\lambda$  provides information on viscous transonic flow

#### Metric & Others

Kerr metric in equatorial plane of BH (Novikov & Thorne 1973)

$$ds^{2} = g_{\mu\nu}dx^{\mu}dx^{\nu} = -\frac{r^{2}\Delta}{A}dt^{2} + \frac{A}{r^{2}}(d\phi - \omega dt)^{2} + \frac{r^{2}}{\Delta}dr^{2} + dz^{2}$$

$$\Delta = r^2 - 2r + a^2 \qquad A = r^4 + r^2 a^2 + 2ra^2 \qquad \omega = \frac{2ar}{A}$$
$$g_{tt} = \frac{A\omega^2}{r^2} - \frac{r^2 \Delta}{A} \qquad g_{t\phi} = -\frac{A\omega}{r^2} \qquad g_{\phi\phi} = \frac{A}{r^2}$$

$$g_{t\phi} = -\frac{A\omega}{r^2} \qquad \qquad g_{\phi\phi} = \frac{A}{r^2}$$

 $\square$  Angular velocity,  $\Omega$ 

$$\Omega = \frac{u^{\phi}}{u^{t}} = -\frac{\left(g_{t\phi} + \lambda g_{tt}\right)}{\left(g_{\phi\phi} + \lambda g_{t\phi}\right)}$$

■ 4<sup>th</sup> Component of velocity:

$$v_{t} = \left[\frac{\Delta}{(1-u^{2})(1-\Omega\lambda)(g_{\phi\phi}+\lambda g_{t\phi})}\right]^{\frac{1}{2}}$$

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## Methodology

Solve conservation eqns. to get *du/dr* 

$$\frac{d\varepsilon}{dr} = 0, \frac{d\dot{\Xi}}{dr} = 0.$$

$$\frac{du}{dr} = \frac{N}{D} = fn(r, u, a_s, \varepsilon, \lambda, \gamma, a)$$

Start from sonic point and integrate to get flow properties at any r

By setting N=0, D=0 get sonic point quantities:

- $\blacksquare \quad \mathcal{U}_{o}, \ a_{s|c}$
- Quadratic eqn. for  $(du/dr)_c$
- For some  $[\varepsilon \lambda \gamma a] \rightarrow \text{get 3 sonic points on solving eqns.}$

Observational Consequences
 Debate: Determining BH spin from observations

 Most popular approach: study of skew shaped fluorescent iron lines

Our approach presents the potential to deal with this problem

We predict behavior of flow properties
 For hot, low-λ prograde accretion flow & high-λ retrograde flow

#### Results

- For some  $[\mathbf{\mathcal{E}}-\lambda-\gamma-a]$  -- get 3 sonic points on solving eqn
  - $\mathbf{I} \mathbf{r}_{\mathrm{out}} > \mathbf{r}_{\mathrm{mid}} > \mathbf{r}_{\mathrm{in}}$
- $\blacksquare$   $r_{out}$ ,  $r_{in}$ : X-type sonic points
- r<sub>mid</sub>: O-type sonic pt (unphysical no steady transonic soln passes thru it)
- Multitransonic Accretion:  $\Xi(r_{in}) > \Xi(r_{out})$
- Multitransonic Wind  $: \Xi(r_{in}) < \Xi(r_{out})$
- General astrophysical accretion  $\rightarrow$  Flow thru  $r_{out}$
- Flow thru  $r_{in}$  possible only in case of a shock
  - If supersonic flow thru r<sub>out</sub> is perturbed to produce entropy = [Ξ(r<sub>in</sub>) Ξ(r<sub>out</sub>)], it joins subsonic flow thru r<sub>in</sub> forming a standing shock
     <u>Shock details from GR Rankine-Hugoniot</u> conditions





#### E- $\lambda$ Multi-Transonic Space for different *a*



## Angular Momentum

Weakly rotating flows, found in several physical situations:

- Detached binary systems fed by accretion from OB stellar winds
- Semidetached low-mass nonmagnetic binaries
- Supermassive BHs fed by accretion from slowly rotating central stellar clusters
- Turbulence in standard Keplerian accretion disk

We found:

 At higher values of angular momentum multi-transonicity is more common for retrograde flow





# Analog Hawking Radiation (Das, Bilic & Dasgupta, 2006)

Hawking Temperature

$$T_{H} = \frac{\hbar c^{3}}{8\pi K_{B}GM_{BH}}$$

Analog Hawking Temperature

$$T_{AH} = \frac{\hbar}{4\pi K_B} \left[ \frac{1}{c_s} \frac{du_{\perp}^2}{d\eta} \right]_{\text{Acoustic Horizon}}$$

$$\tau = \frac{T_{AH}}{T_{H}}$$

#### Variation of the ratio of analog to actual Hawking temperature with black hole spin

