

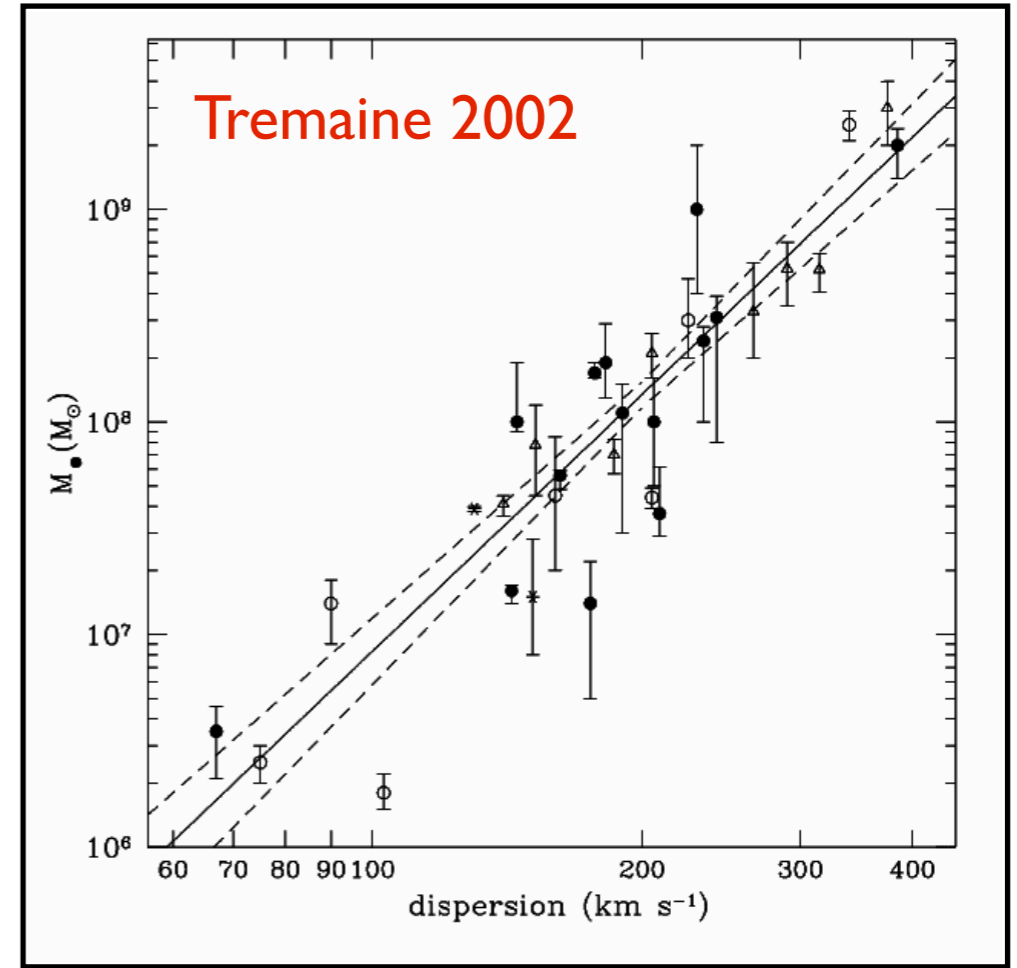
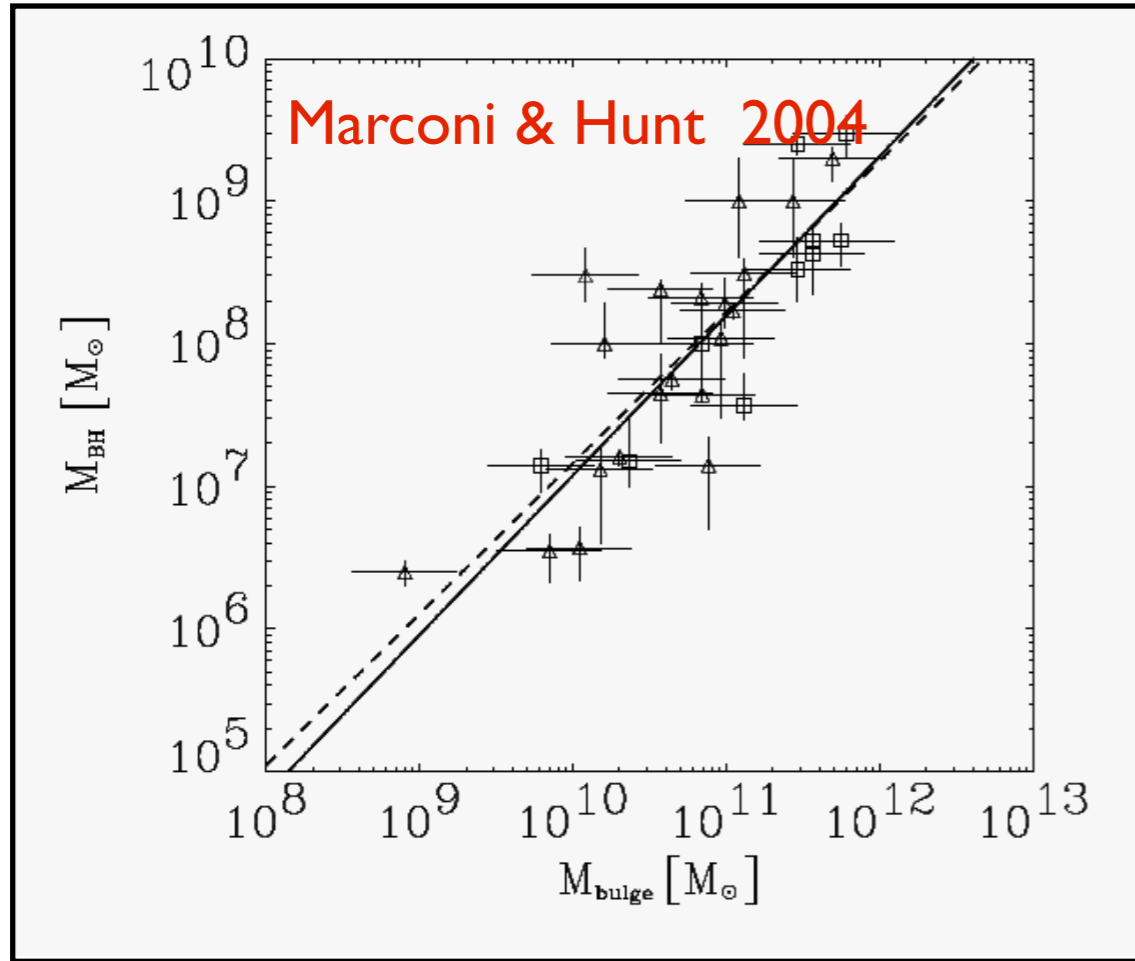
# The Formation of Massive Black Holes: a Review

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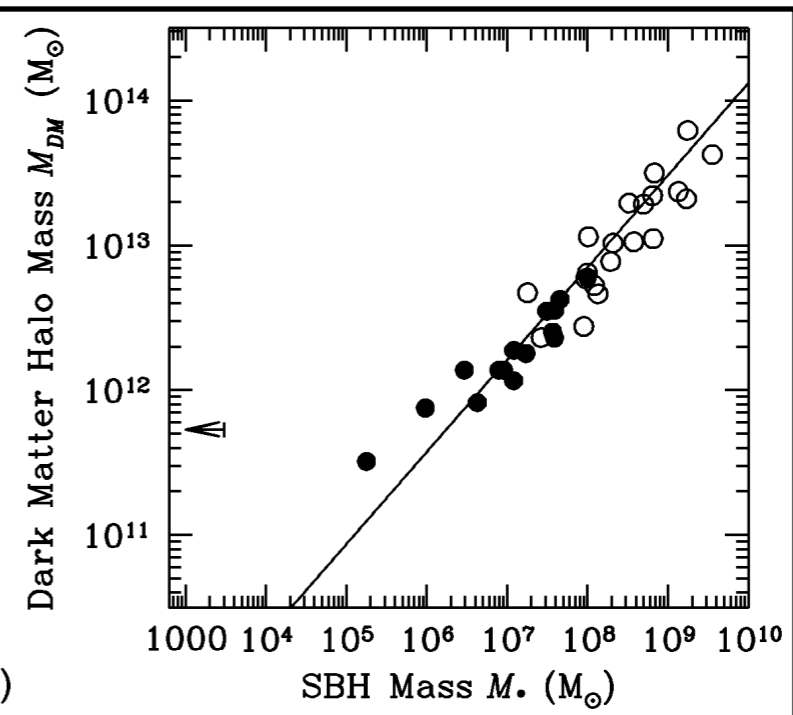
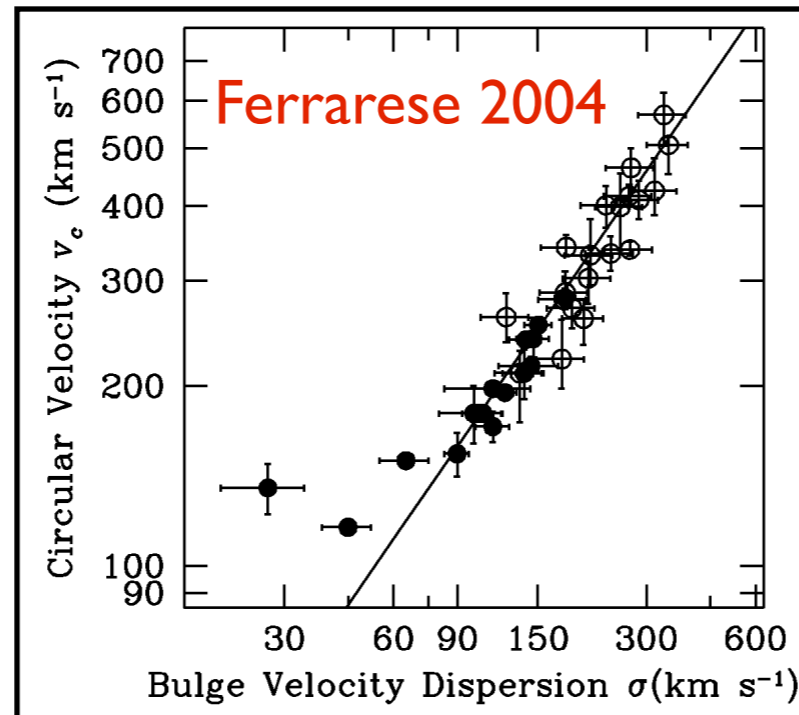
# Understanding the formation and growth of MBHs along the Cosmic history.

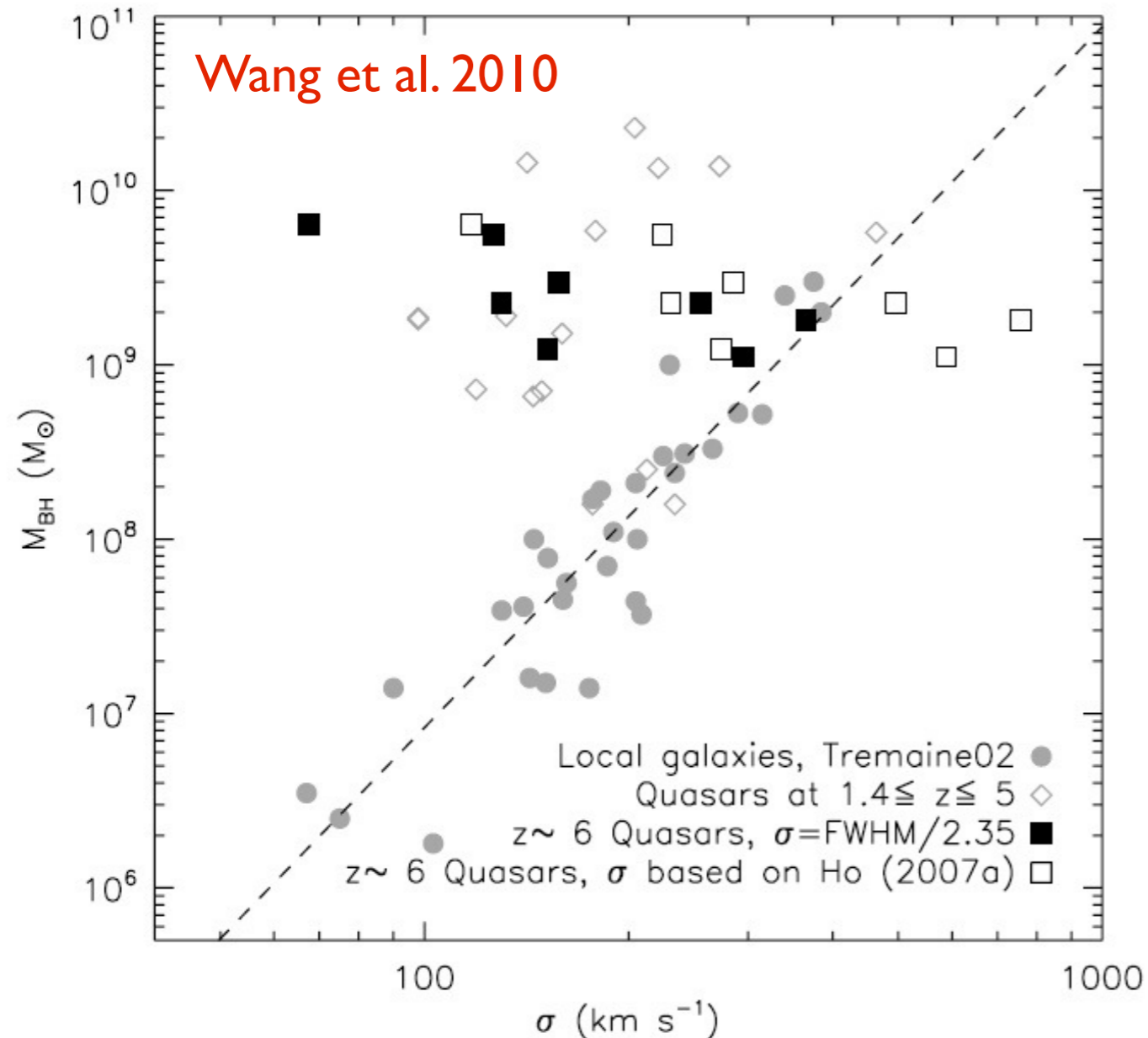
1. How fast can MBHs grow to  $10^9$ - $10^{10} M_{\odot}$ ?
2. How many of them along the Cosmic history?

We need to understand the physics of black hole *and* galaxy mergers.



**LOCAL SCALING RELATIONS**

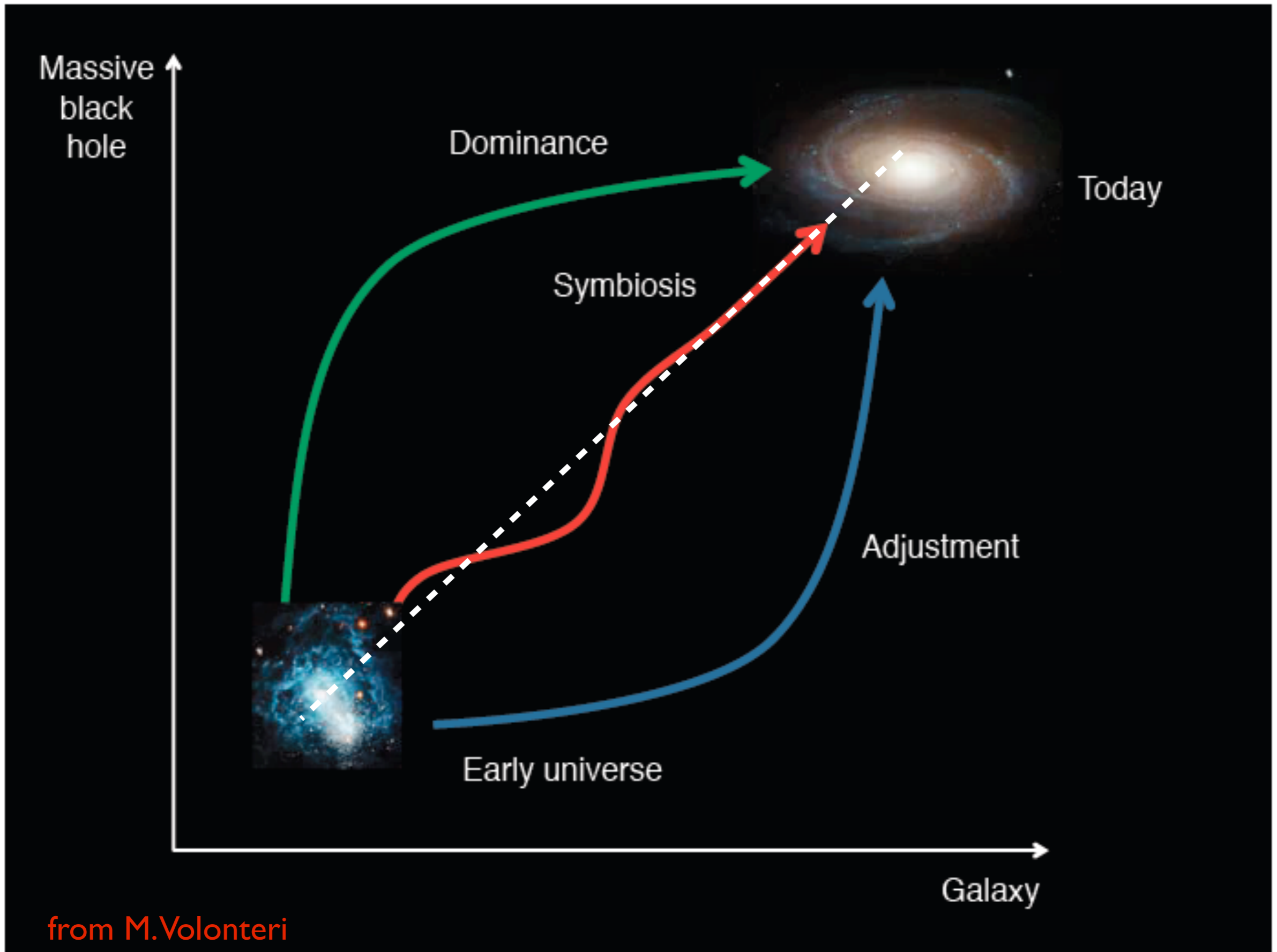




1. There seems to be little or no correlation between mass and velocity dispersion (Wang et al. 2010).

2. Typically black holes are ‘over massive’ at fixed mass/velocity dispersion compared to  $z=0$  counterparts (e.g. Walter et al 2004; Decarli et al. 2010; Merloni 2010).

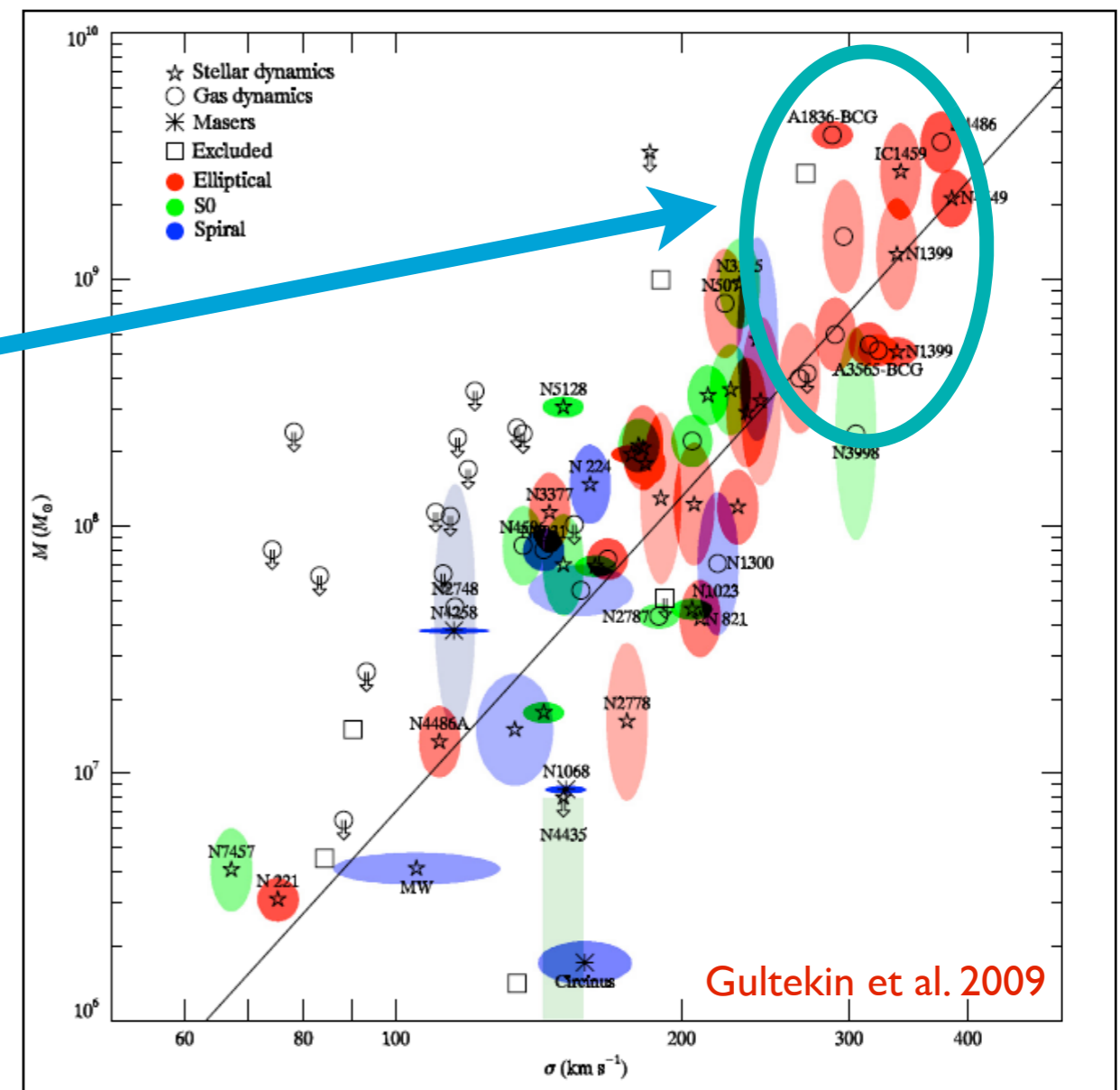
3. Studies suggest that either many massive galaxies do not have MBHs or these BH are less massive than expected (Willott et al. 2010).



# When do you make the first SMBHs?

The QSO ULASJ1 12010641  
at  $z=7.1$  has an estimated SMBH mass  
 $M_{\text{BH}} \sim 2 \times 10^9 M_{\odot}$  (Mortlock et al. 2011)

As massive as the largest  
SMBHs today, but when the  
Universe was only 0.75  
Gyrs old



Accretion time needed by a BH to reach a given final mass:

$$M(t) = M_0 e^{f_{\text{Edd}} \frac{t}{\tau} \frac{1-\epsilon}{\epsilon}}$$

$$t_{\text{acc}} = 0.45 \text{ Gyr} \frac{\epsilon}{1-\epsilon} f_{\text{Edd}}^{-1} \ln(M_{\text{fin}}/M_{\text{in}})$$

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$$M_{\text{fin}} = 10^9 M_{\odot}$$

$$M_{\text{in}} = 10^2 M_{\odot} \quad (10^5 M_{\odot})$$

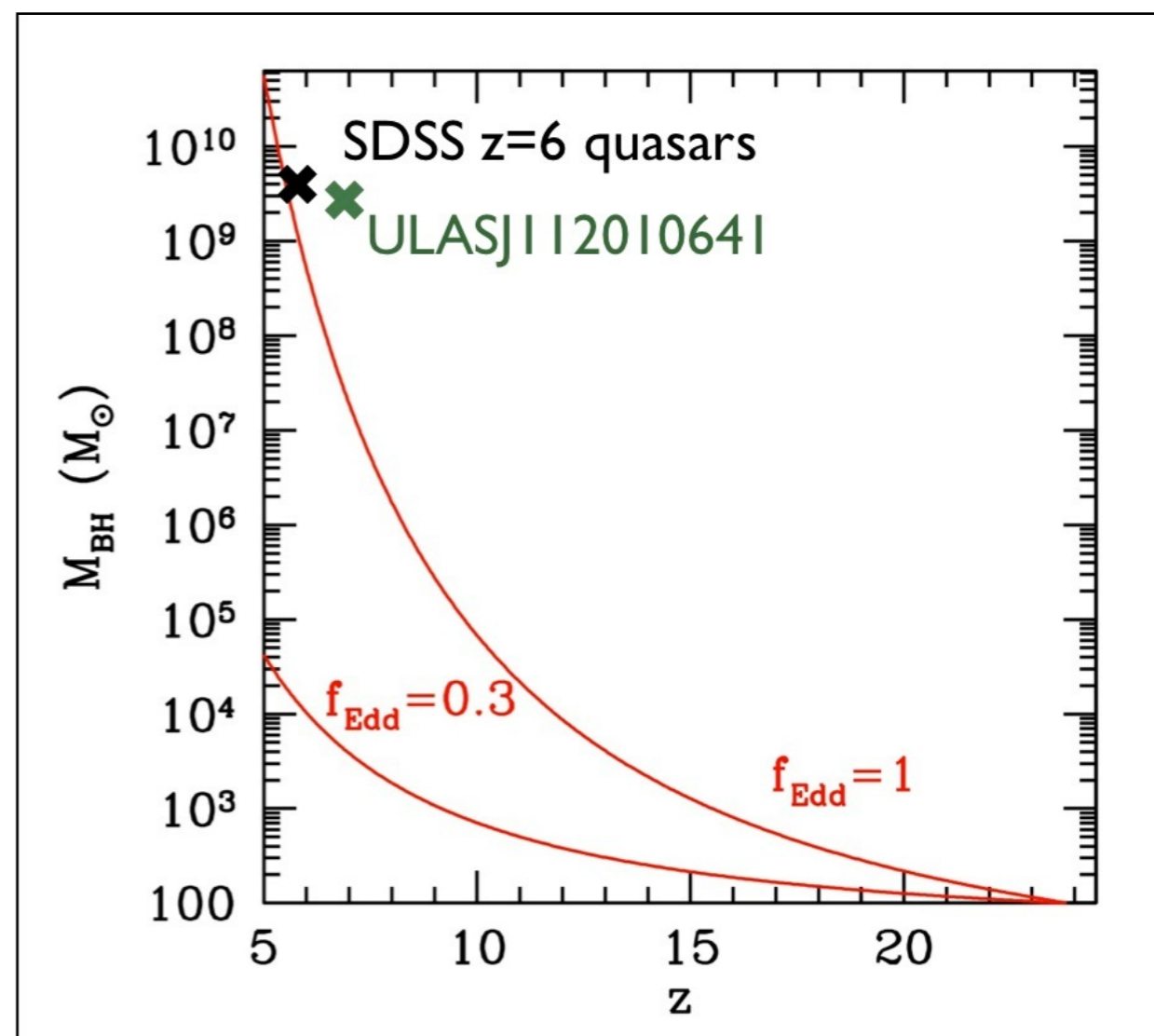
$$f_{\text{Edd}} = 1$$

$$\epsilon = 0.1$$

$$\Rightarrow t_{\text{acc}} = 0.8 \text{ Gyrs} \quad (0.45 \text{ Gyrs})$$

$$t_{\text{H}}(z=7) = 0.75 \text{ Gyrs}$$

$$t_{\text{H}}(z=6) = 0.9 \text{ Gyrs}$$





How do we make SMBH seeds?

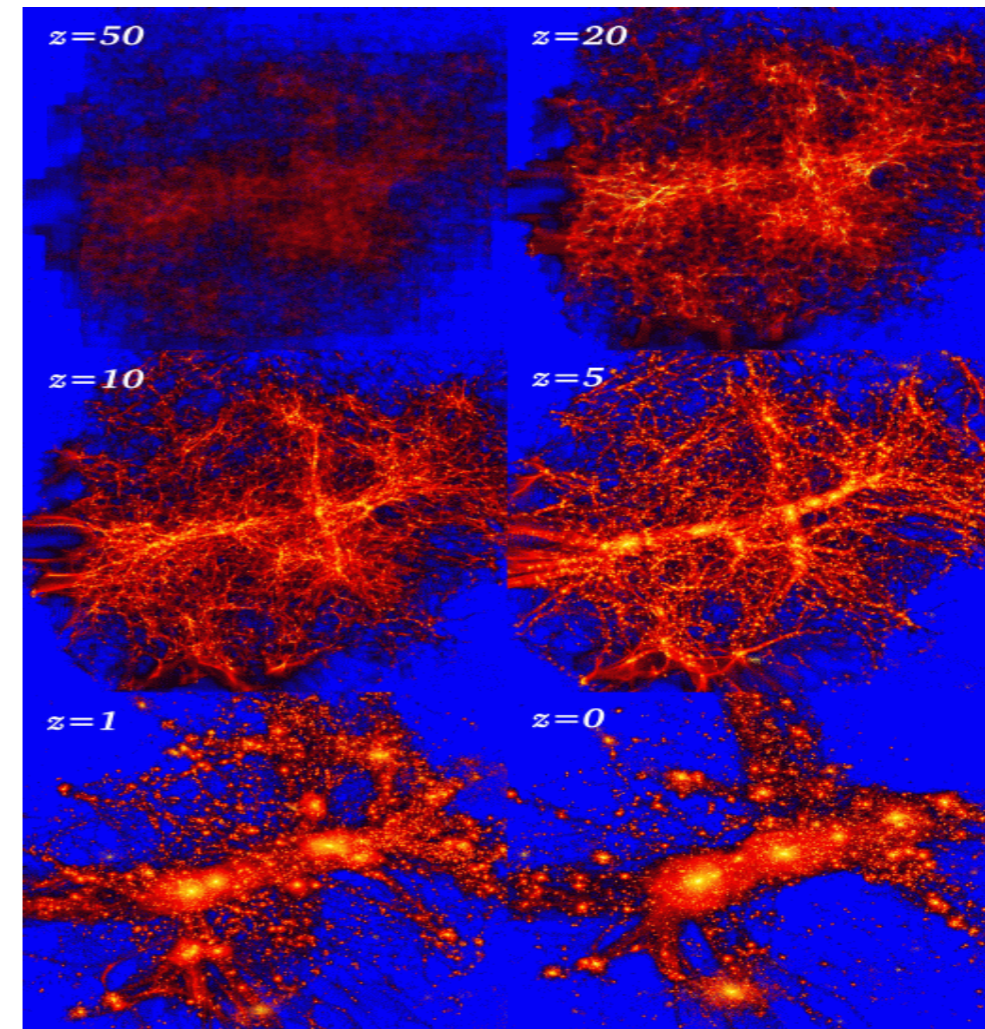
# DARK AGES

## The fate of DM:

CDM: Small scales collapse first  
BOTTOM-UP HIERARCHY

DM-HALOS COLLAPSE AS:

$$M = 10^{15} / (1+z)^6 M_{\odot}$$



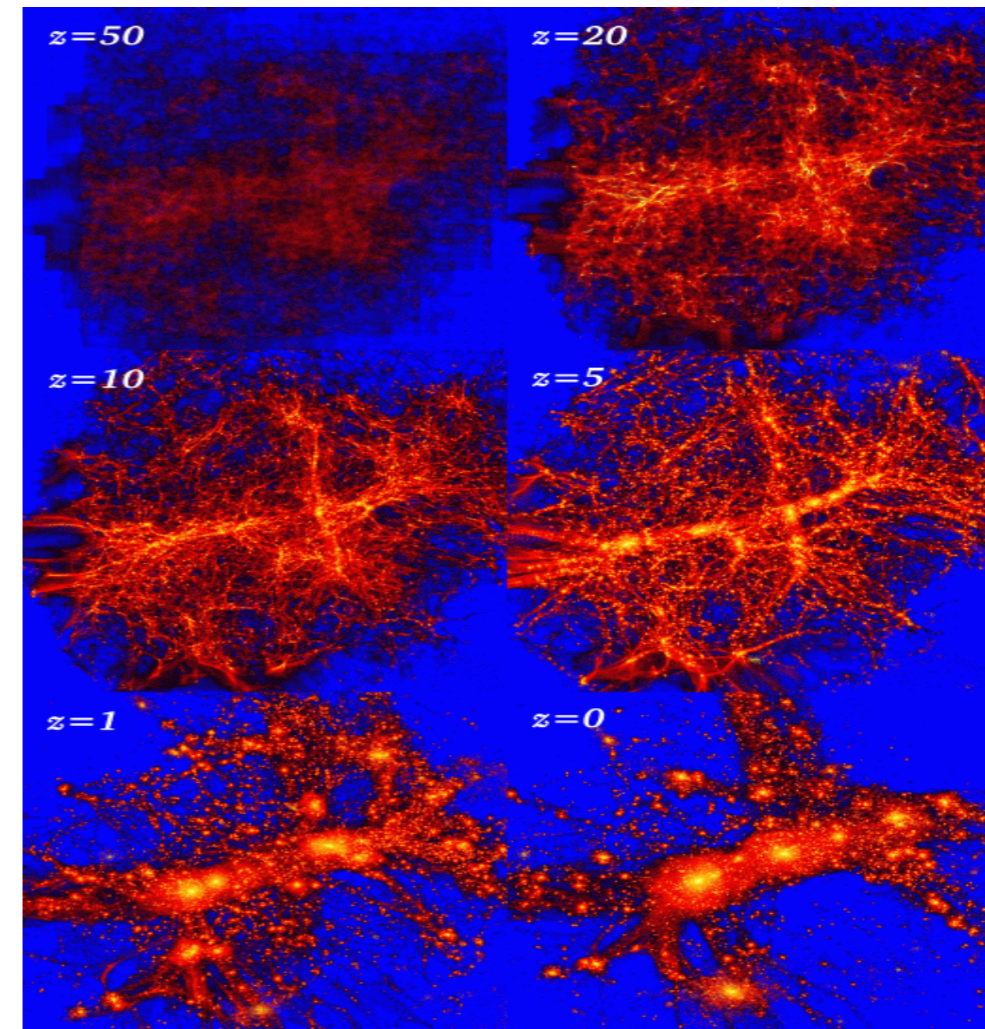
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## The fate of baryons:

At  $z \approx 130$  baryons are free to fall into DM halos

COOLING TIME  $\gg$  HUBBLE TIME  $\rightarrow$  ADIABATIC COLLAPSE

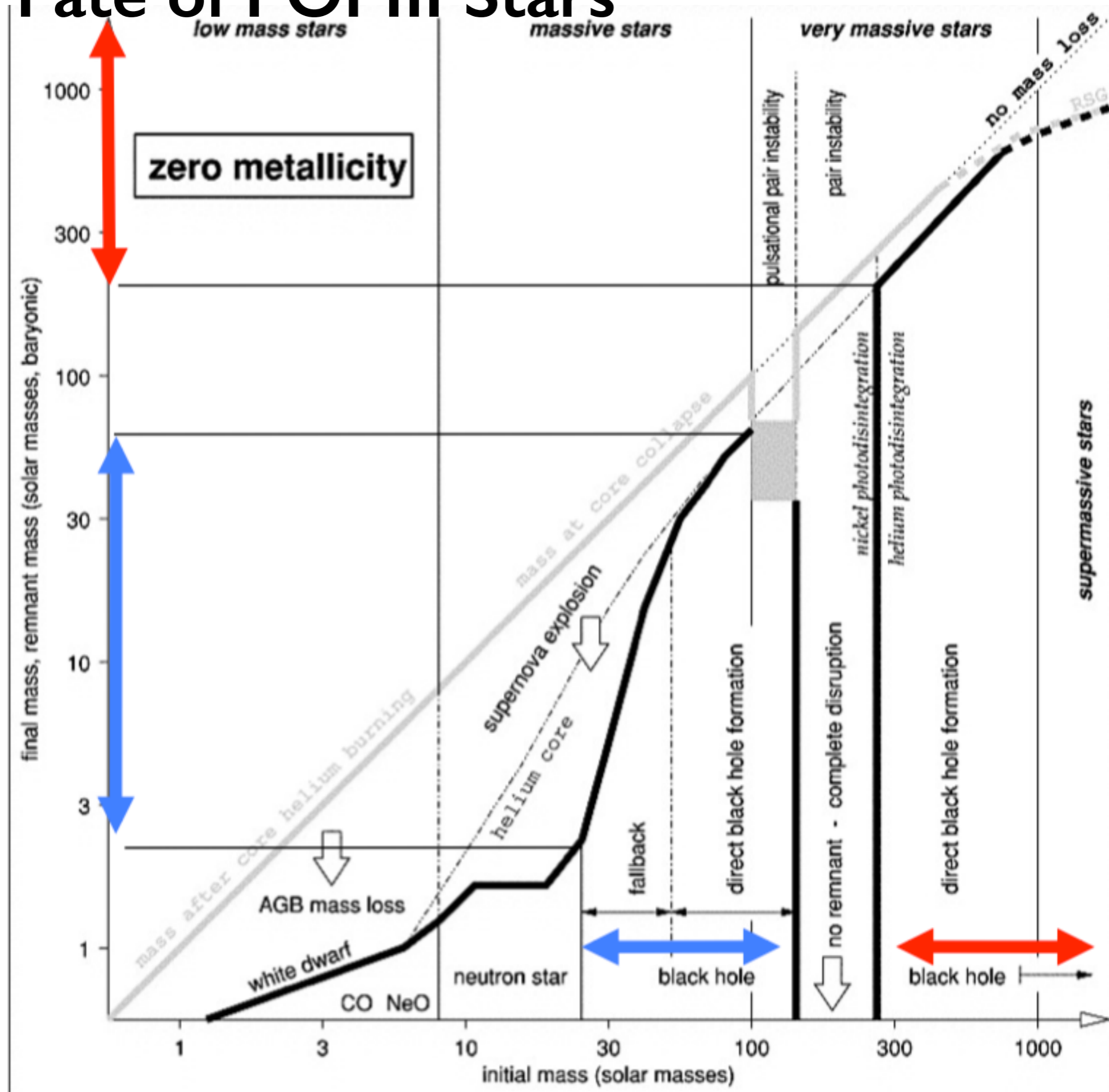
Baryons virialize as DM particles

COOLING TIME  $\ll$  HUBBLE TIME  $\rightarrow$  ISOTHERMAL COLLAPSE

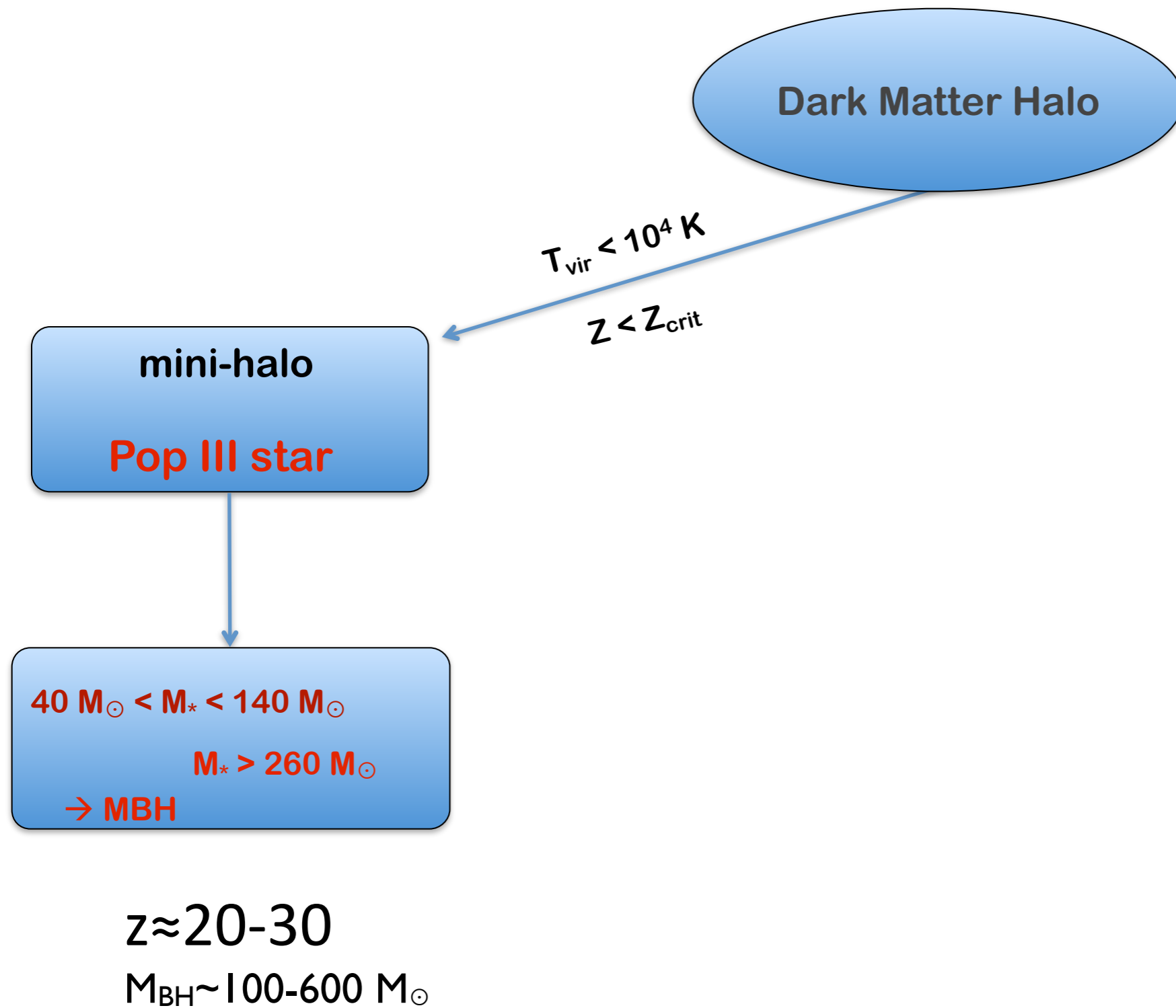
Baryons fall into DM potential wells  $\rightarrow$  Self-gravitating baryonic objects (POPIII)

$\rightarrow$  END OF THE DARK AGES ( $z \approx 25$ )

# The Fate of POP III Stars

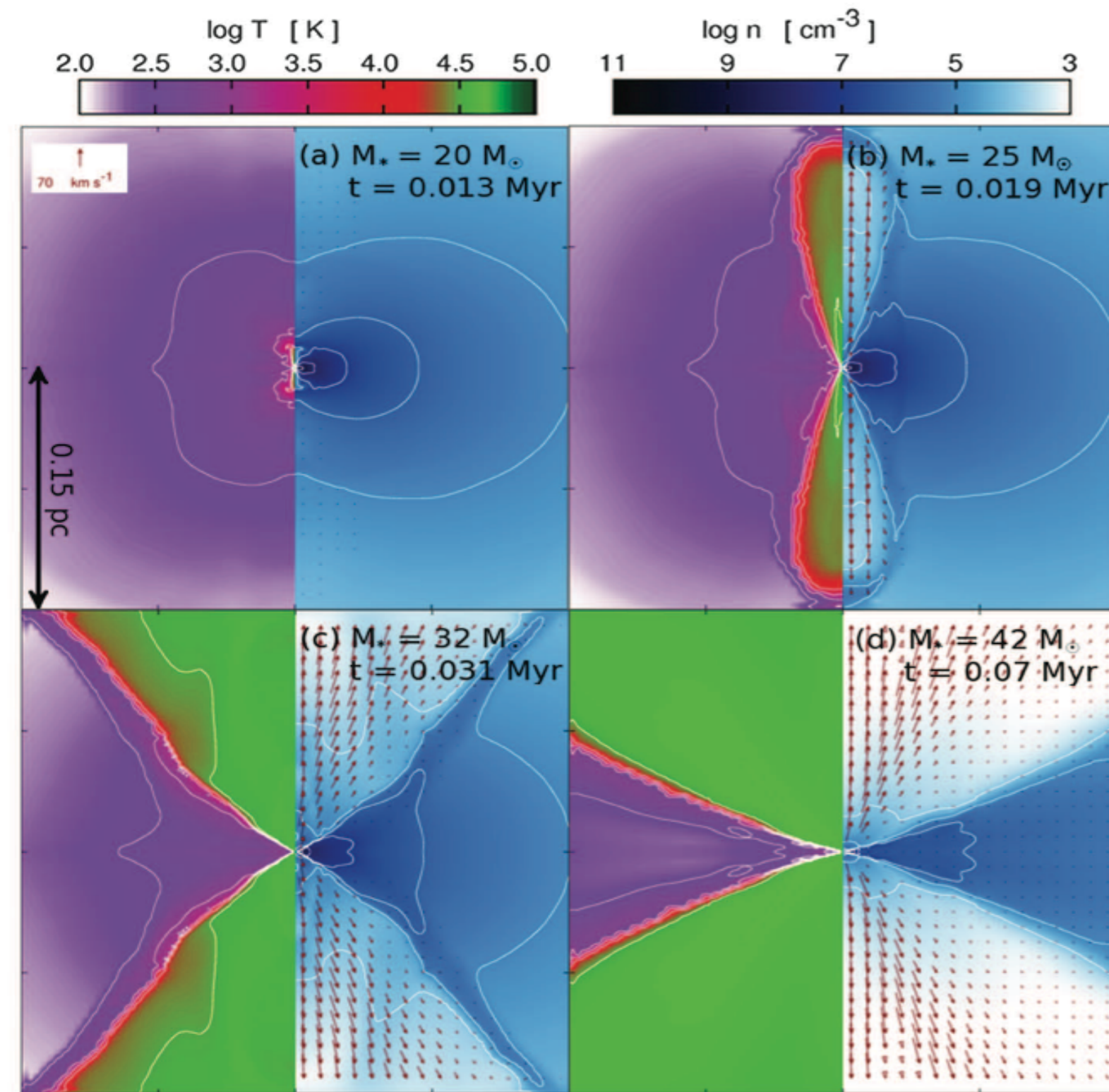


# The Nature of BH seeds



# POPIII Stars

Hosokawa et al. 2011



Protostar  $M \sim 0.01 M_{\odot}$  surrounded by  $1 M_{\odot}$  molecular envelope.

Accretion of envelope. Further out molecular disk ( $\sim 400$  AU)

Accretion through the disk at  $\sim 10^{-3} M_{\odot}/\text{yr}$

Ionizing bipolar flux (disk self-shielded)

HII region breaks the accreting envelope

Bipolar outflows of gas

Shocks propagates into the envelope

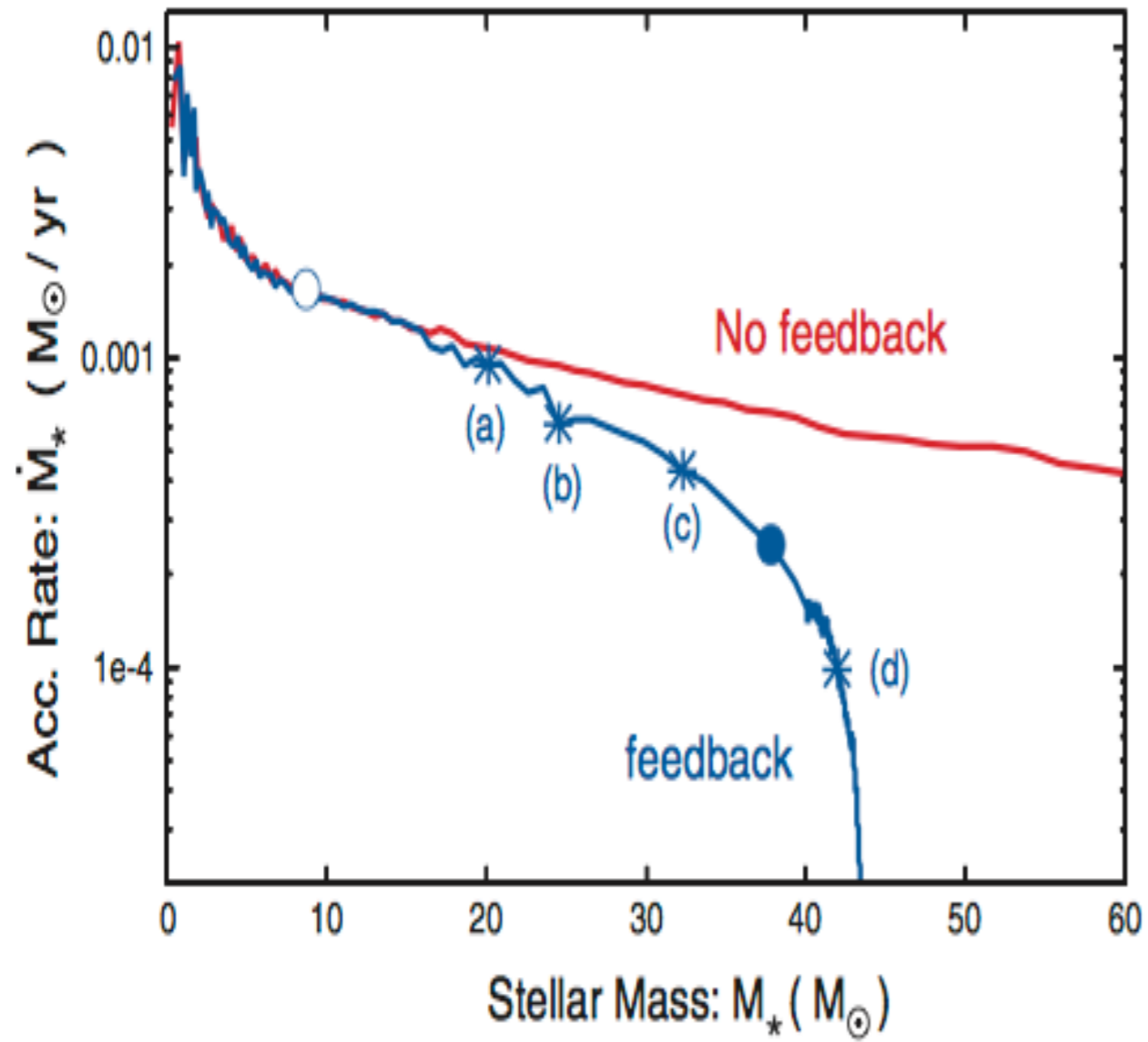
Accretion rate falls down

Disk exposed to UV. Photoevaporation.

Accretion stops.  $M \sim 45 M_{\odot}$ . Limit set by self-regulating radiation feed-back

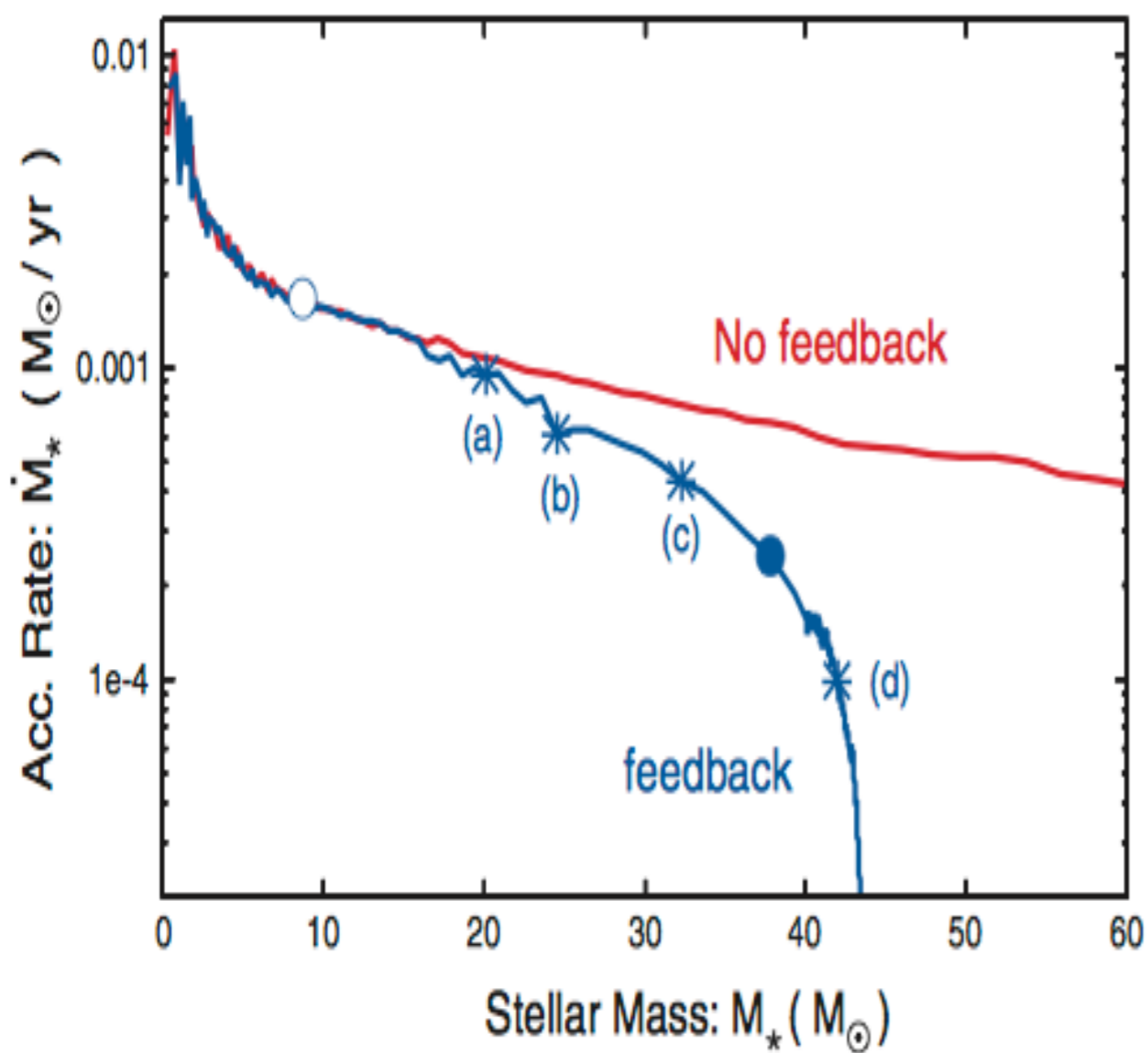
**Fig. 2.** UV radiative feedback from the primordial protostar. The spatial distributions of gas temperature (left), number density (right), and velocity (right, arrows) are presented in each panel for the central regions of the computational domain. The four panels show snapshots at times when the stellar mass is  $M_{\star} = 20 M_{\odot}$  (A),  $25 M_{\odot}$  (B),  $35 M_{\odot}$  (C), and  $42 M_{\odot}$  (D). The elapsed time since the birth of the primordial protostar is labeled in each panel.

# The Fate of POPIII Stars: Feedback Effects

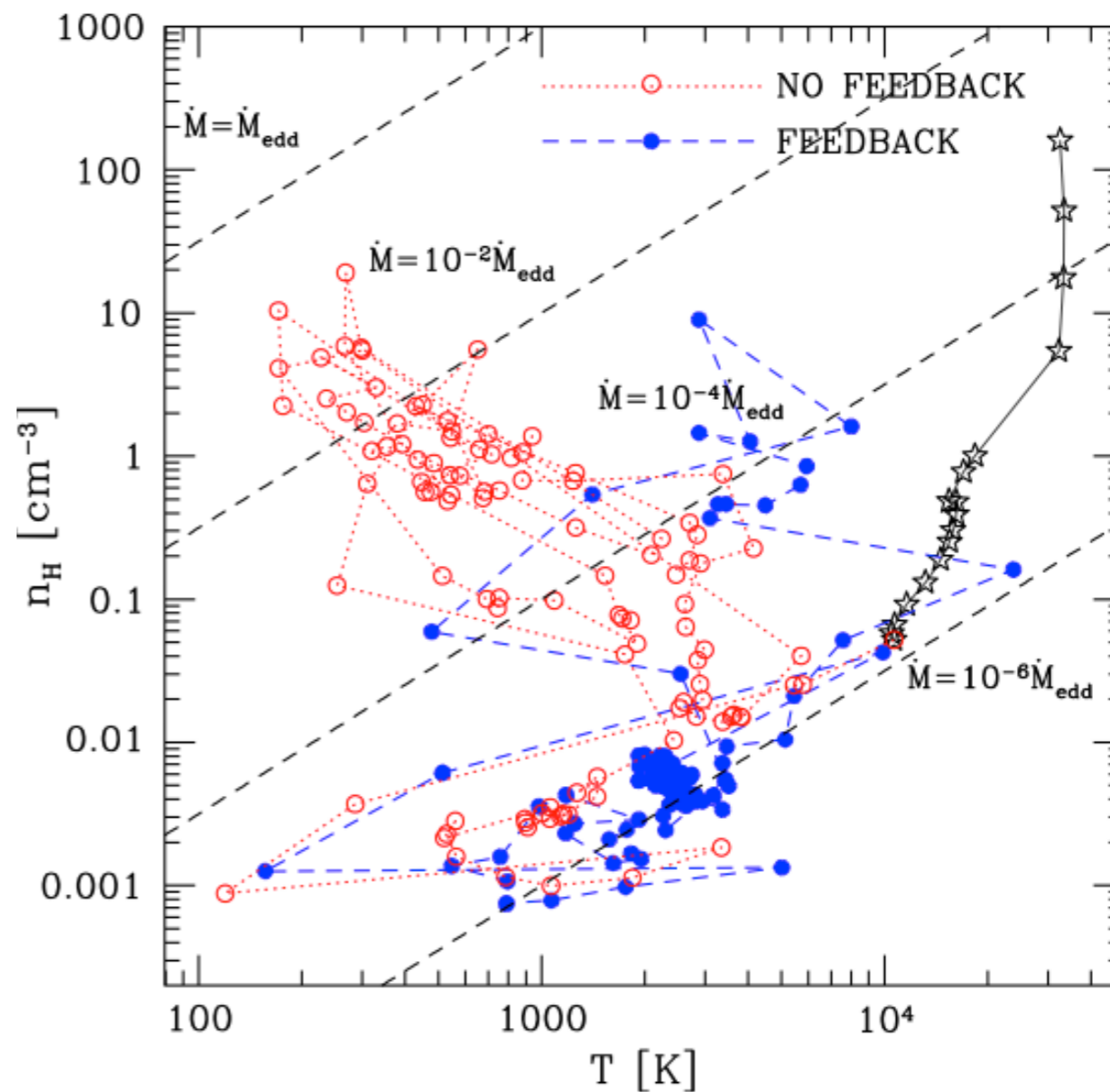


Hosokawa et al. 2011

# The Fate of POPII Stars: Feedback Effects



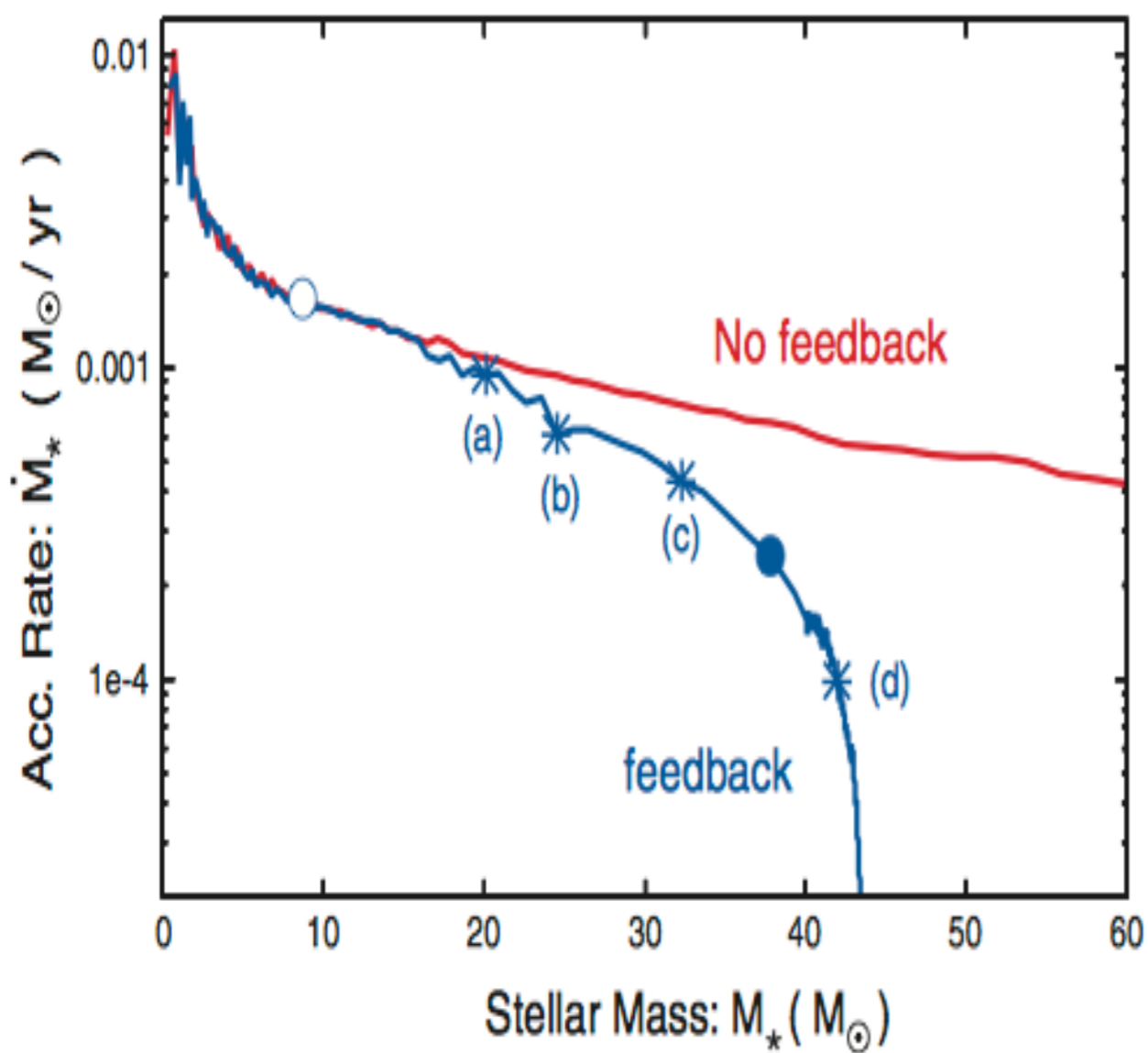
Hosokawa et al. 2011



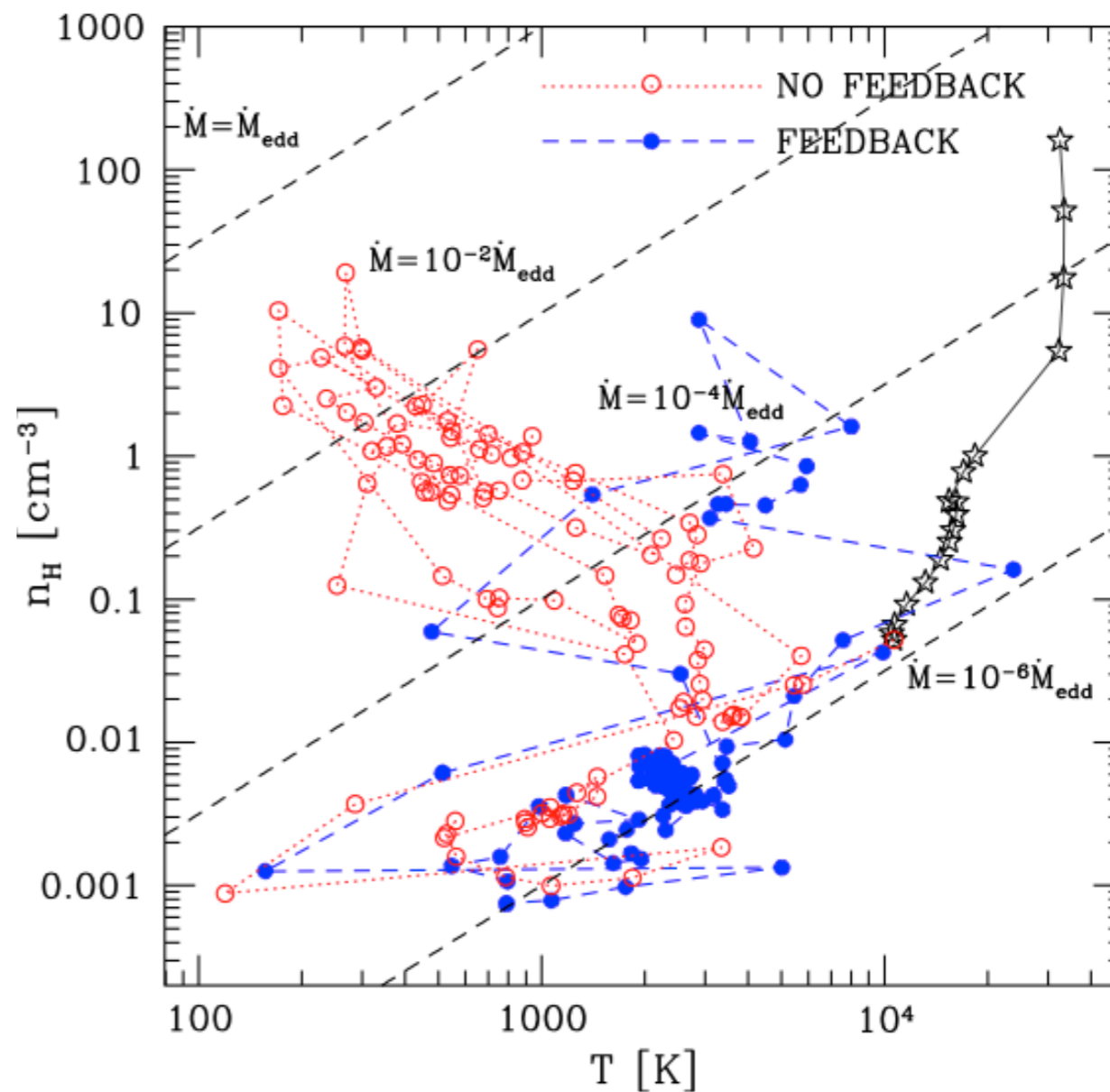
Alvarez et al. 2009



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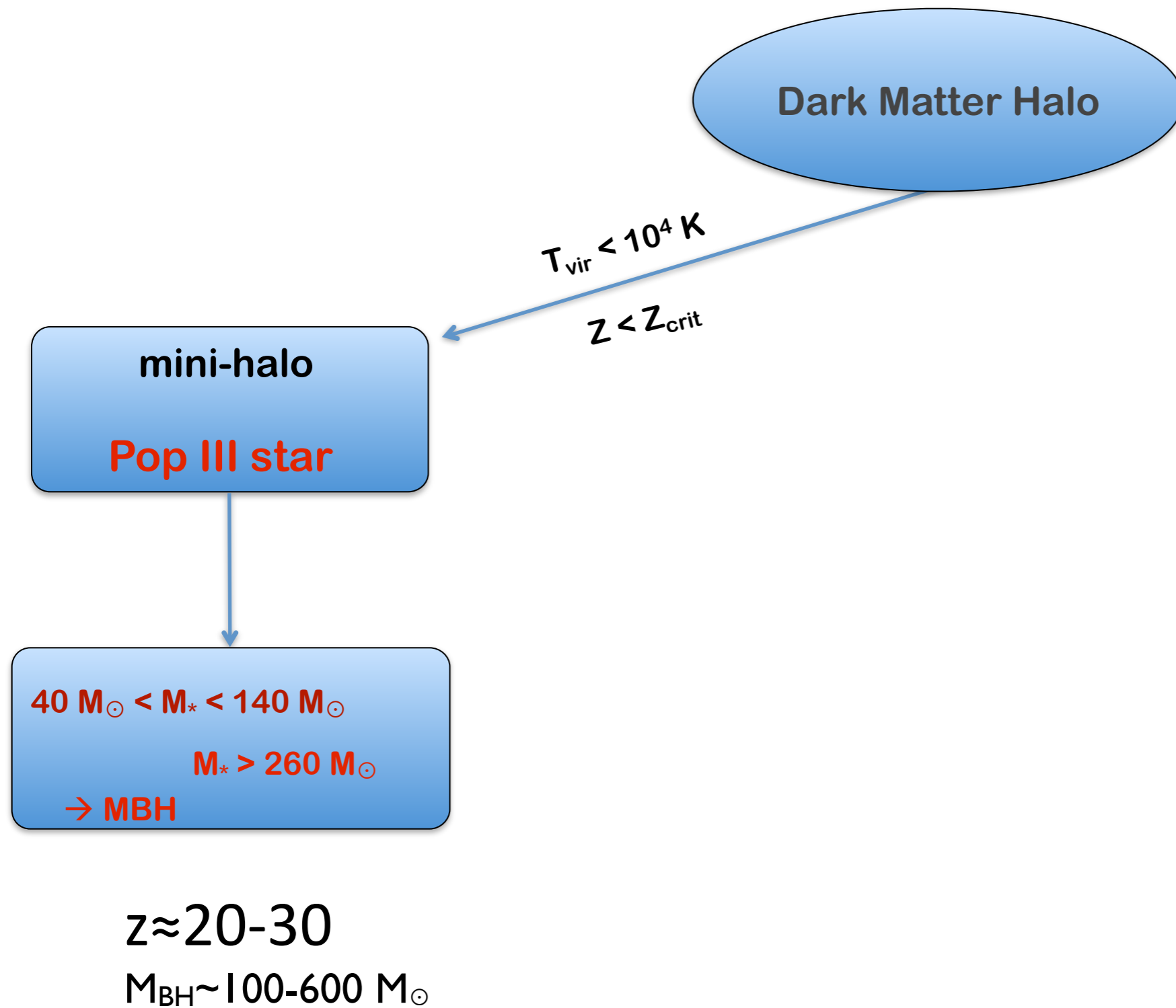
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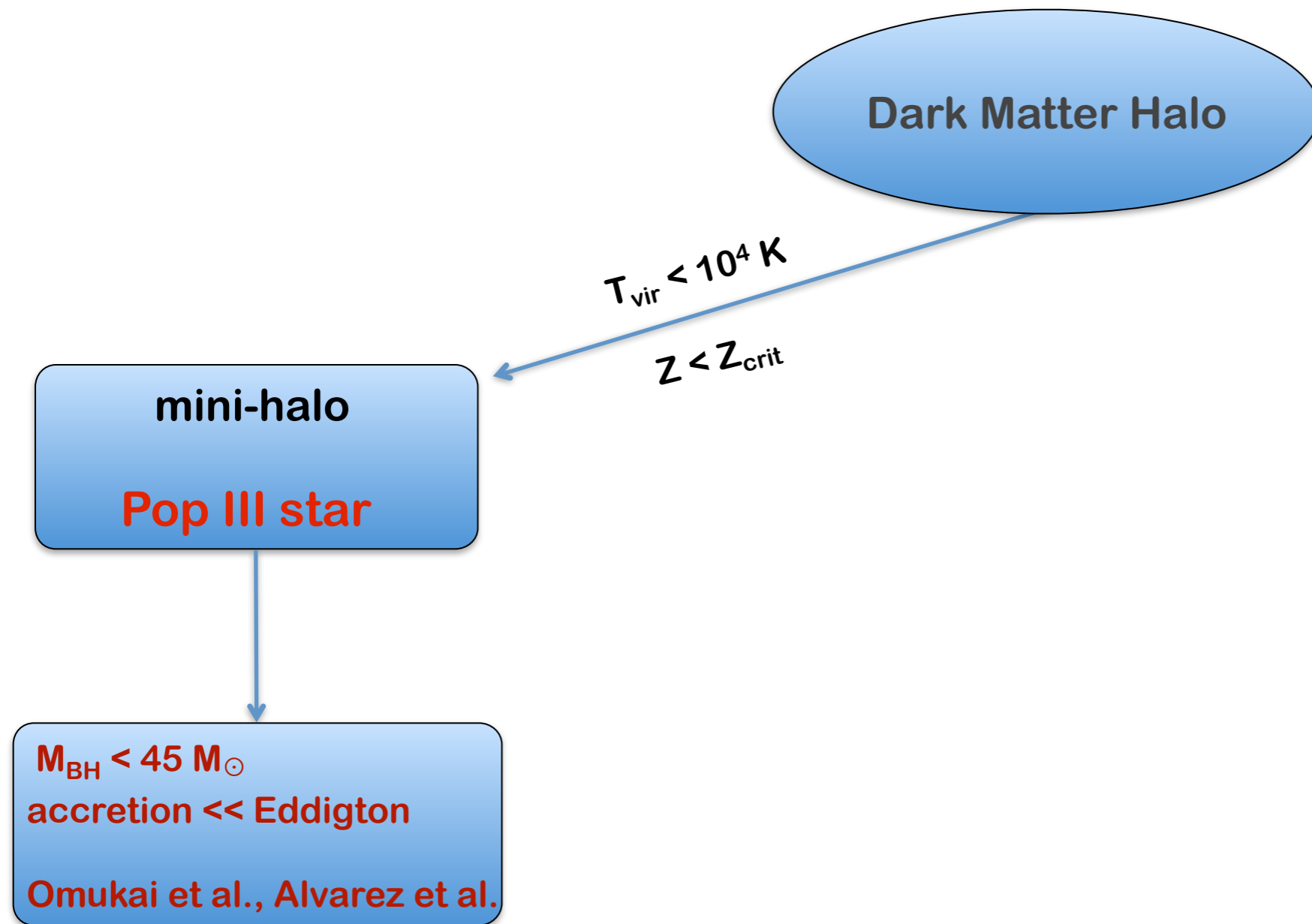
Alvarez et al. 2009

POPIIIs: massive but not very massive ( $\sim 45 M_\odot$ ).  
 Accretion rate onto the remnant: very sub-Eddington.

# The Nature of BH seeds



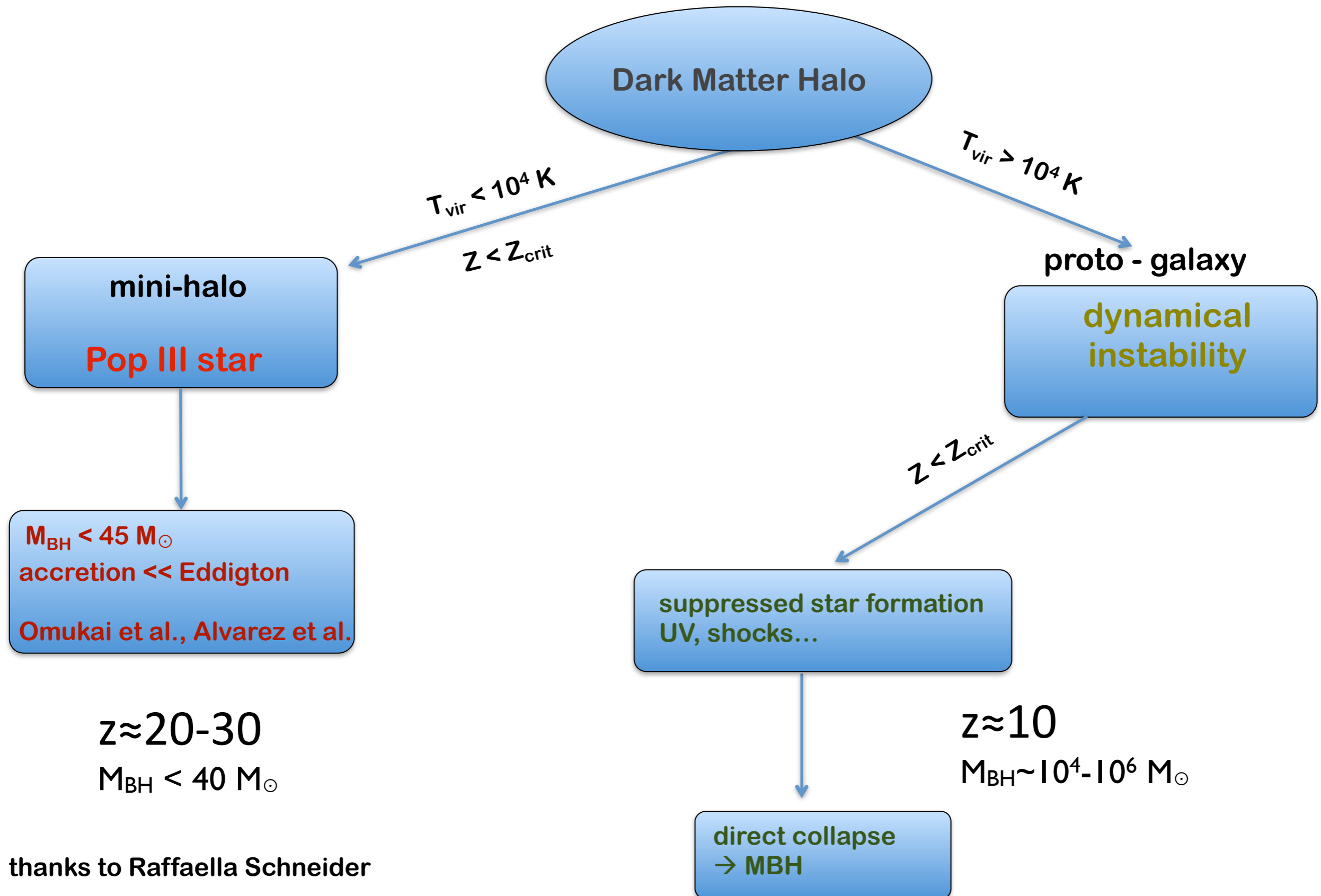
# The Nature of BH seeds



$z \approx 20-30$

$M_{\text{BH}} < 40 M_{\odot}$

# The Nature of BH seeds



# Direct Gas Collapse

e.g., Bromm & Loeb 2003, Begelman, Volonteri & Rees 2006

- Deep potential well for gas infall and collapse
- Global dynamical instabilities to trigger inflow and dissipate angular momentum
- Inflow and formation of a super-massive star (quasistar) that collapses into a BH

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## Problems

- Angular momentum dissipation
- Star Formation along the way down

# Direct Gas Collapse

- Efficient gas collapse in halos with  $T_{\text{vir}} > 10^4$  K (metal free and no  $\text{H}_2$  cooling)
- No fragmentation and star formation (Haehnelt & Rees 1993, Begelman et al. 2006)
- Atomic cooling down to 4,000 K (collisionally excited Ly $\alpha$  emission), isothermal collapse, then adiabatic
- Local Lyman-Werner UV to prevent  $\text{H}_2$  formation (o.w. turbulent medium)
- Collapsing gas settles into a rotationally supported disk.
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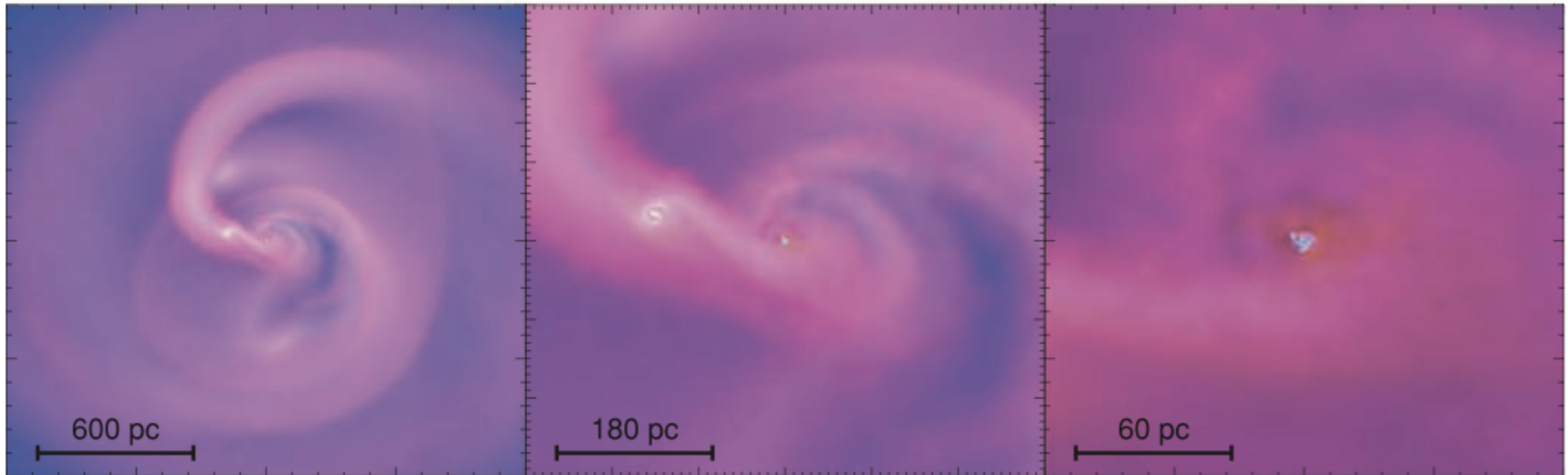
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- Angular momentum problem solved if: a) DM halo with very low spin ( $\lambda_{\text{spin}} \ll 0.04$ ); b) presence of low angular momentum gas; c) onset of global dynamical instabilities, **bars-within-bars** (Shlosman et al. 1989 Begelman et al. 2006).
- A bar can transport angular momentum outwards on a dynamical timescale via gravitational and hydrodynamical torques, allowing the radius to shrink.
- If gas can cool, this shrinkage leads to even greater instability on shorter timescale, and the process cascades

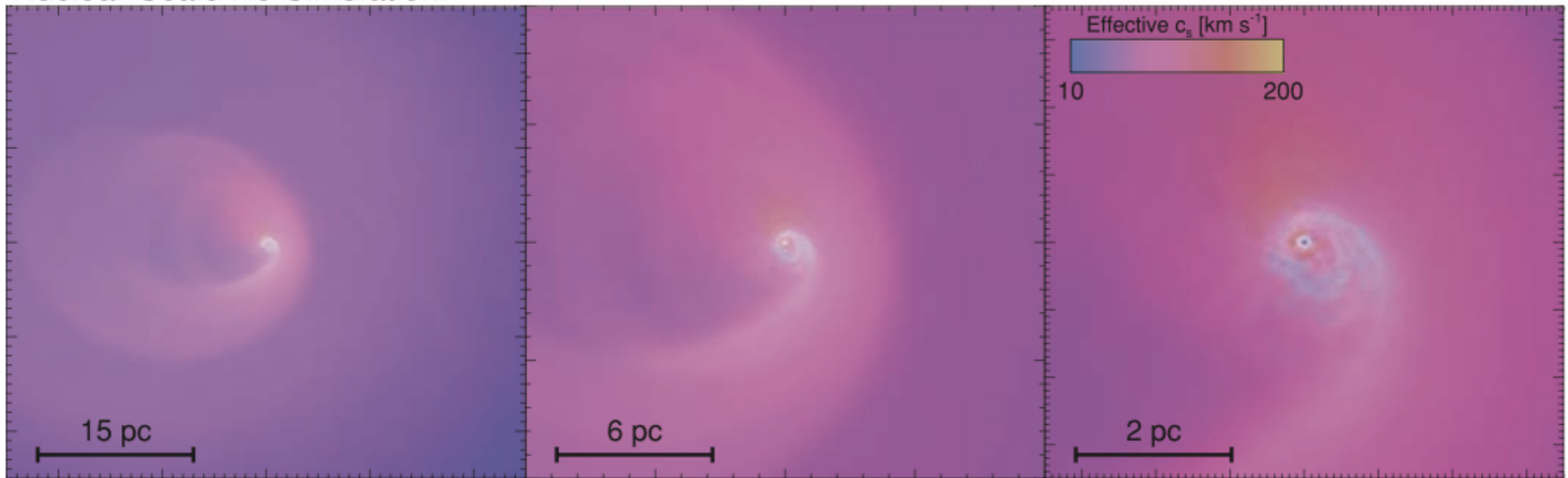


Hopkins &amp; Quataert 2010

## Intermediate-Scale Re-Simulation:



## Nuclear-Scale Re-Simulation:



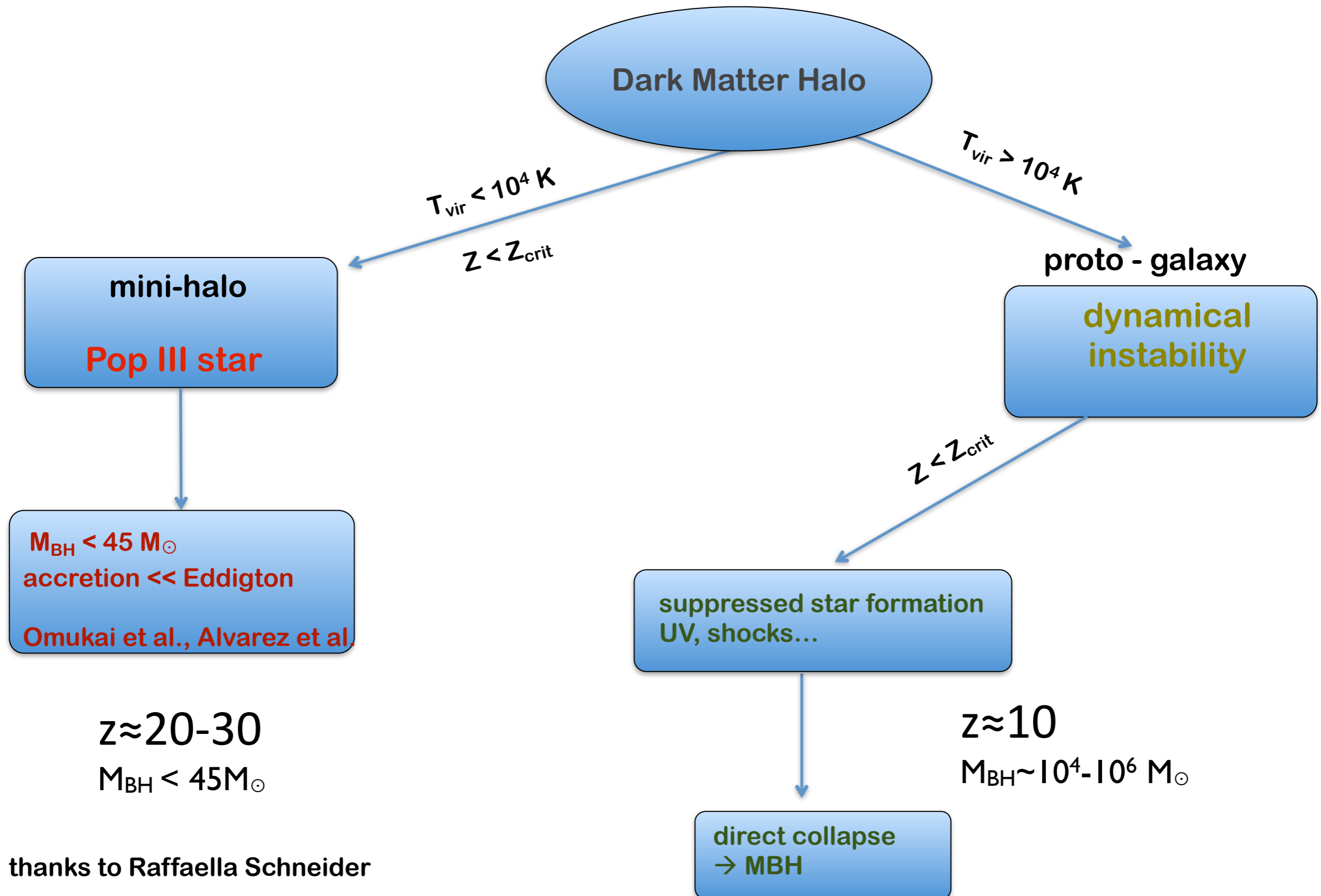
- If this all works, typical gas masses within the central few pc can reach  $10^4$ - $10^5 M_{\odot}$
- A Very Massive Star forms with  $M_{\text{star}} \sim 10^4 M_{\odot}$ . Numerical full-GR calculations predict the core of the VMS collapses into a few  $M_{\odot}$  Kerr BH (via neutrino cooling). A *Quasistar* is formed: a stellar-size BH embedded into a very massive stellar envelope
- Ultrafast accretion at Eddington rate (where mass is the Quasistar mass)
- The  $10 M_{\odot}$  BH can swallow 1-10% of the Quasistar envelope: MBH reaches  $10^4$ - $10^5 M_{\odot}$

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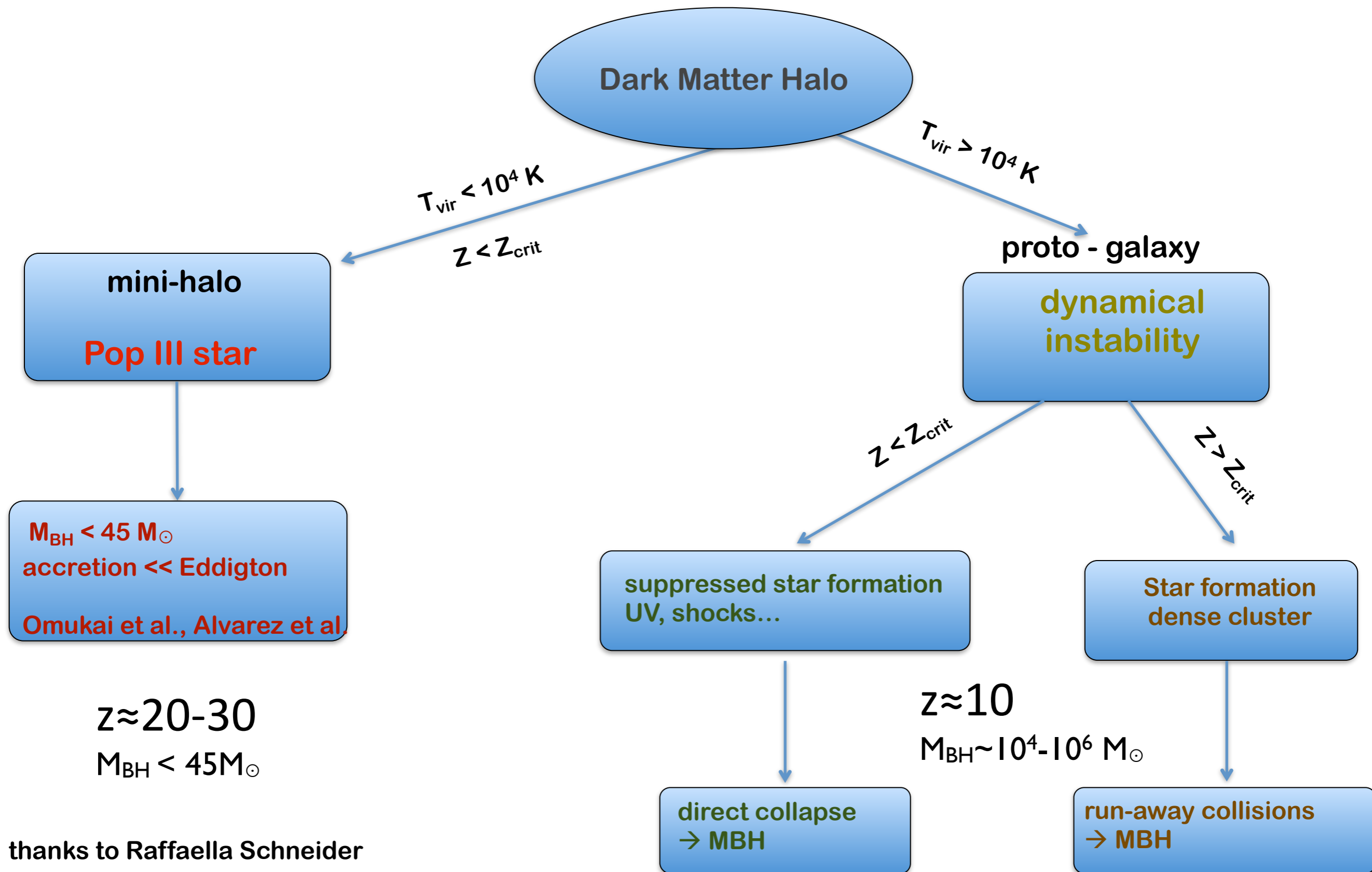
Begelman et al. 2006

Instead of going into BH formation, collapsing gas can fragment and form stars

# The Nature of BH seeds



# The Nature of BH seeds



# Star Formation and Run-Away Collapse

e.g., Devecchi & Volonteri 2009

- If gas is metal enriched by first POpIII stars can not avoid fragmentation
- Local dynamical instabilities lead to mass infall
- The inflow produces an inner, compact, very dense star cluster
- Mass segregation: more massive stars sink to the centre of the cluster
- Stellar collisions form a VMS
- The VMS accretes gas. BH remnant of  $10^3 M_{\odot}$  (Begelman & Rees 1978, Freitag et al. 2006)

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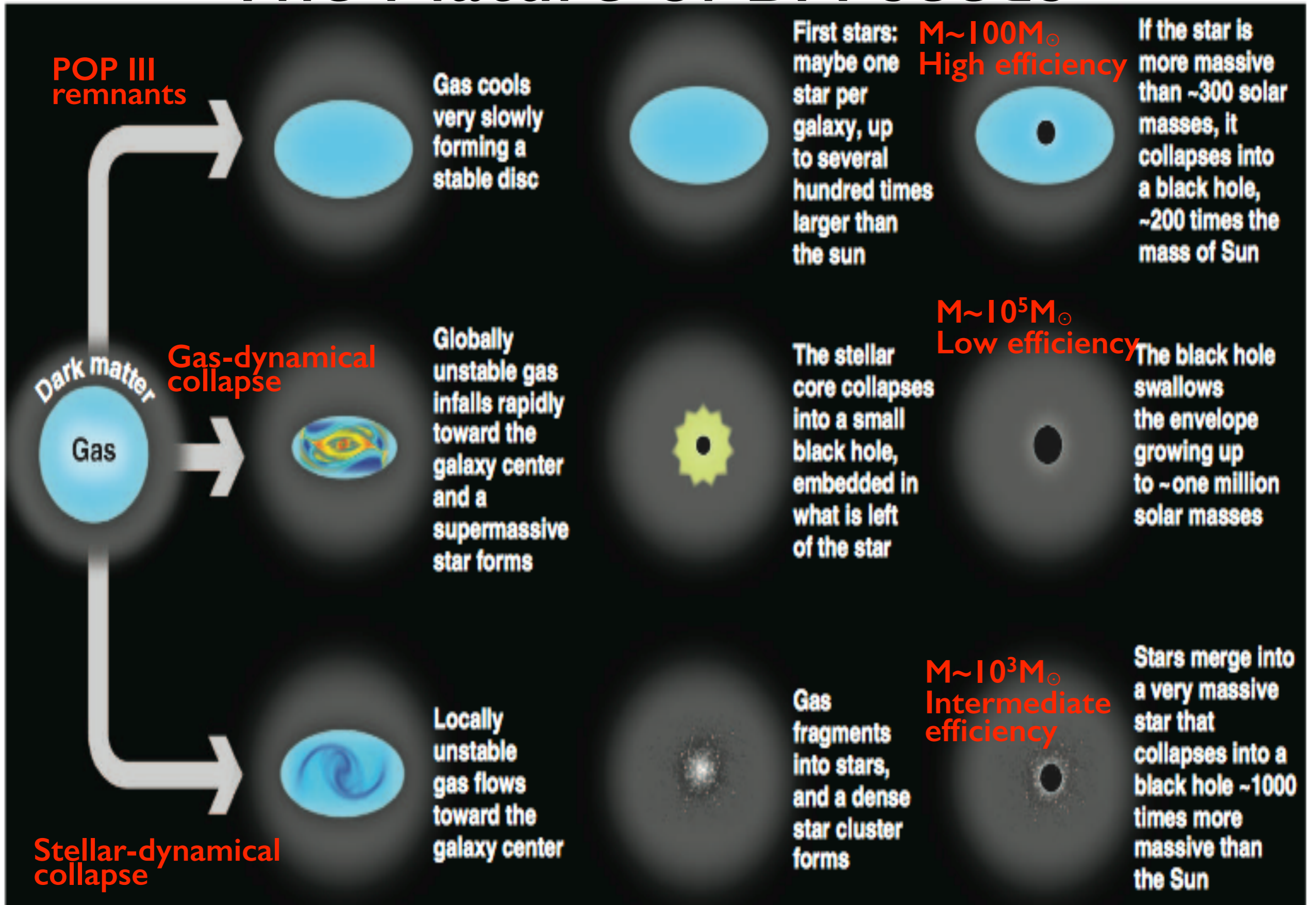
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- This route happens above a critical metallicity:  $Z > 10^{-5} Z_{\odot}$ . But  $Z$  needs not to be too high to avoid excessive gas consumption by star formation
- Time for mass segregation must be  $\ll$  main-sequence timescale (3 Myrs) to avoid SN driven mass-loss, and formation of compact objects (small cross section)

# The Nature of BH seeds

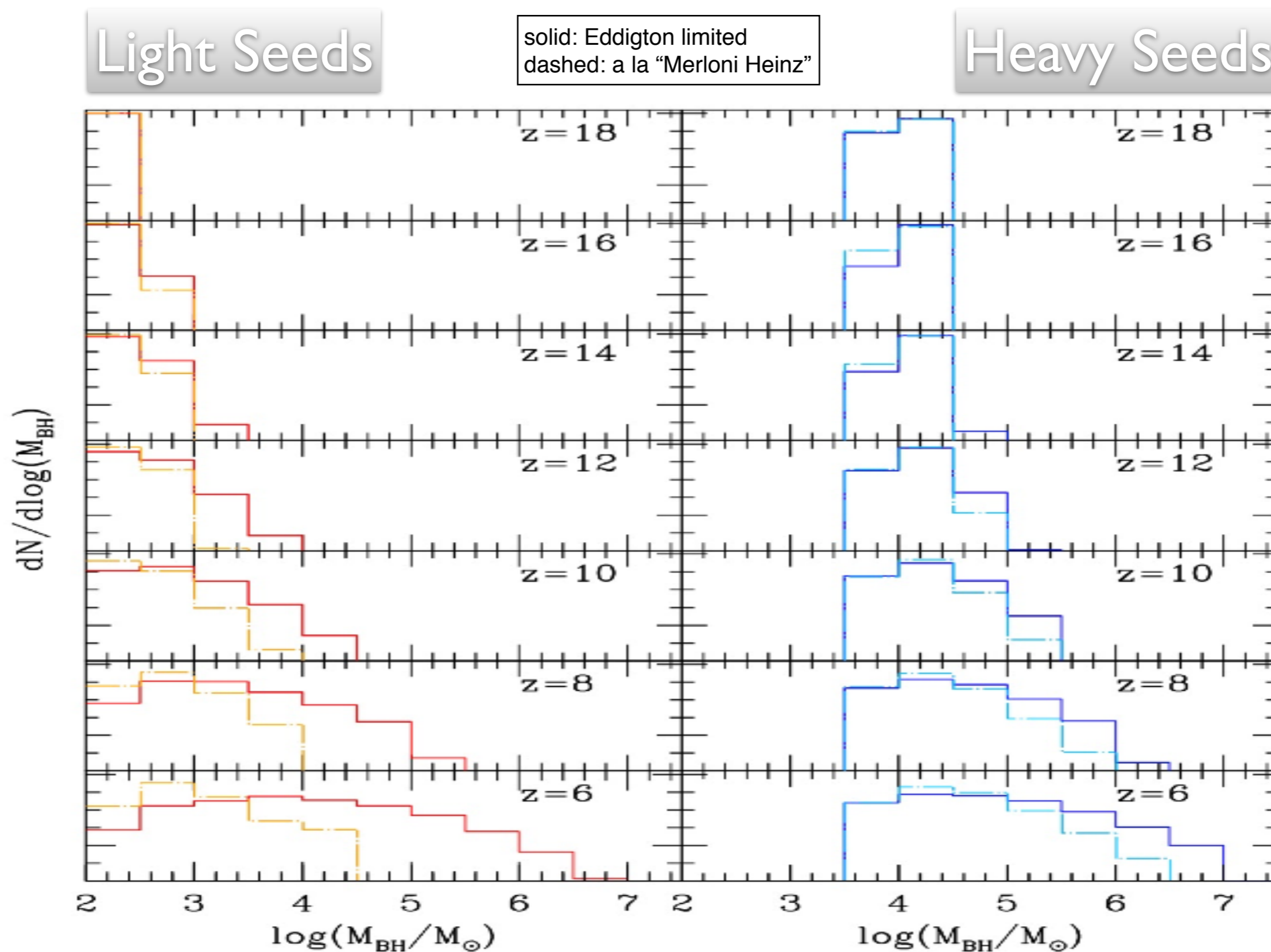
from Volonteri 2010





# Low Mass Seeds vs High Mass Seeds

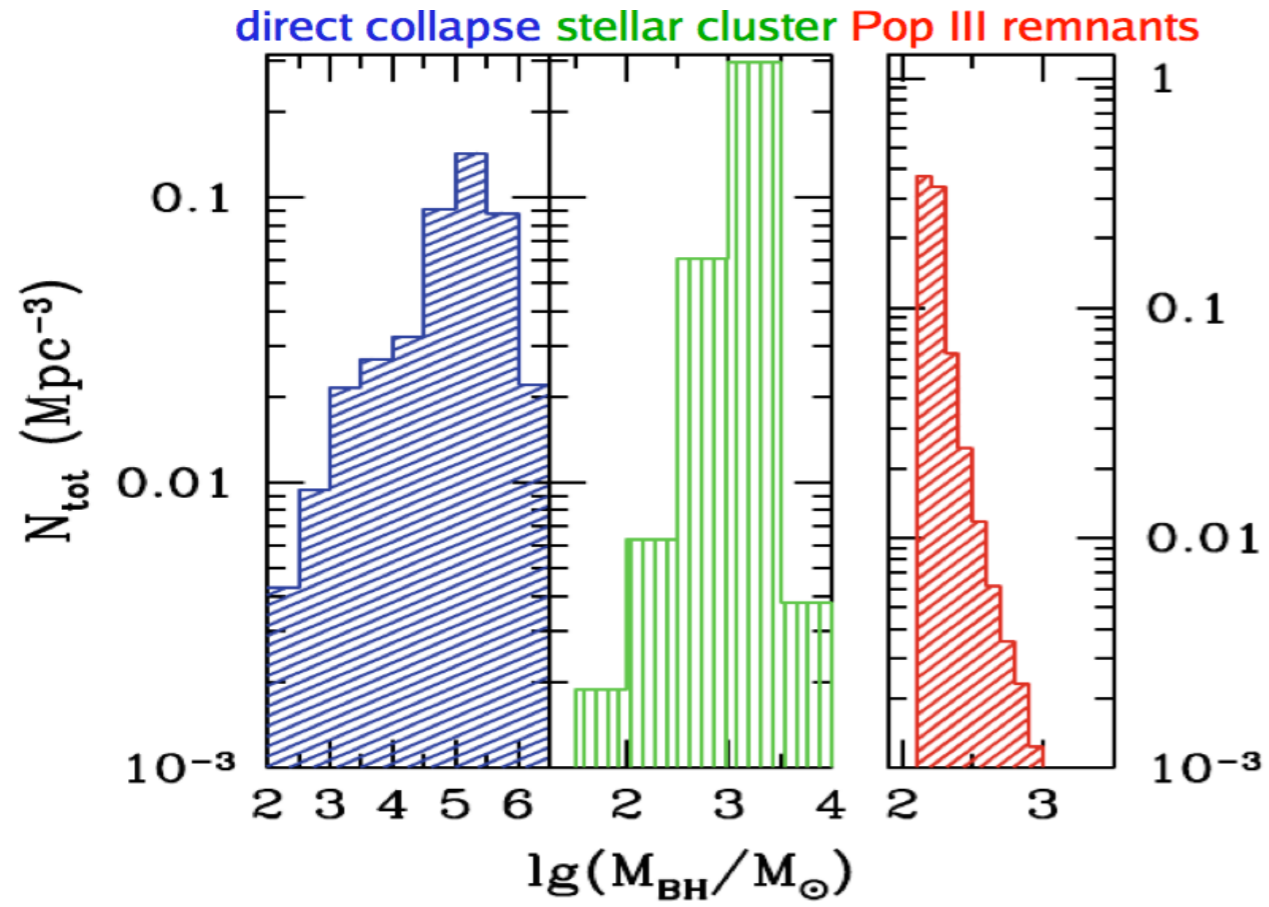
Volonteri &amp; Gnedin 2009



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Volonteri 2010

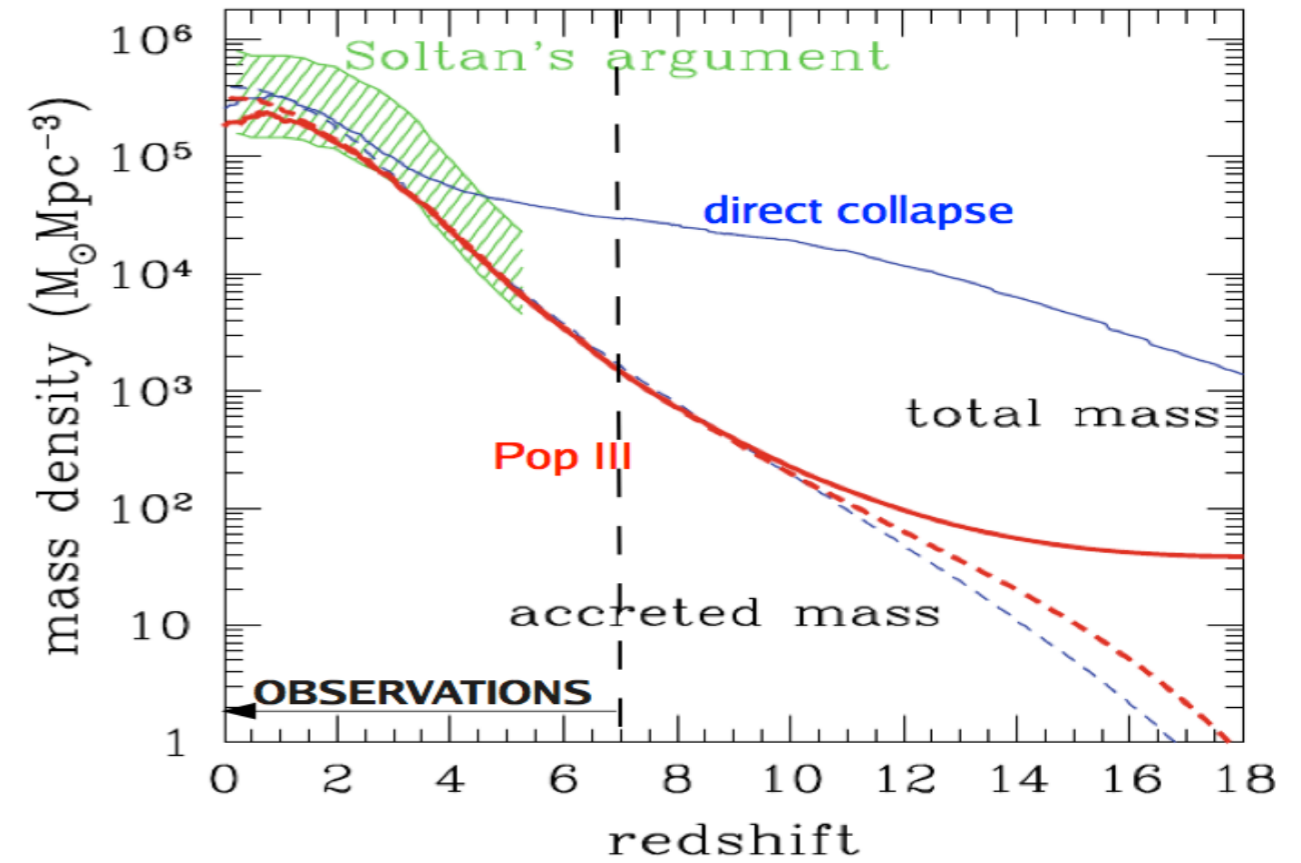
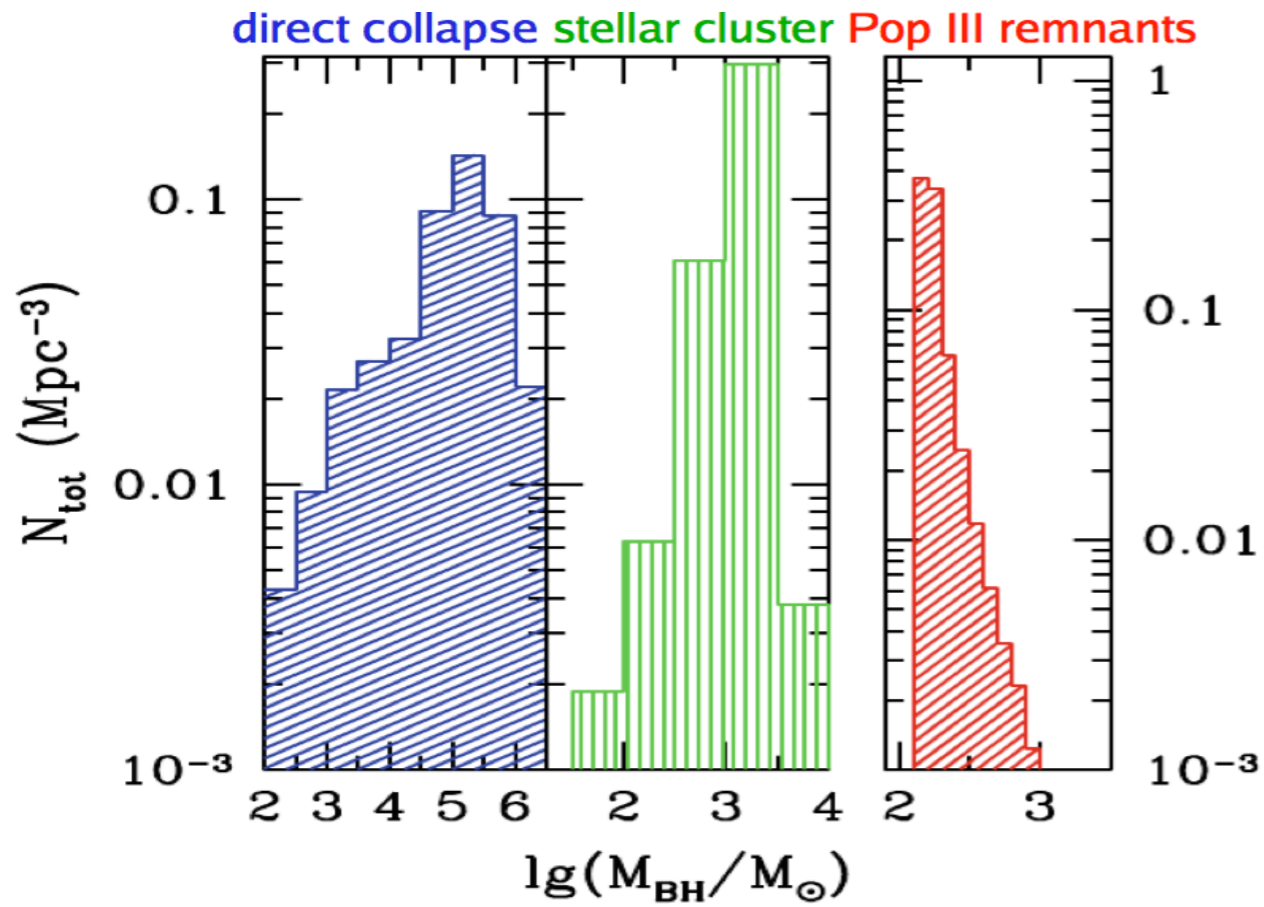
Madau &amp; Rees 2001



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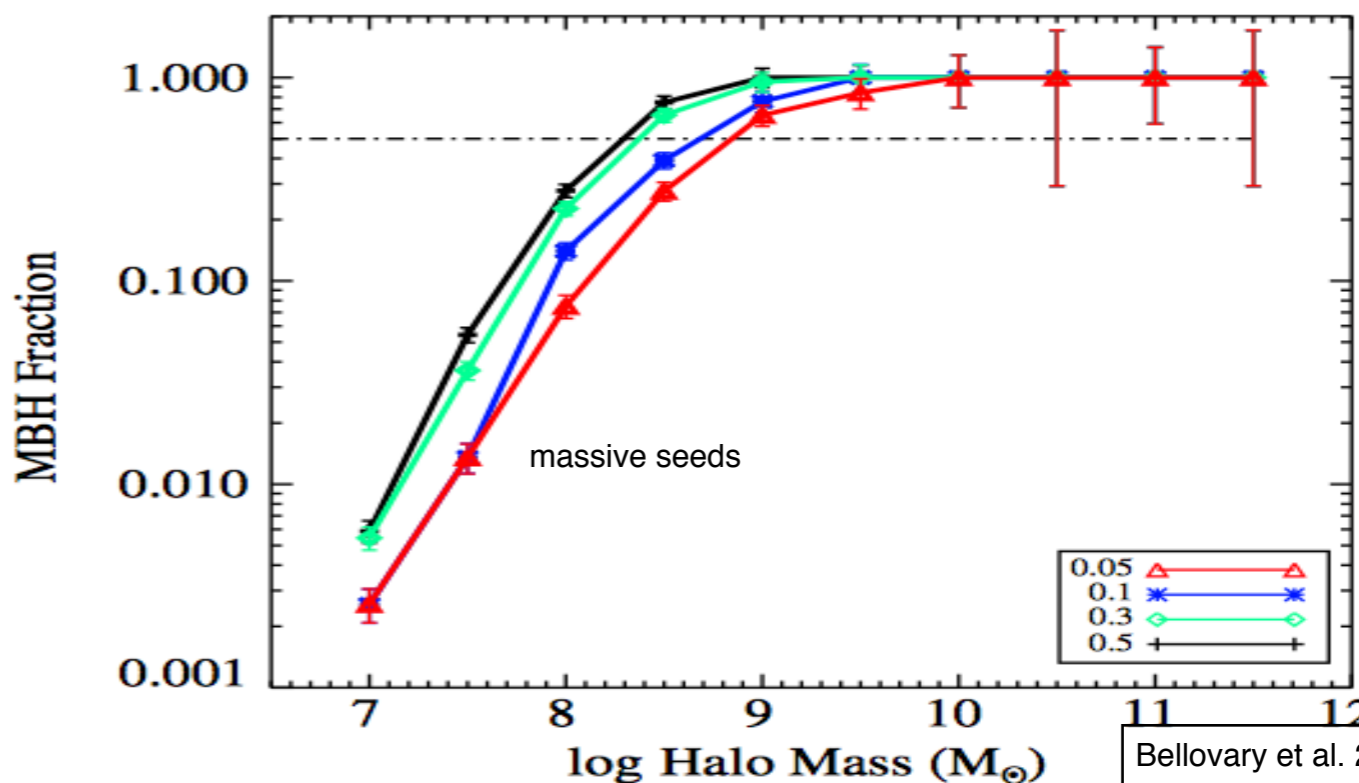
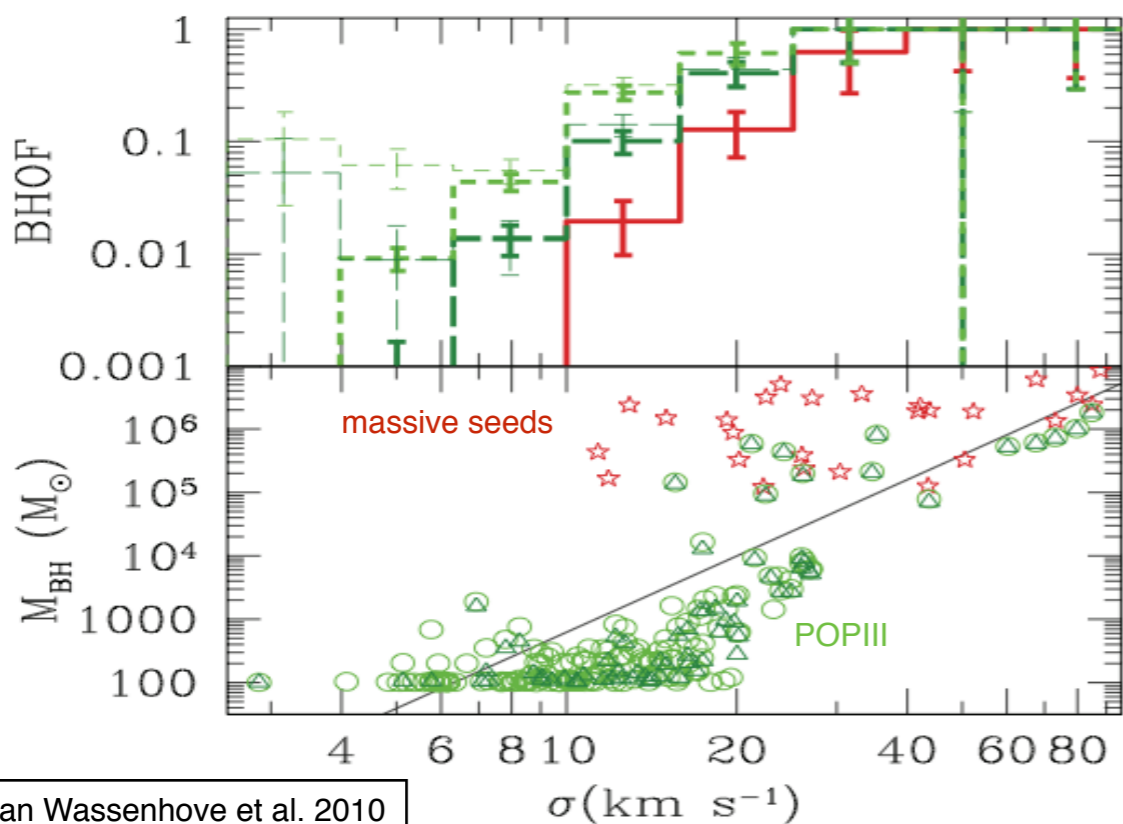
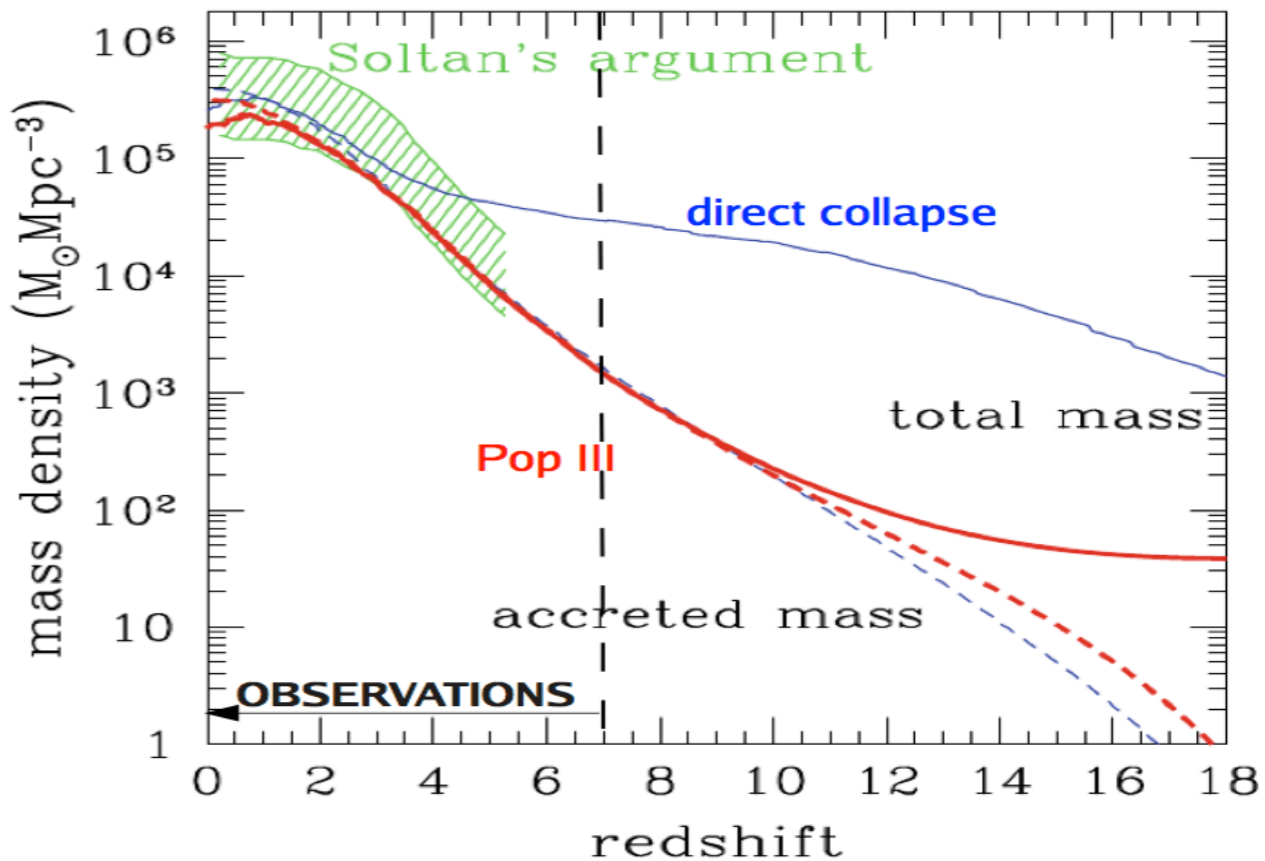
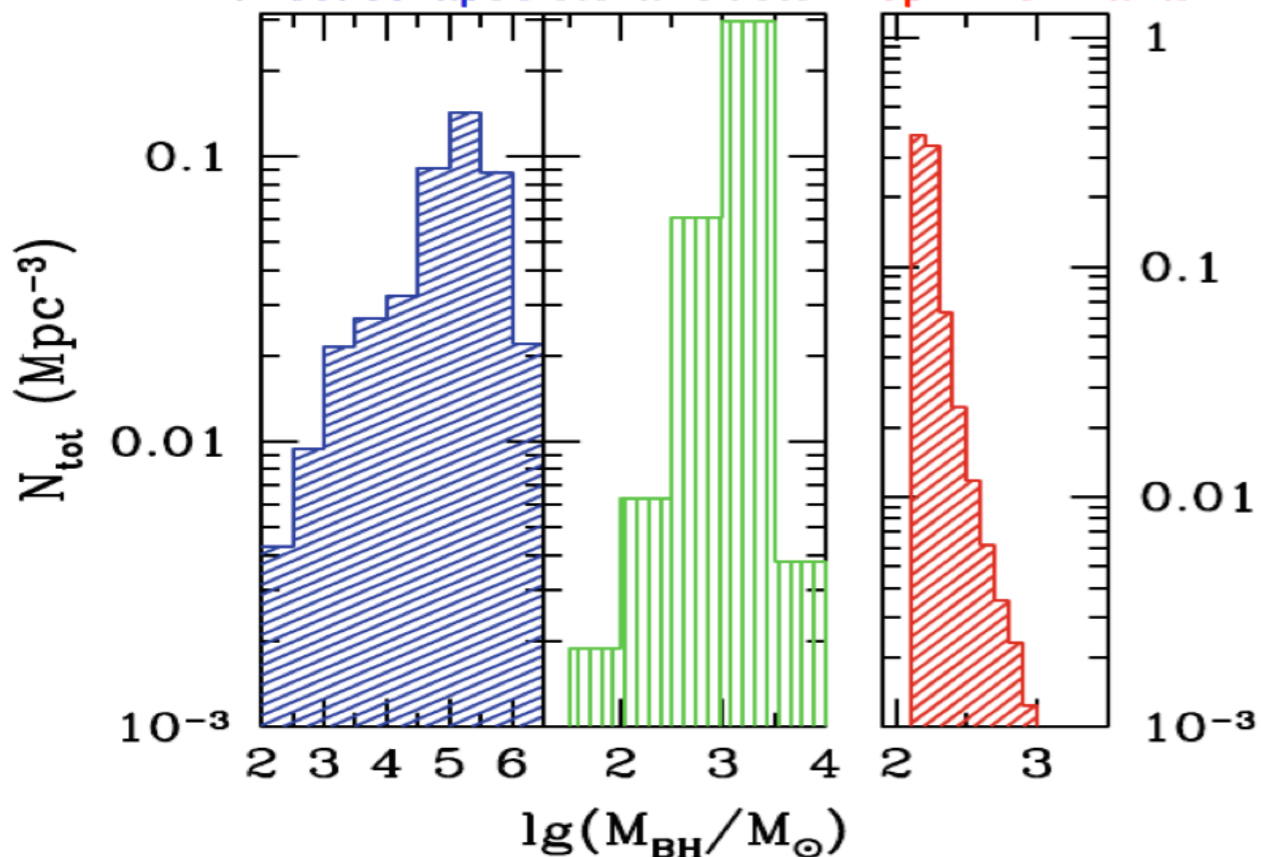


# Low Mass Seeds vs High Mass Seeds

Volonteri 2010

Madau & Rees 2001

direct collapse stellar cluster Pop III remnants



van Wassenhove et al. 2010

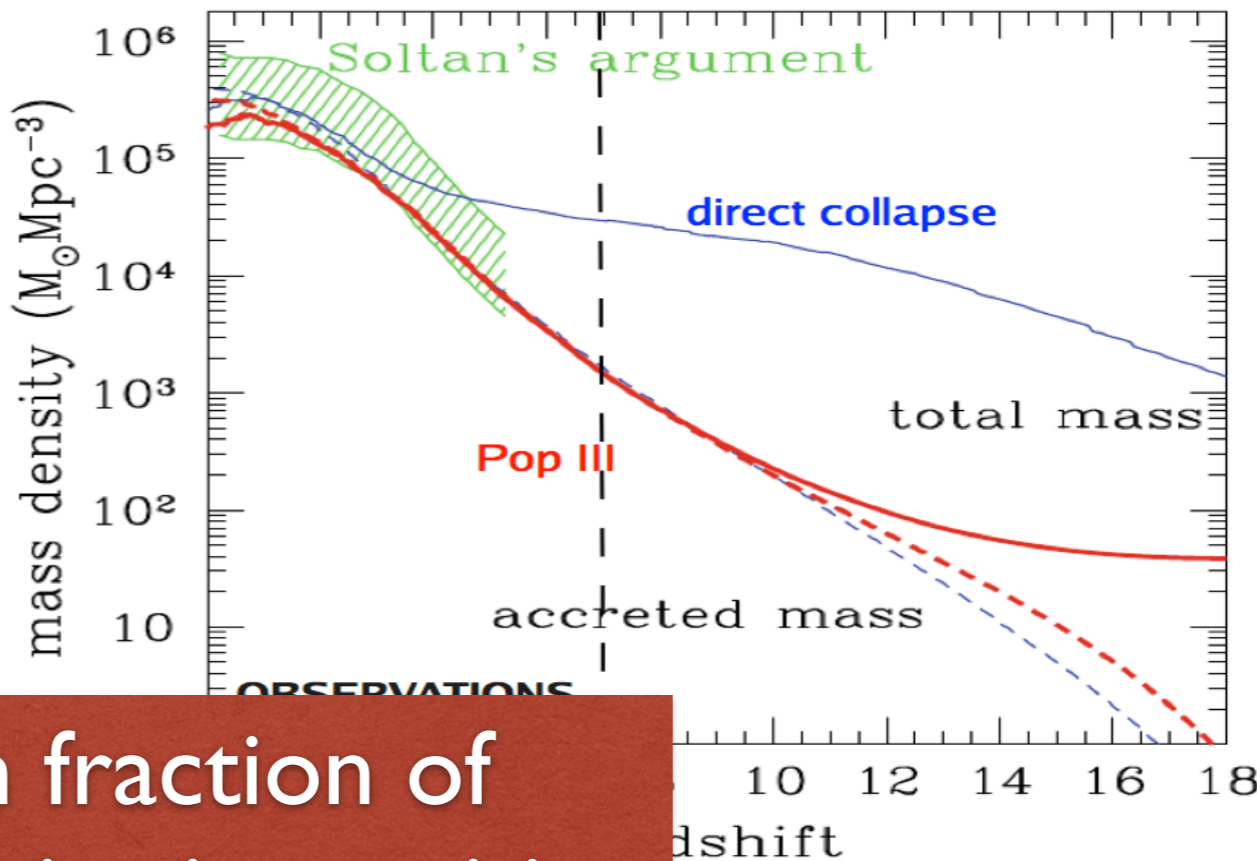
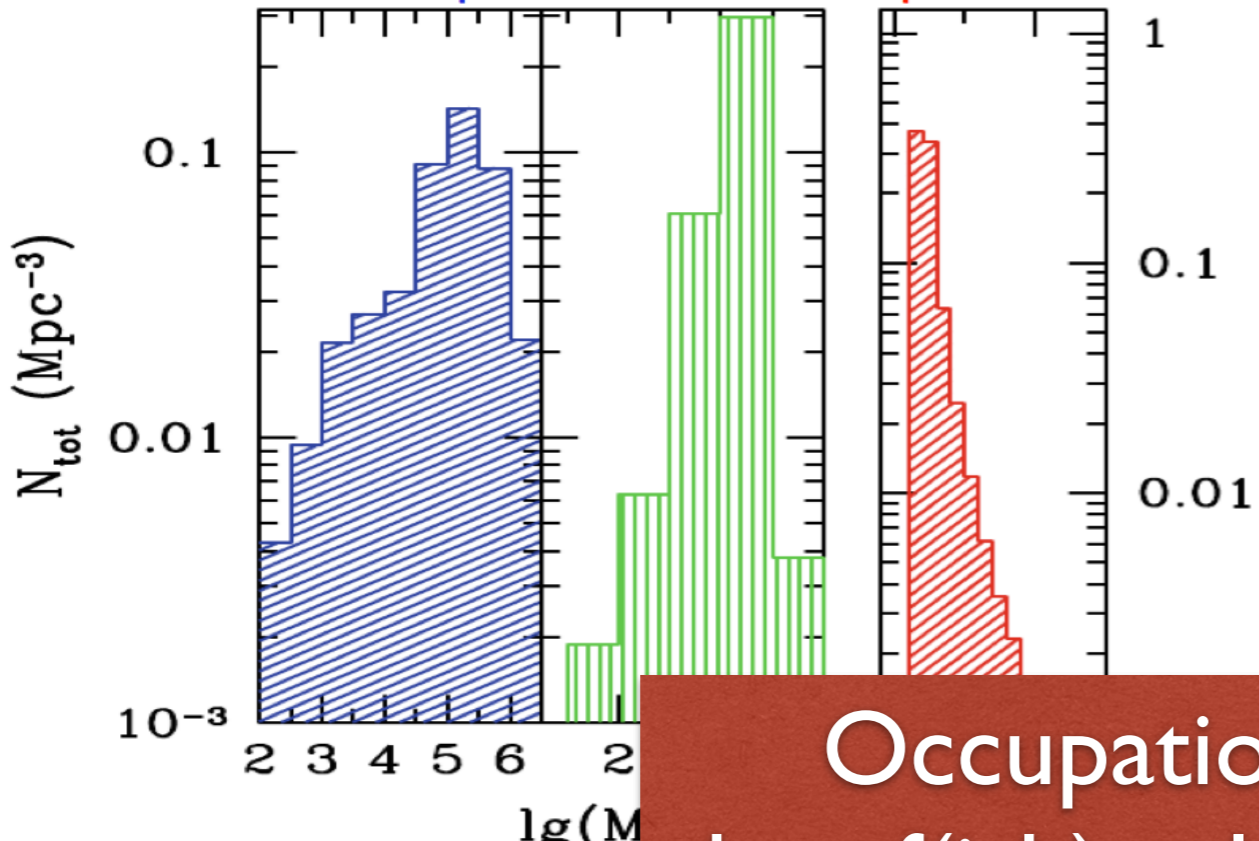
Bellovary et al. 2011

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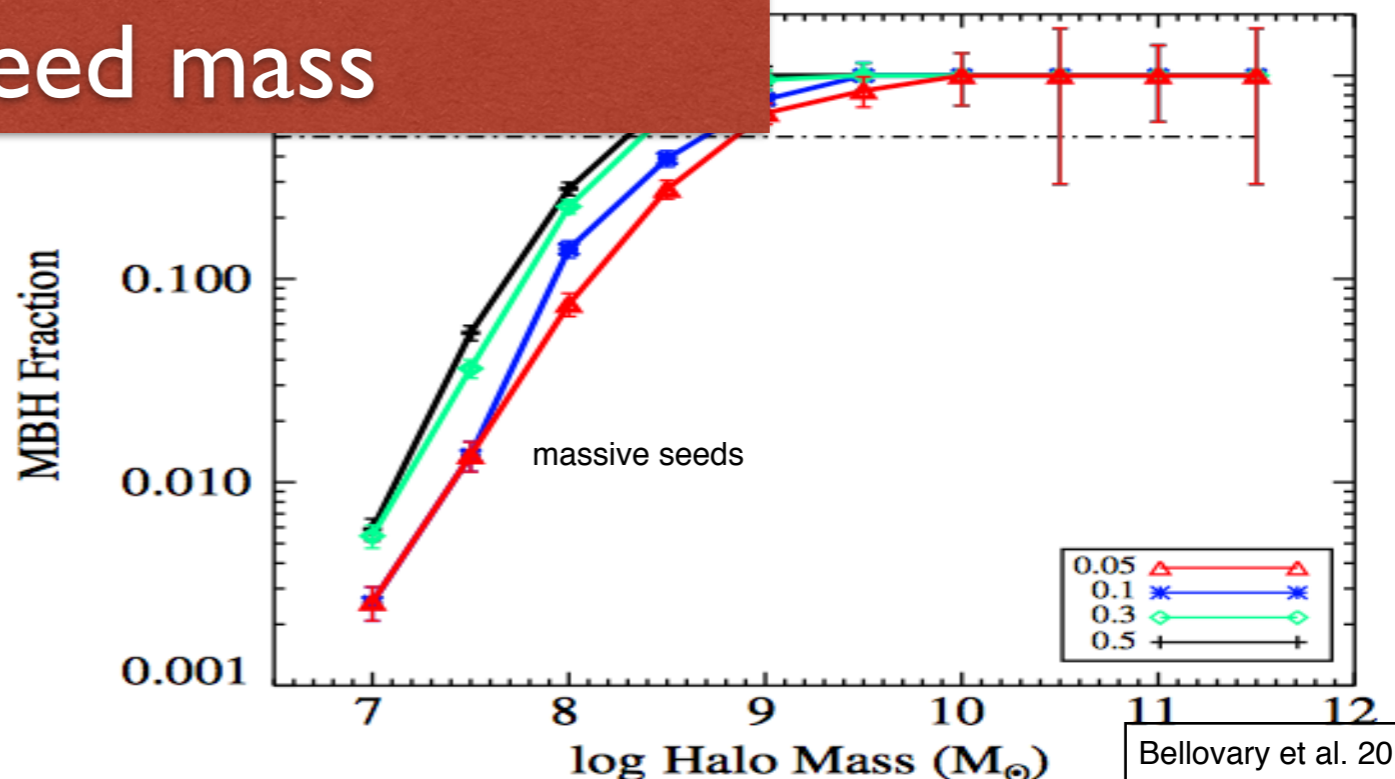
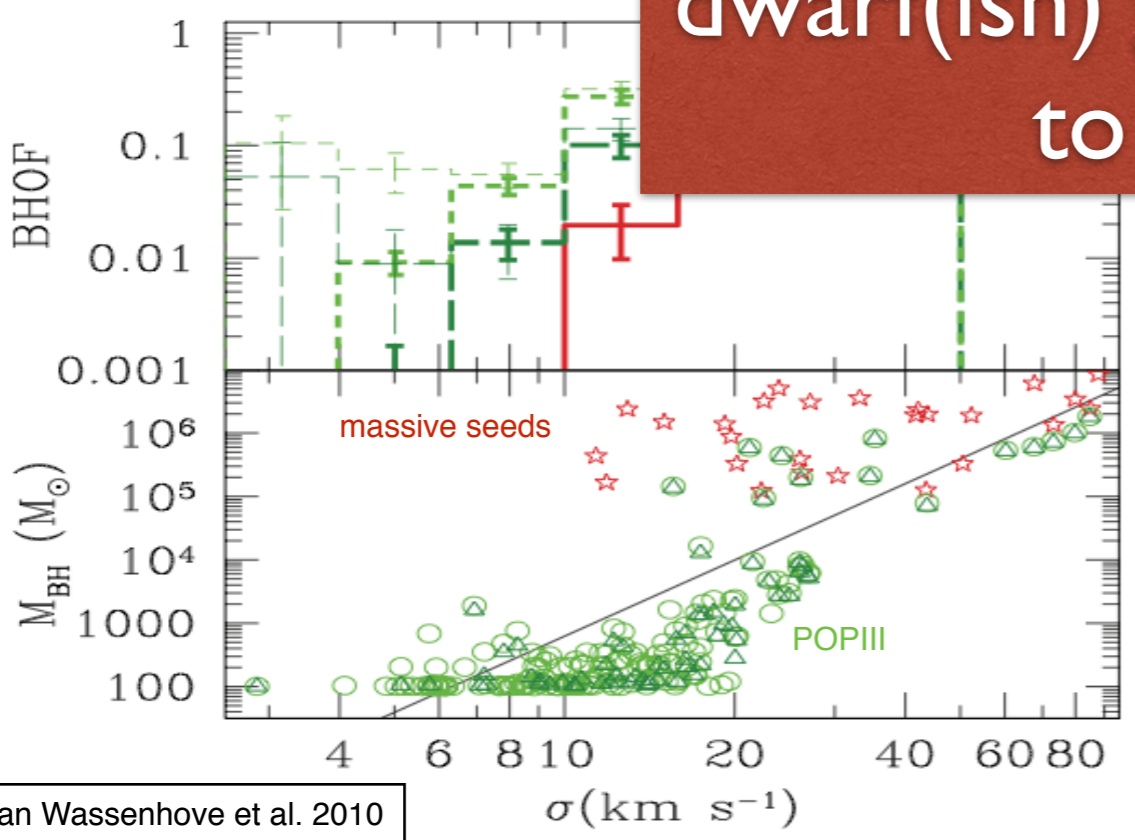
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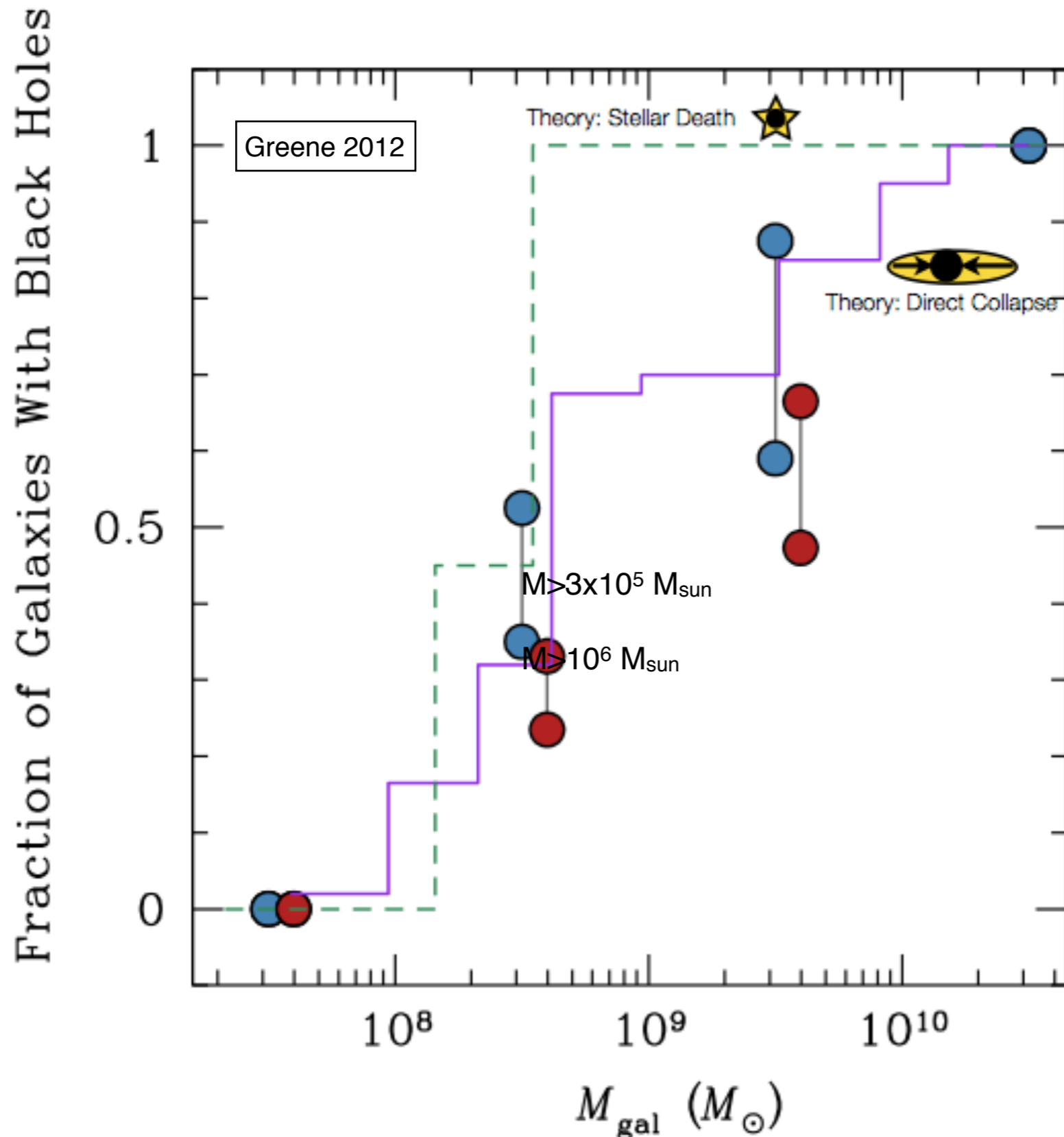
Occupation fraction of dwarf(ish) galaxies is sensitive to seed mass



van Wassenhove et al. 2010

Bellovary et al. 2011

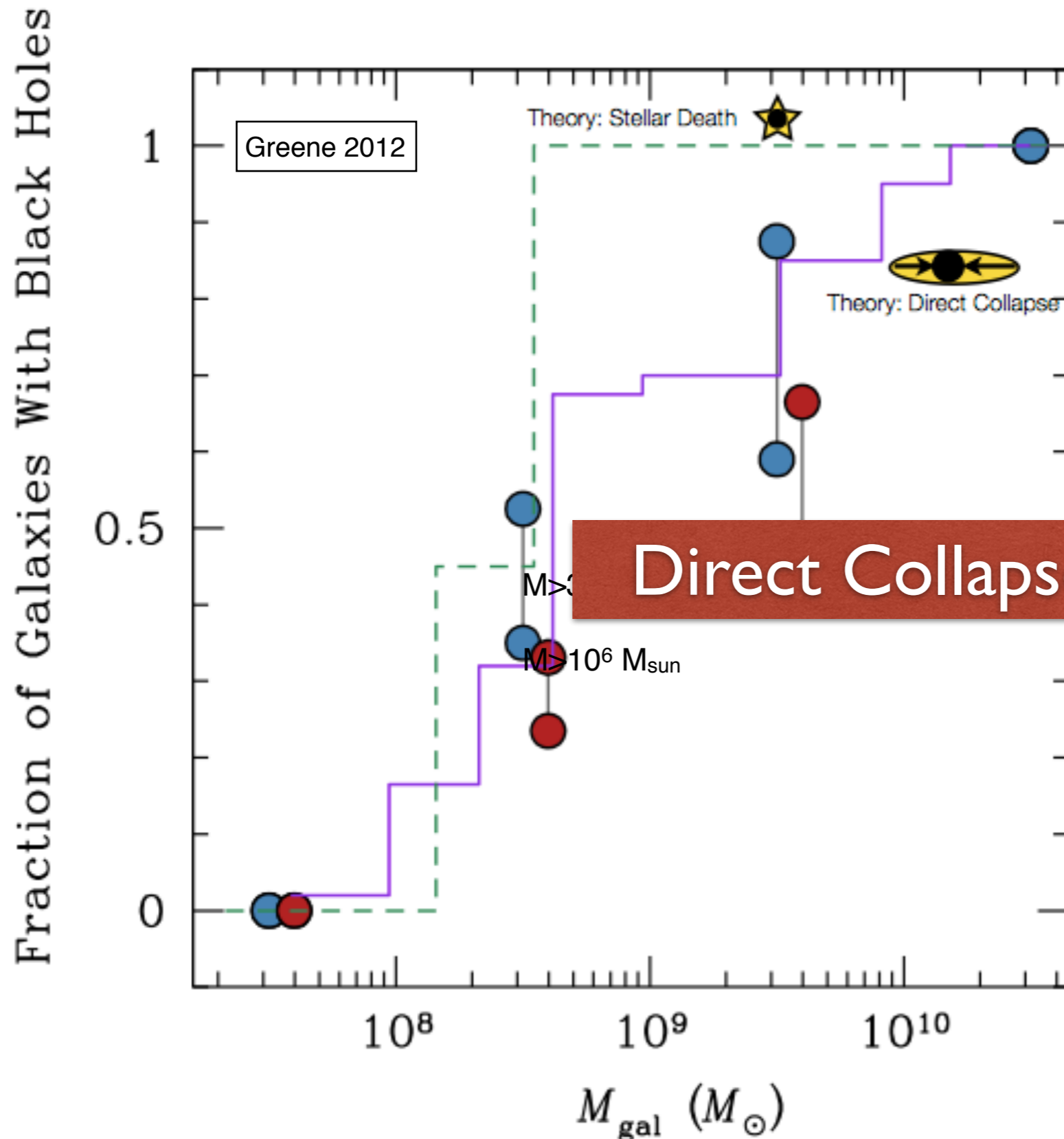
# Low Mass Seeds vs High Mass Seeds (observations)



POPIII seeds (green) and massive seeds (purple) from Volonteri et al. (2008).

Estimates from Desroches et al. (2009, blue) and Gallo et al. (2010, red).

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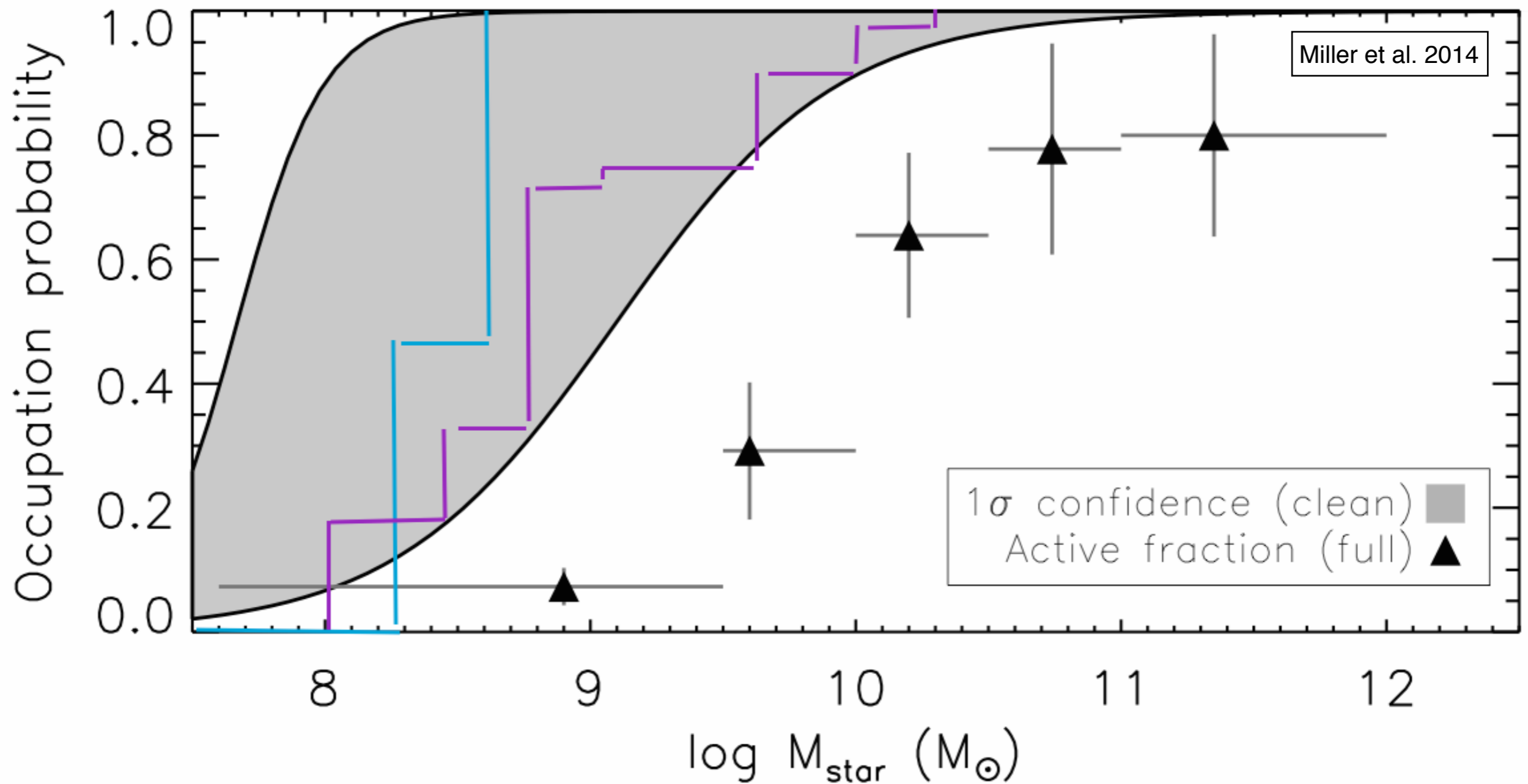


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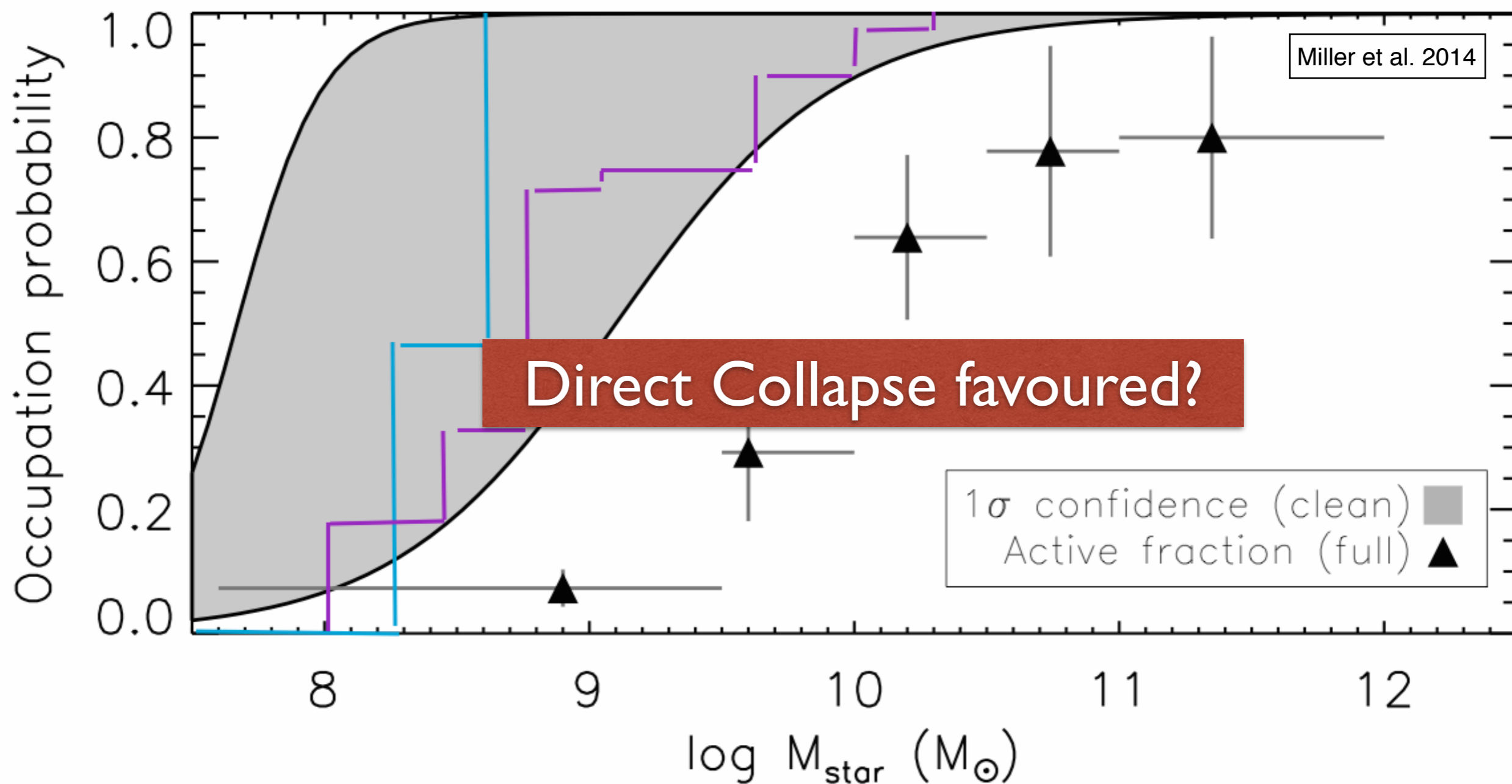
Direct Collapse favoured

# Low Mass Seeds vs High Mass Seeds (observations)





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# SMBHs at “Low” Redshifts by Gas Collapse

- Simulation of two well-formed hi-z MW-like galaxies merger Mayer et al. 2010
- Self-gravitating turbulent nuclear disk ( $r \sim 40$  pc,  $M \sim 2 \times 10^9 M_{\odot}$ ) forms
- Gravo-turbulent instabilities drive the gas toward the centre at very high rate,  $10^4 M_{\odot}/\text{yr}$
- A  $\sim 1$  pc size spherical structure grows by 13% of the disk mass ( $> 10^8 M_{\odot}$ ) in 0.1 Myr.  $T \sim 10^7 \text{K}$  prevents fragmentation.
- The core collapse into a BH + envelope = Quasistar

- The key point:

$$\dot{M} \approx M_J/t_{ff} = \frac{\pi^2}{8G} c_s^3 \simeq 5 \times 10^3 \left( \frac{T}{10^7 \text{K}} \right)^{3/2} M_{\odot} \text{yr}^{-1}$$

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- Would  $T$  be so high with a realistic thermodynamics? (the simulation uses a polytropic EoS) Ferrara, FH & Salvaterra 2013
- On which timescale is turbulence (the source of energy) dissipated?

# Conclusions

## Light Seeds

- POPIII stars form copiously in the high-z Universe ( $z \sim 20$ ).
- $M_{\text{star}} \sim 45 M_{\odot} \rightarrow M_{\text{BH}} < 45 M_{\odot}$
- Accretion rate very low because of UV feedback (but simulations need to be truly cosmological)

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## Heavy Seeds

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- If metal polluted, (i.e., lower  $z$ ) runaway collisions (need  $Z$  fine tuning?)

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## Observations

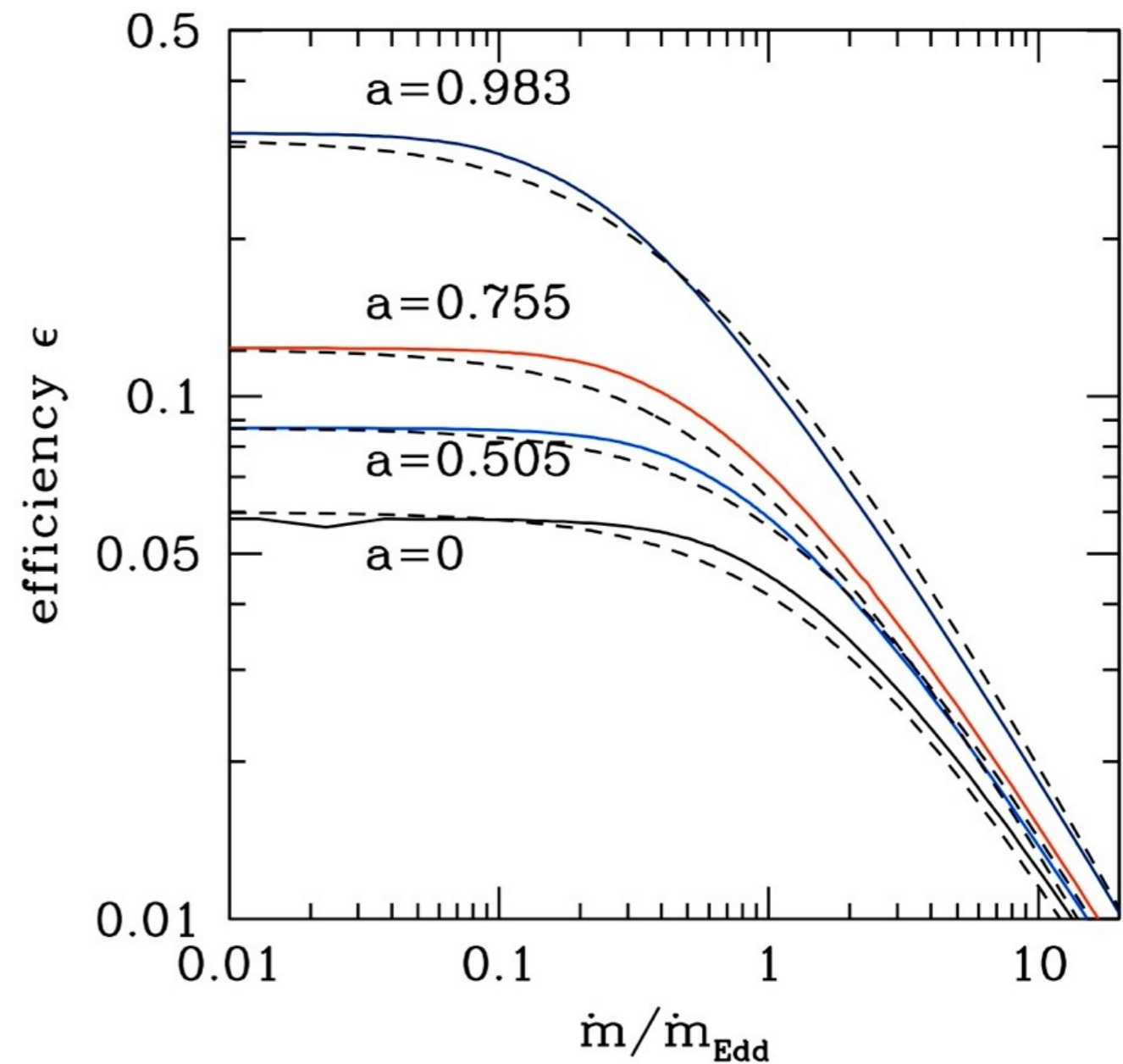
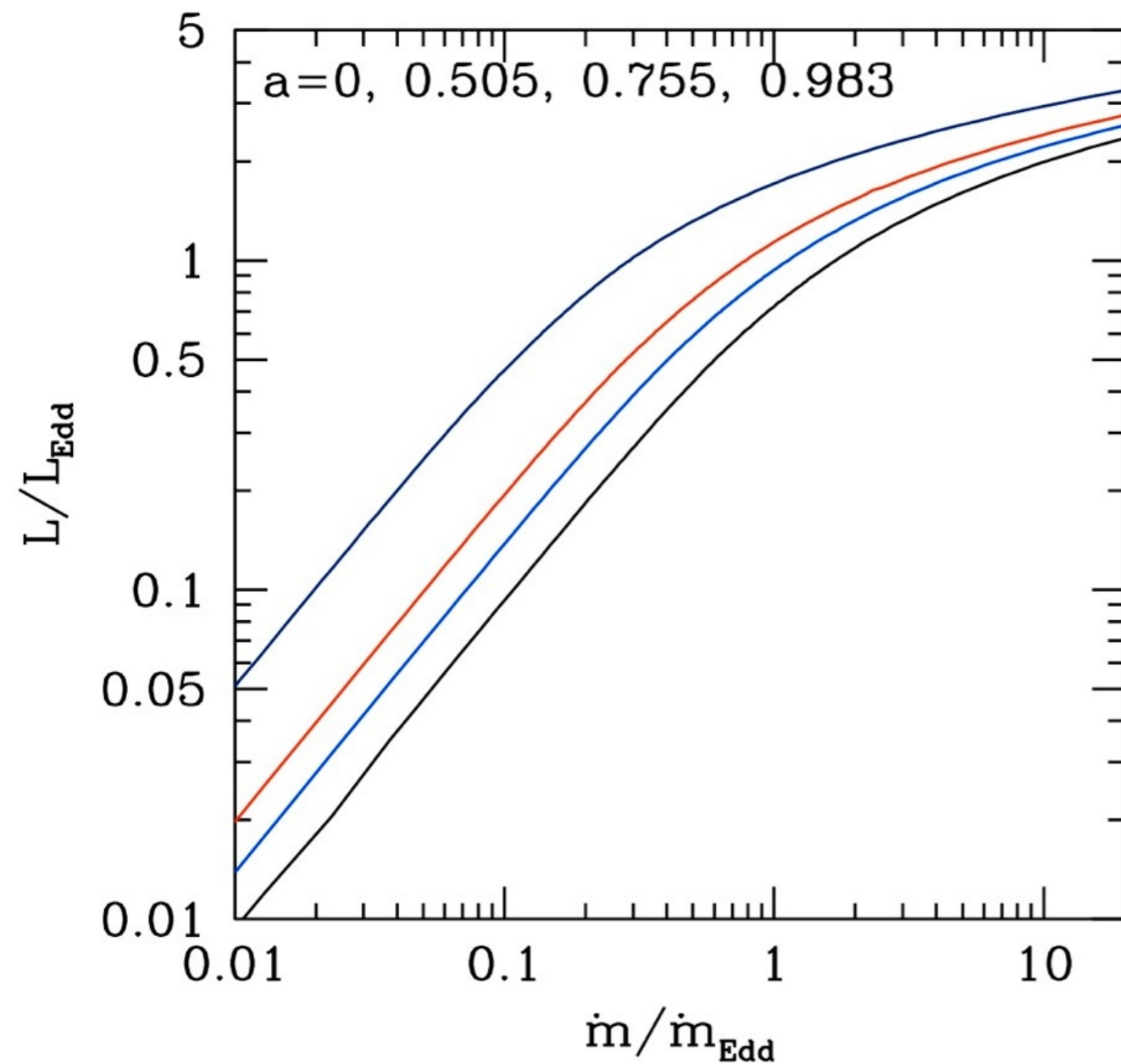
- Imprint of seeds on occupation fraction.
- Observations tell nothing so far.
- Future GW observations (eLISA) of the MBH merger rate at very-high redshifts.

# SuperEddington accretion (slim disks)

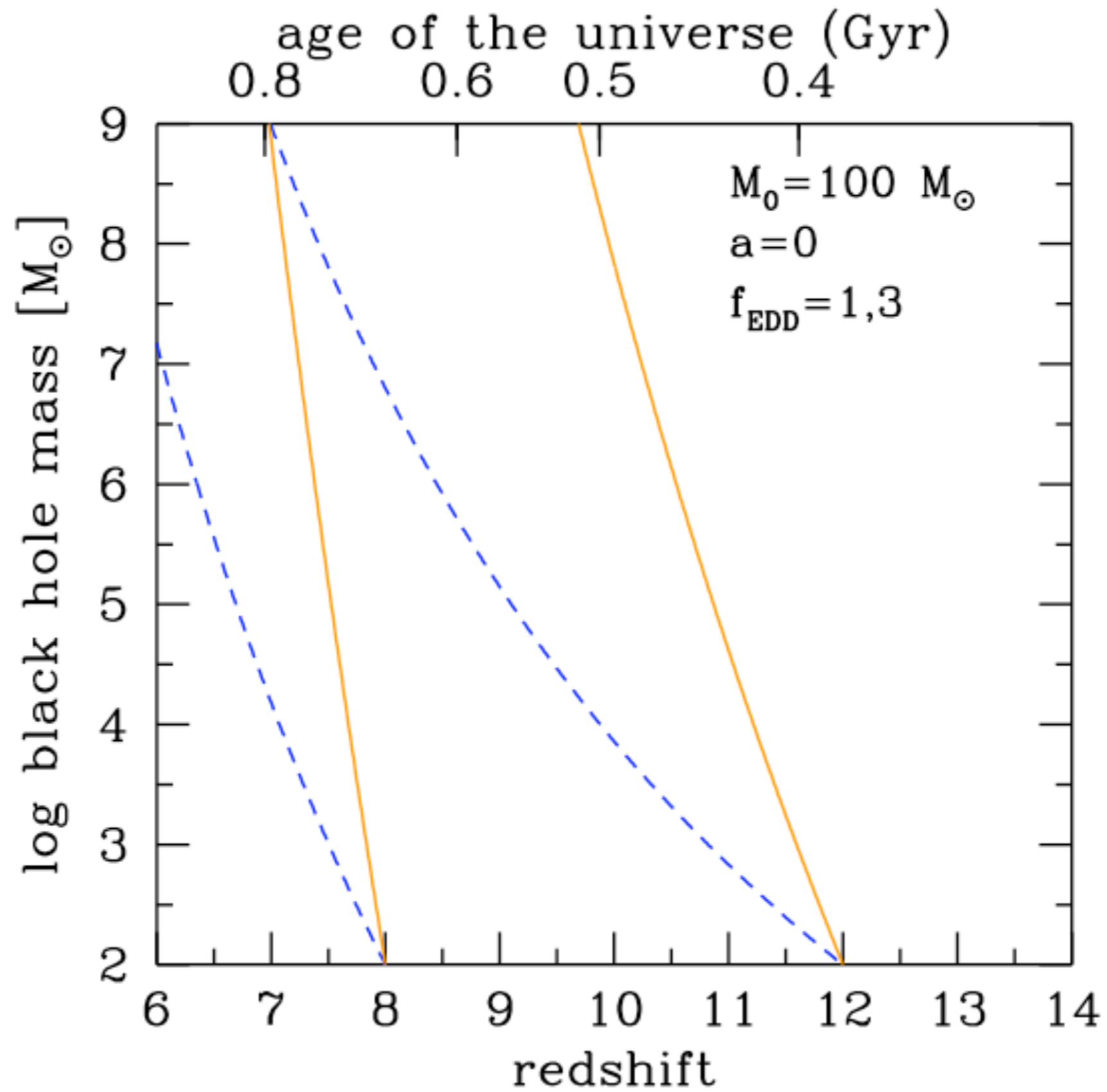
Madau, FH & Dotti 2014

$$M(t) = M_0 e^{f_{\text{Edd}} \frac{t}{\tau} \frac{1-\epsilon}{\epsilon}}$$

$\tau \cong 0.44$  Gyrs

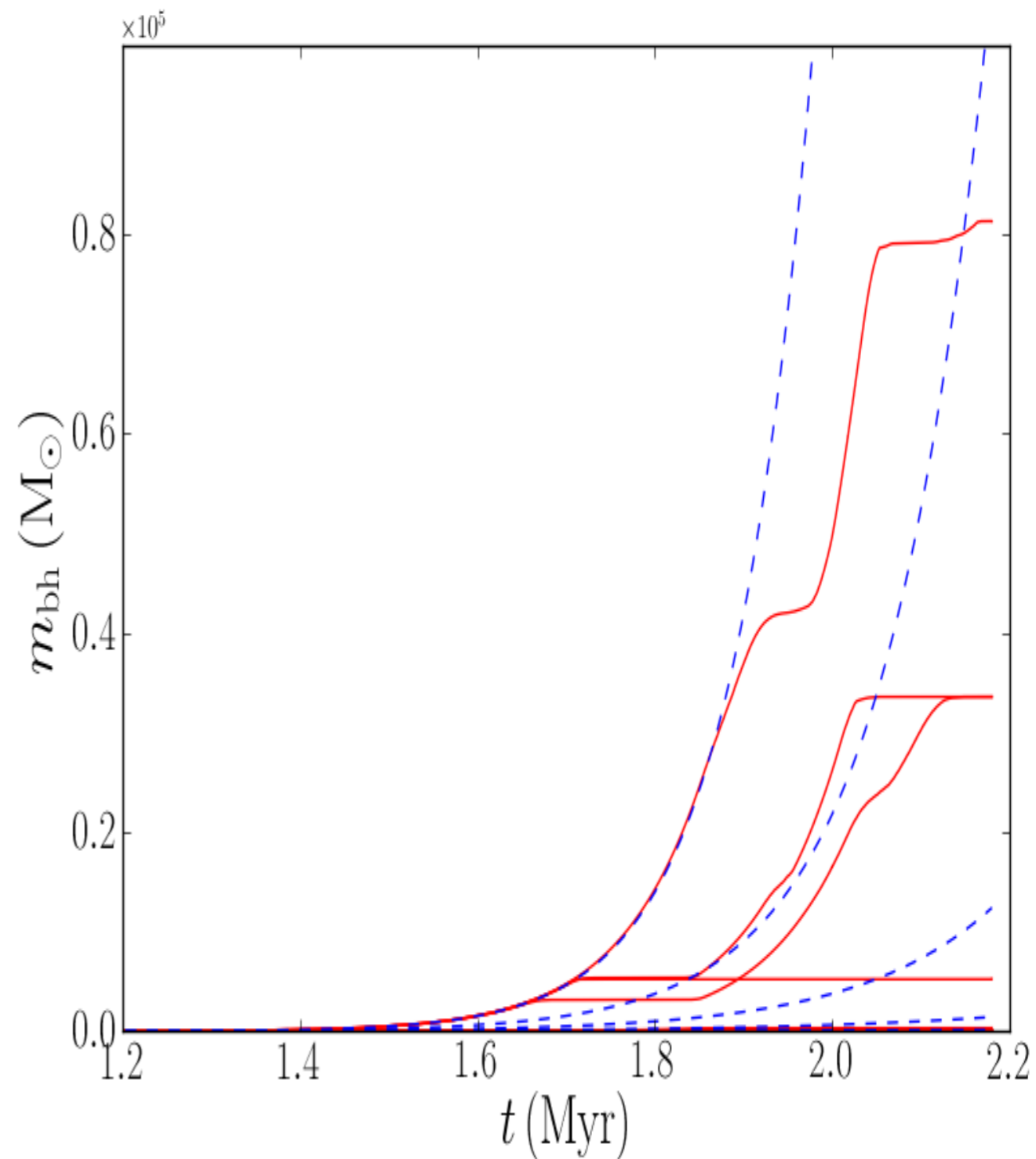
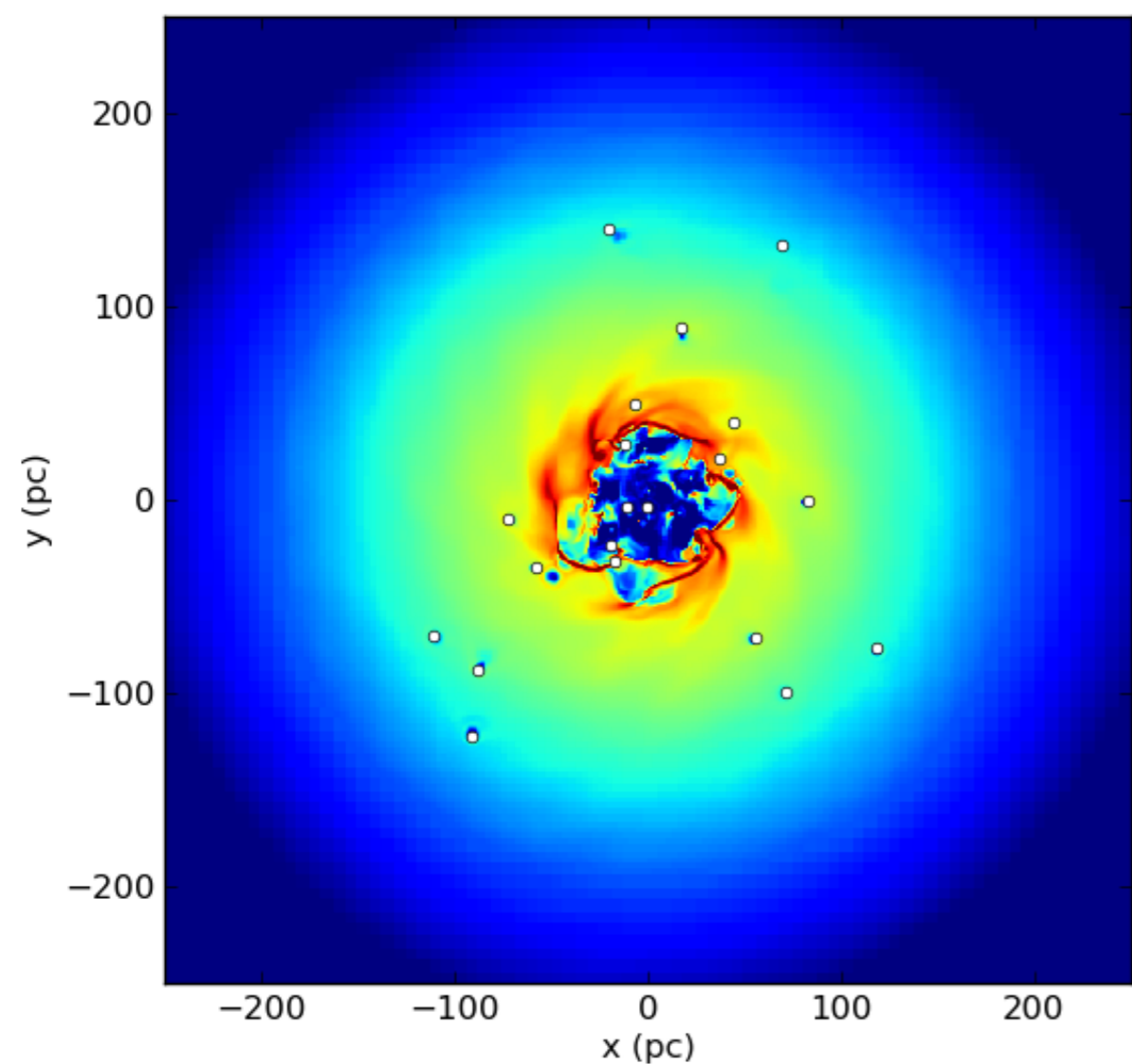


Madau, FH &amp; Dotti 2014





# SuperEddington accretion at not-so high $z$



Lupi et al., in prep