

**“Hydrogen is a light, odorless gas, which,
given enough time, turns into people.”**

Edward Robert Harrison



Measuring the Jeans Scale of the Intergalactic Medium with Close Quasar Pairs

Alberto Rorai, Joseph Hennawi, Martin White

OUTLINE

- Structure formation, reionization and thermal history of the IGM
- Probing the IGM with the Lyman- α forest
- Statistics and degeneracies in the Ly- α forest: a parameter study
- A new statistic: phase differences in quasar pairs
- Properties and constraining power of phase statistic
- Preliminary results at $z=2$ and $z=3$

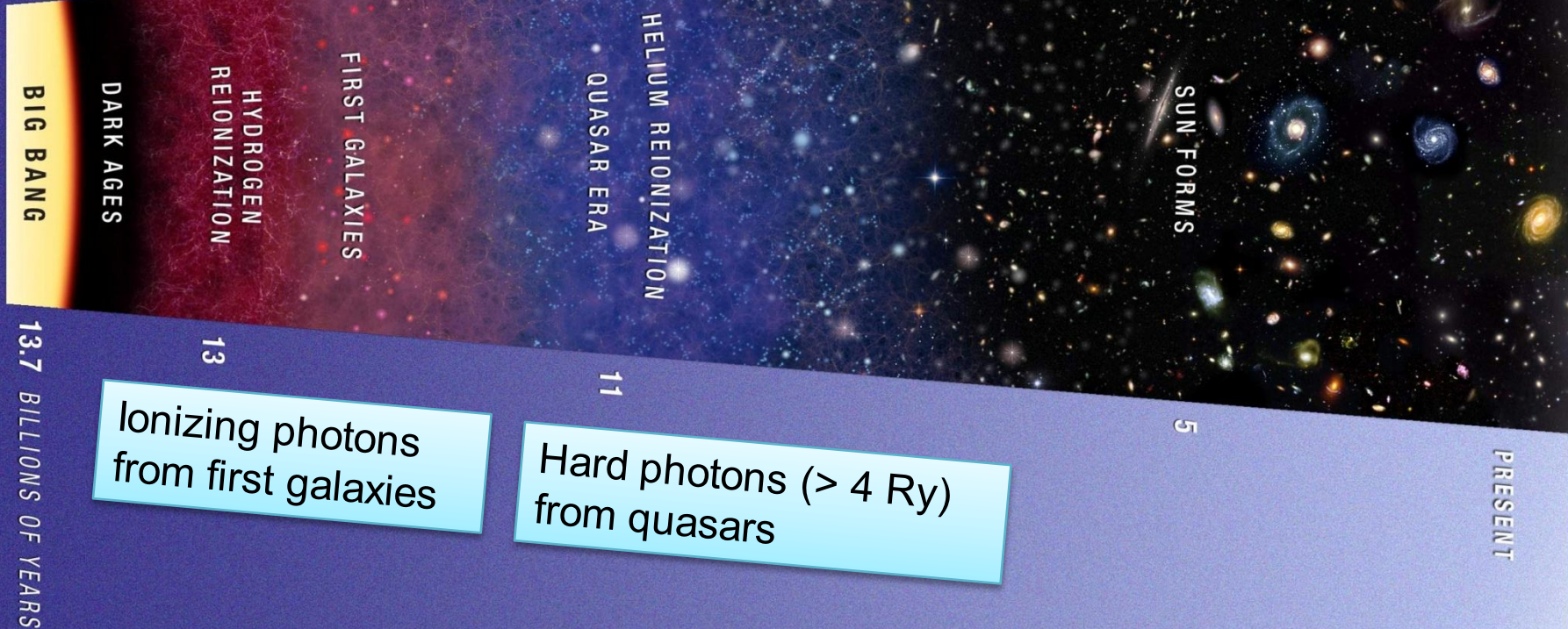
The Epoch(s) of Reionization

Cold, neutral IGM

Hot, ionized IGM

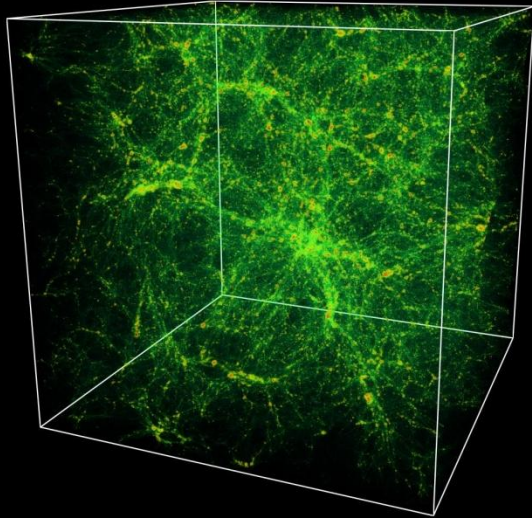
$Z \sim 30 - 6$: H and He singly ionized

$Z \sim 2-4$: He II ionized

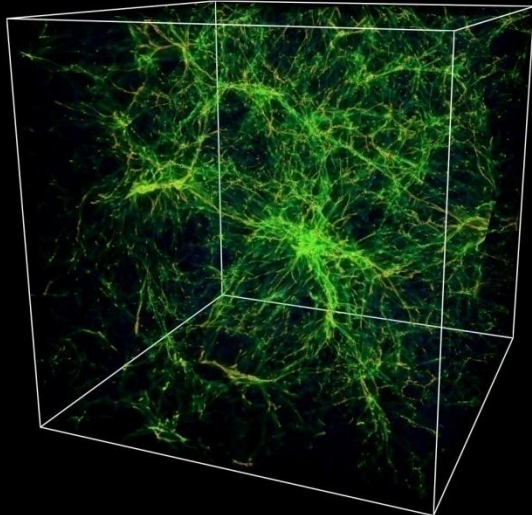


Density Distribution in the Universe

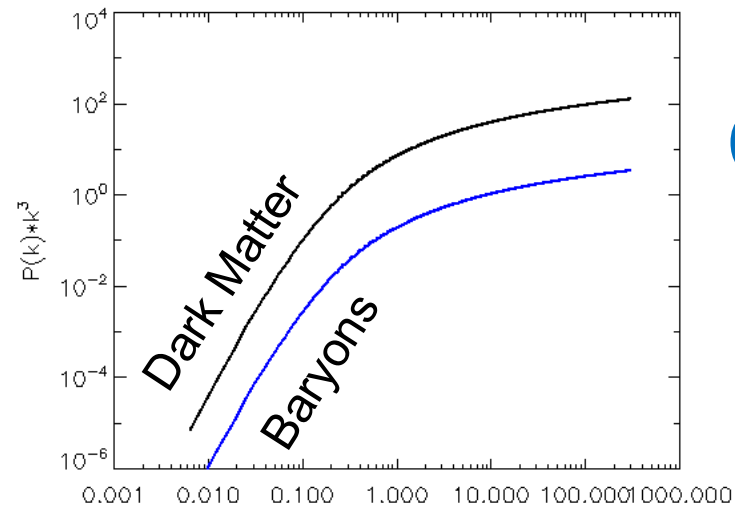
Dark Matter



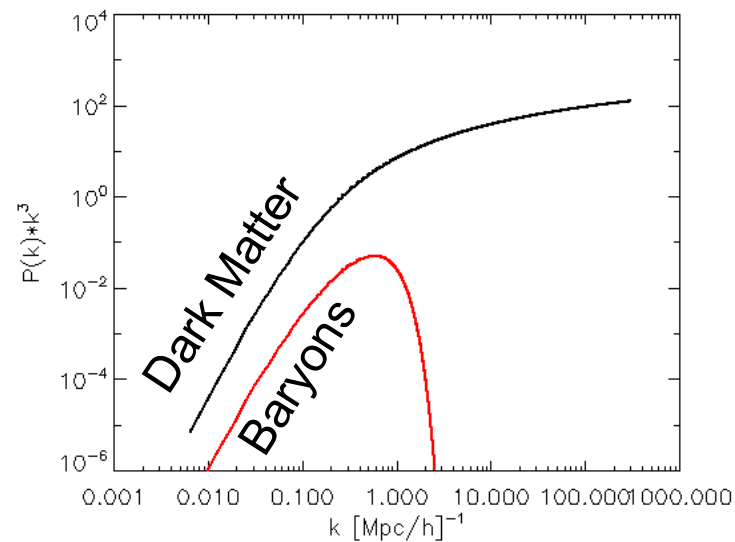
Baryons



Credit: R. Cen



Cold IGM
 $T=0 \text{ K}$



Hot IGM
 $T=10^4 \text{ K}$

The Jeans Scale

- Self-gravitating structure of size λ collapses if
free-fall time $(G\rho)^{-1/2} < \lambda c_s$ **sound-crossing time**
critical λ is called “*Jeans scale*”

$$\lambda_J = \frac{c_s}{1+z} \sqrt{\frac{\pi}{G\rho}} \approx 1.1 \text{ Mpc} \left(\frac{T}{10^4 \text{K}} \right)^{1/2} \left(\frac{1+z}{1+3} \right)^{-1/2}$$

- Small-scale-structure growth and Galaxy formation suppressed below the **Filtering mass**:

$$M_F = \frac{4}{3} \pi \bar{\rho} \lambda_J^3$$

- Jeans filtering scale depends on the *whole thermal history* (Gnedin & Hui, 1998): sound-crossing time λ_J/c_s is comparable to the Hubble time

$$\lambda_J^2(t) = \int_0^t f(T(t')) dt'$$

Factors Governing Reionization and Thermal History

- Population and spectrum of ionizing sources
- Escape fraction
- Clumping factor (also related to Jeans scale!)
- Adiabatic cooling due to expansion
- Perturbations growth

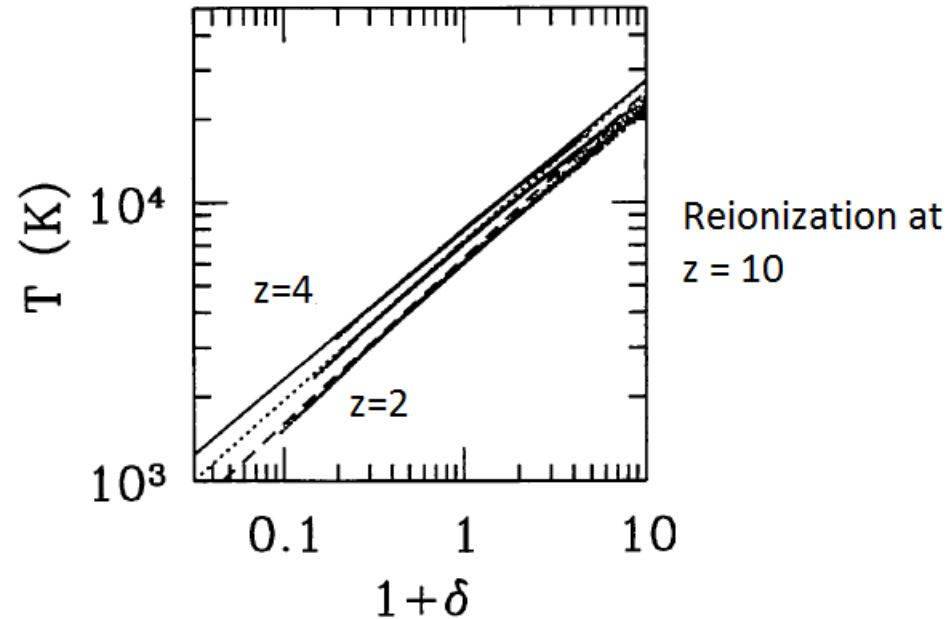
Under the assumption of ionization equilibrium:

$$T(\rho) = T_0 \left(\frac{\rho}{\rho_0} \right)^{\gamma-1}$$

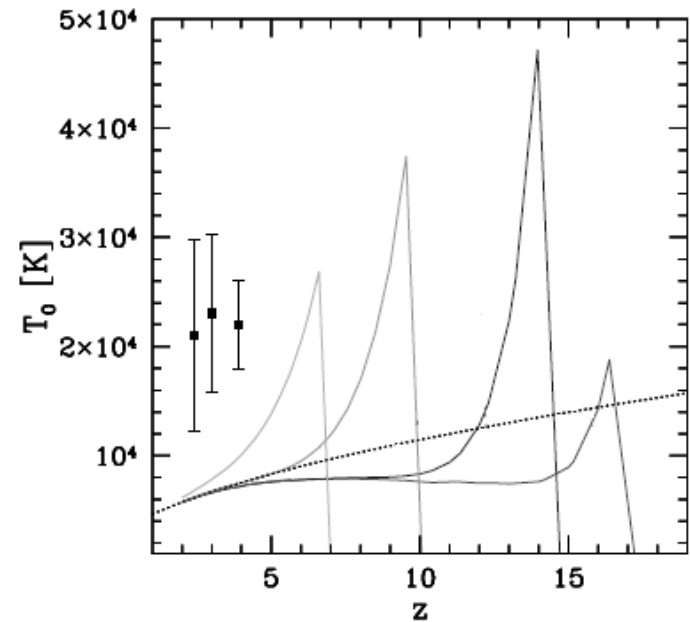
“Equation of state” of the IGM (Hui and Gnedin, 1997)

The Thermal State of The IGM

$$T(\rho) = T_0 \left(\frac{\rho}{\rho_0} \right)^{\gamma-1}$$



Hui and Haiman 2003 , measurements from Zaldarriaga et al. 2001



The Instantaneous temperature of the IGM provides little information about reionization heating occurred in the distant past

Summary: Galaxies and IGM

Galaxies

Intergalactic Medium

- **Reionize** Hydrogen and Helium in the IGM

- Energetic photons deposit energy: higher IGM **temperature**

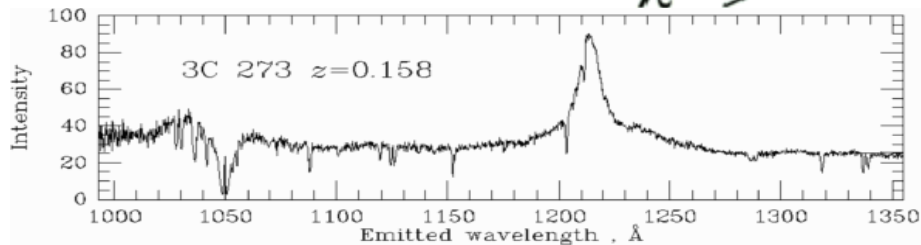
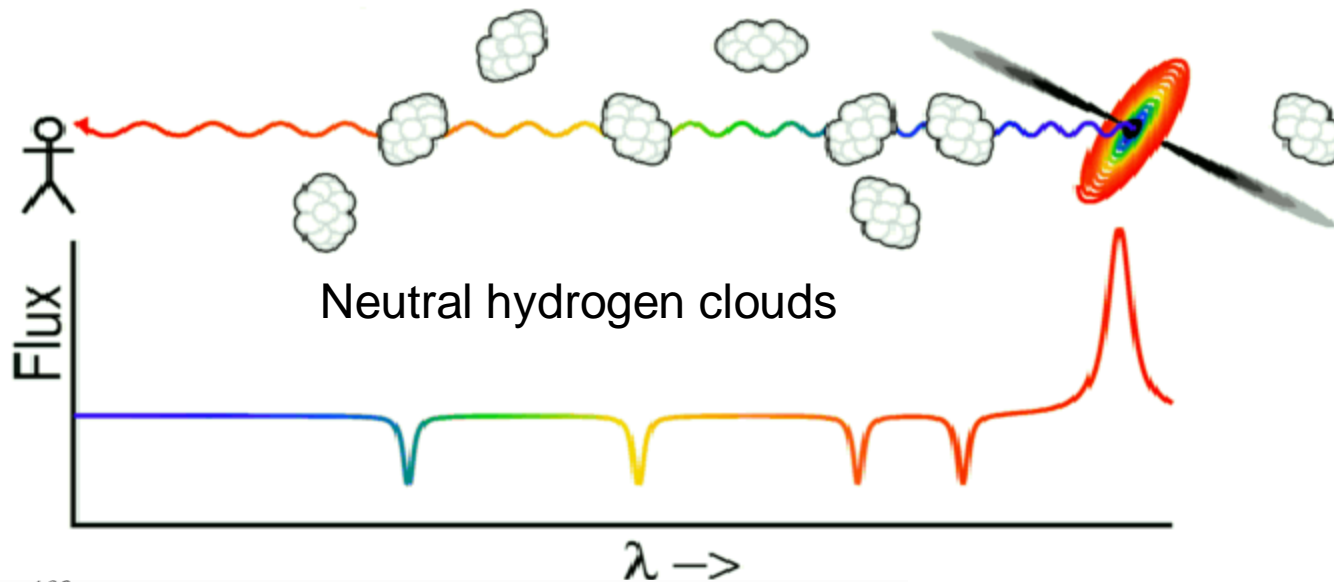
- Pressure increase changes the **Jeans scale** of the IGM

- Fuel for **galaxy formation**

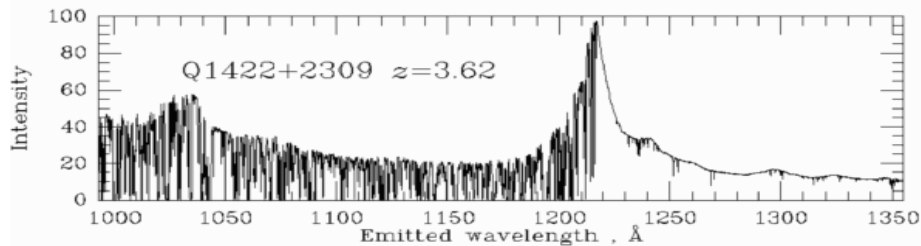
- Pressure determines the **minimum-mass** galaxy that can collapse

- Clumps absorb ionizing photons and **slow down reionization**

Probing the IGM: the Lyman α Forest



Low redshift:
- Lower density
- Higher ionized fraction



High redshift:
- Higher density
- Lower ionized fraction

Physics of the Lyman α Forest

Photoionization equilibrium

$$n_{HI} \approx n_H^2 \frac{\alpha(T)}{\Gamma}$$



Equation of state

$$T \propto (1 + \delta)^{\gamma-1}$$

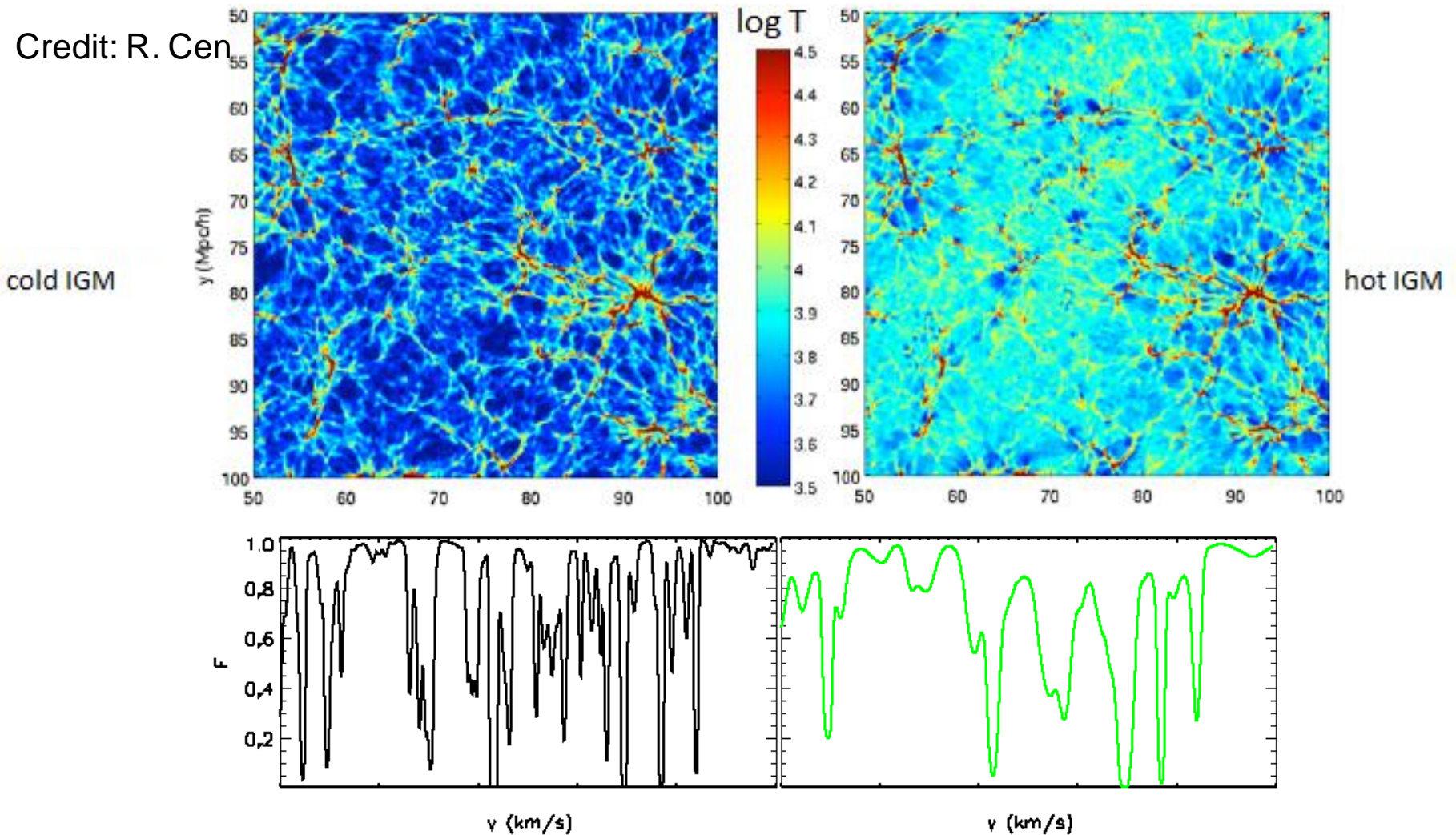


Fluctuating Gunn-Peterson approximation

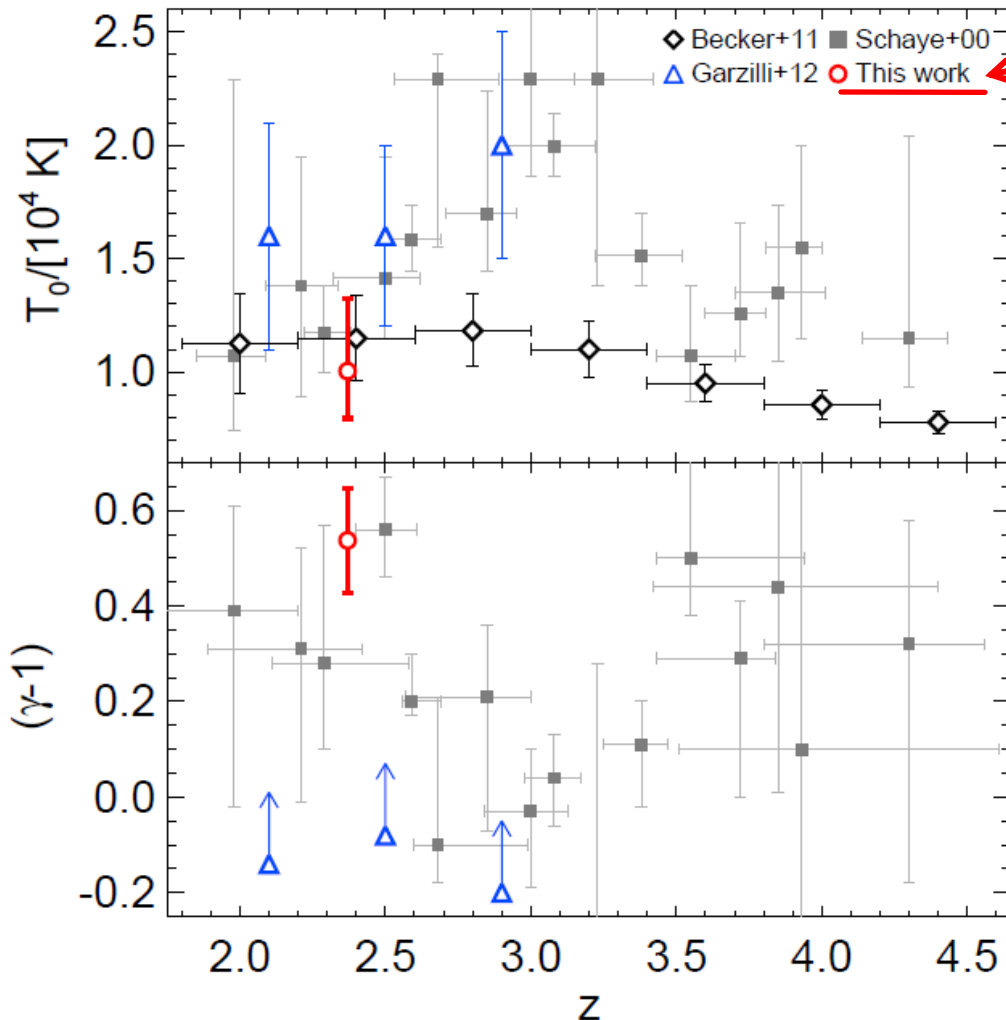
$$\tau \propto (1 + \delta)^{2-0.7(\gamma-1)}$$

- Physics relatively easy to understand
- Not biased by light-to-mass assumptions
- Probes smaller scales than CMB

Relation Forest – Thermal State



Constraints from the Lyman α Forest on the Thermal State



From Bolton et al., 2013

Line broadening must be connected to temperature \rightarrow

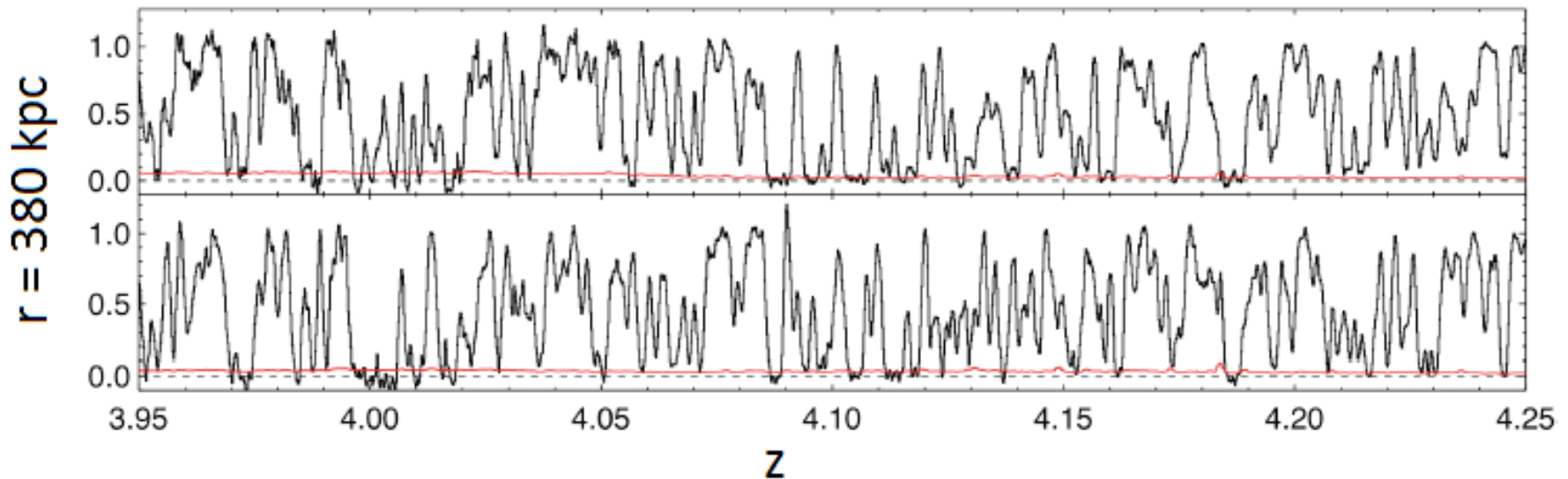
Several statistical tool used to constrain temperature:

- Power spectrum
- Flux Probability distribution
- Wavelet analysis
- Line fitting
- Curvature

The Jeans scale also contributes to line width, and must be assumed or modeled in these measurements.

Probing the IGM with Quasar Pairs

There are ~ 300 known pairs at $1.6 < z < 4.3$ at $r < 1$ Mpc
(Hennawi et al. 2004, 2006, 2009)



- Close separation (380 comoving kpc)
- High degree of coherence



We are observing the
Jeans scale at $z=4$

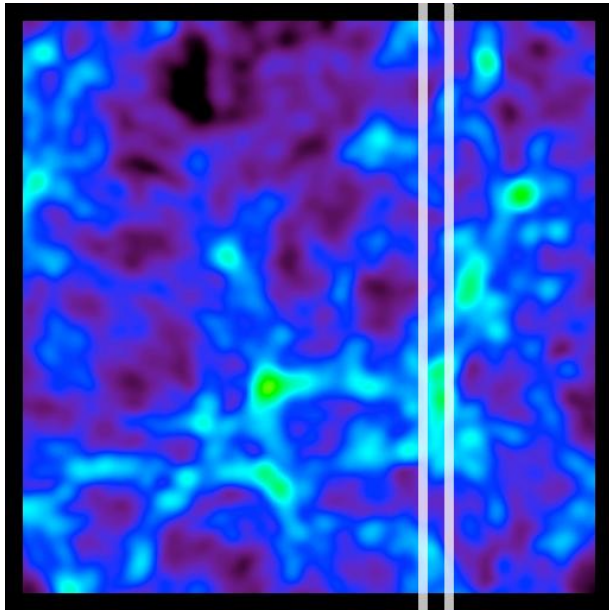
Key Questions

cold and **steady** IGM: Lyman- α forest traces density exactly.

Real IGM: **thermal broadening** and **redshift-space distortion** introduce degeneracies among density distribution and thermal properties of the IGM.

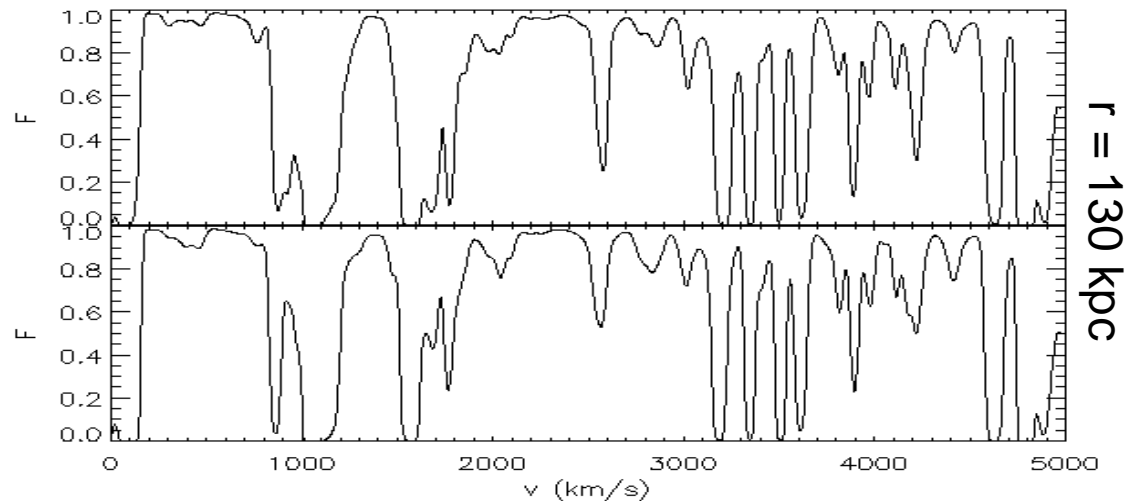
- To which extent does pressure affect the statistics of the Lyman- α forest lines?
- Which is the optimal statistic for measuring the Jeans scale with pairs?
- What additional constraints do quasar pairs put on the parameters governing the IGM thermal state?

Methodology: Simulations and Parameter Study



1. Dark matter only simulation, snapshot at $z=3$ ($Box=50 \text{ Mpc}/h$, $n=1500^3$)
2. Mimic pressure support by a smoothing density with a kernel of radius λ_j .
3. Definition of the equation of state

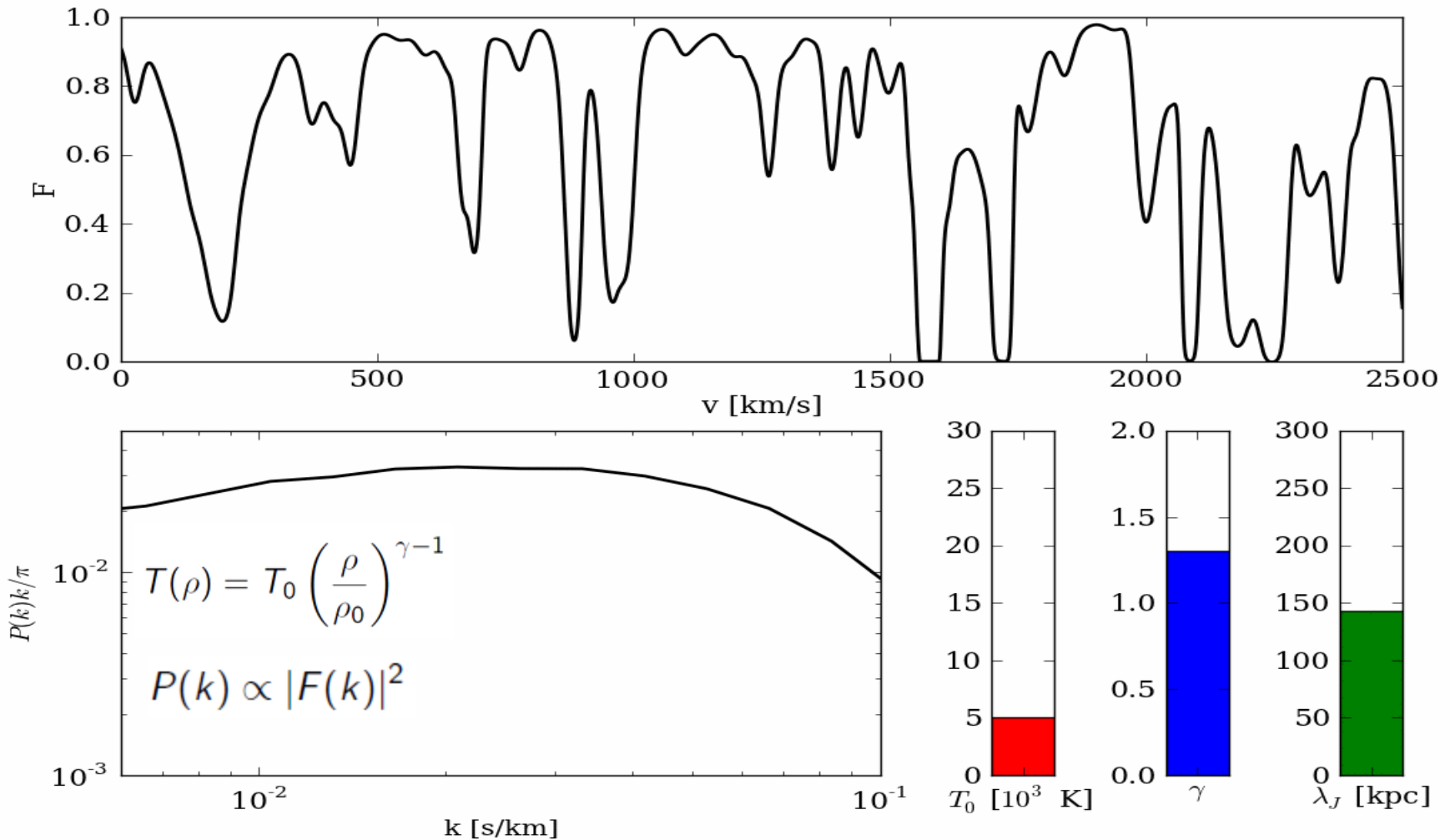
$$T(\rho) = T_0 \left(\frac{\rho}{\rho_0} \right)^{\gamma-1}$$



We calculate 500 different thermal models, defined by T_0 , γ and λ_j

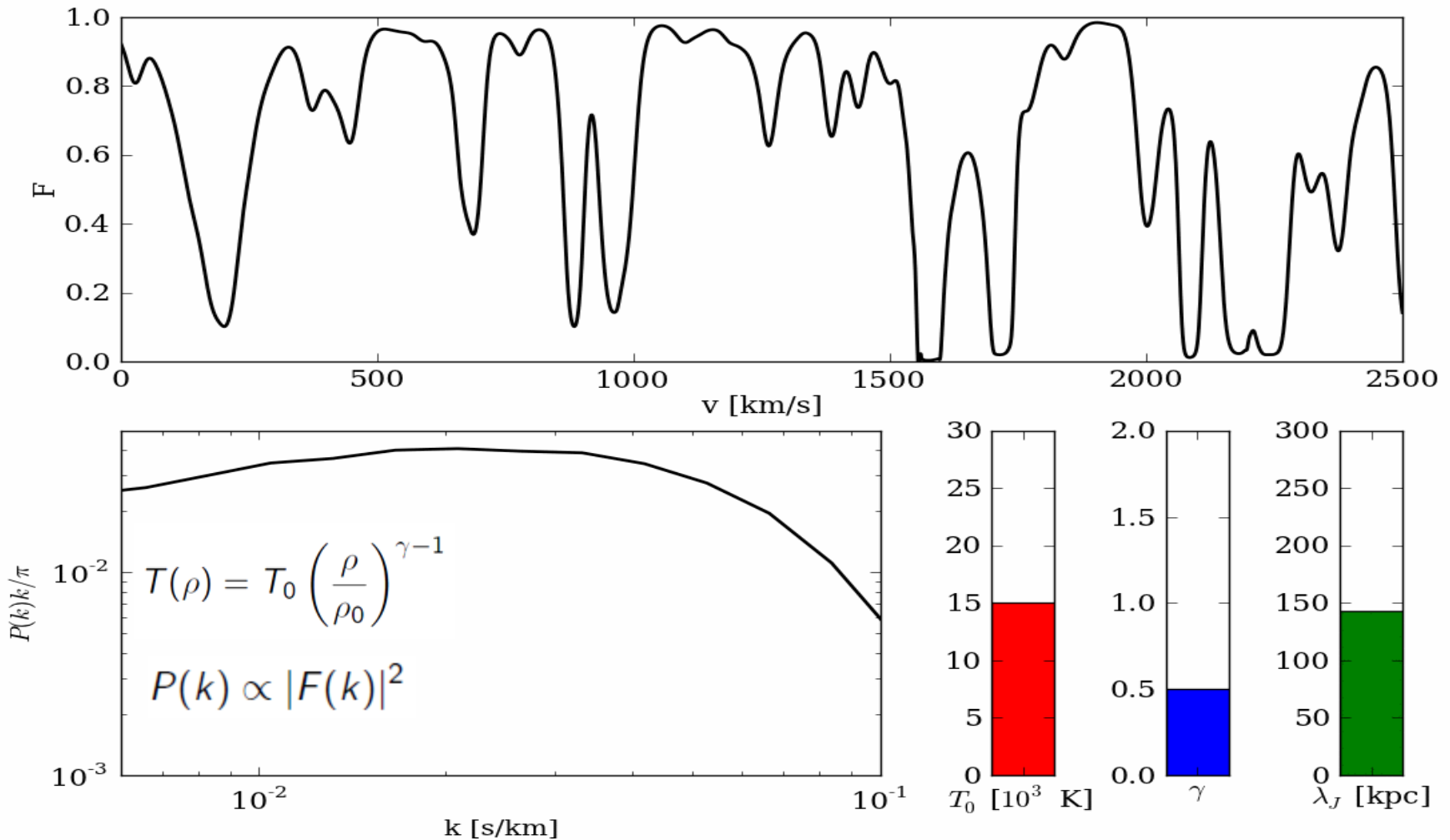
Sensitivity of the LOS Power Spectrum to 1D Smoothing

$z = 3$



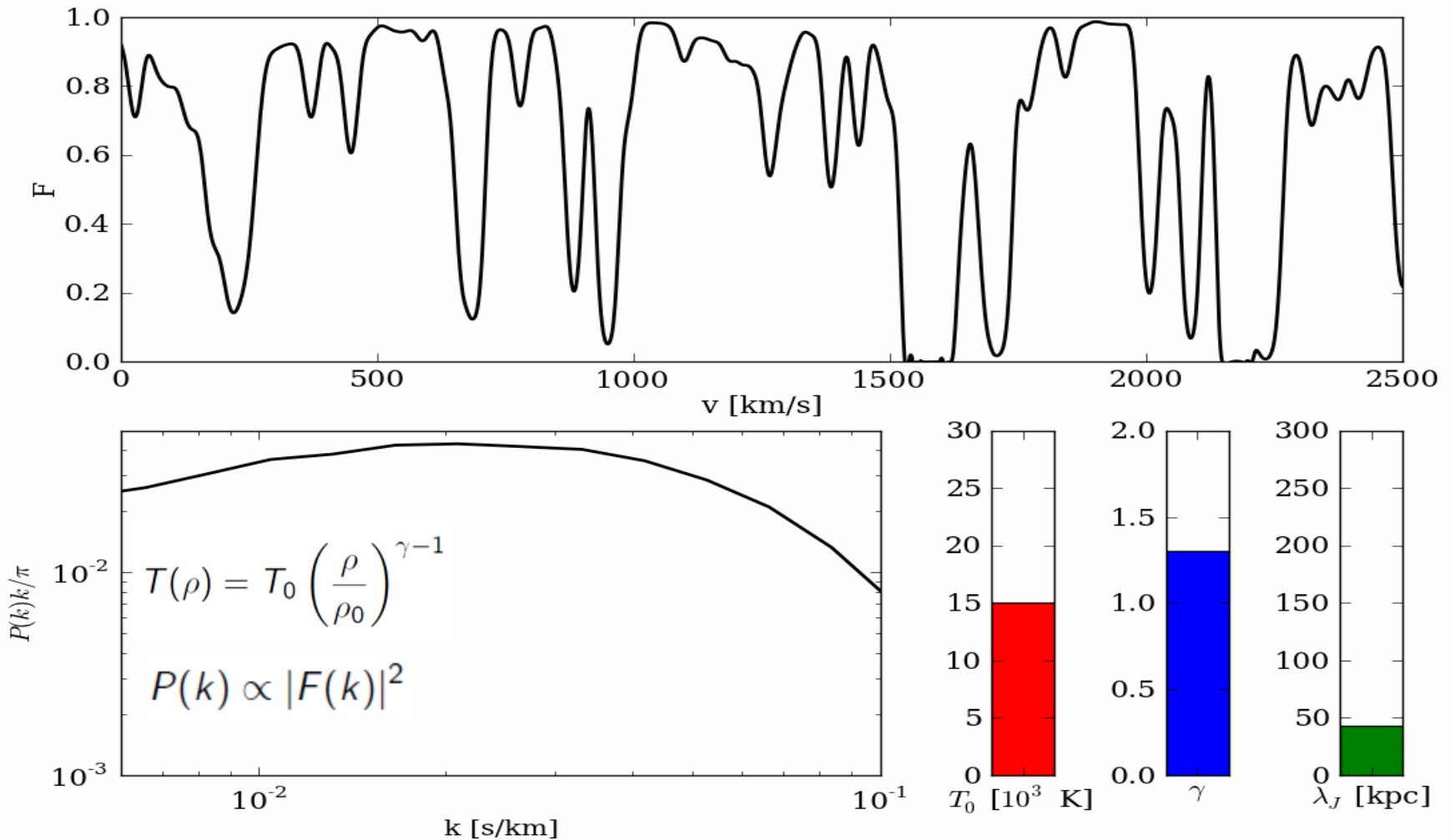
Sensitivity of the LOS Power Spectrum to 1D Smoothing

$z = 3$

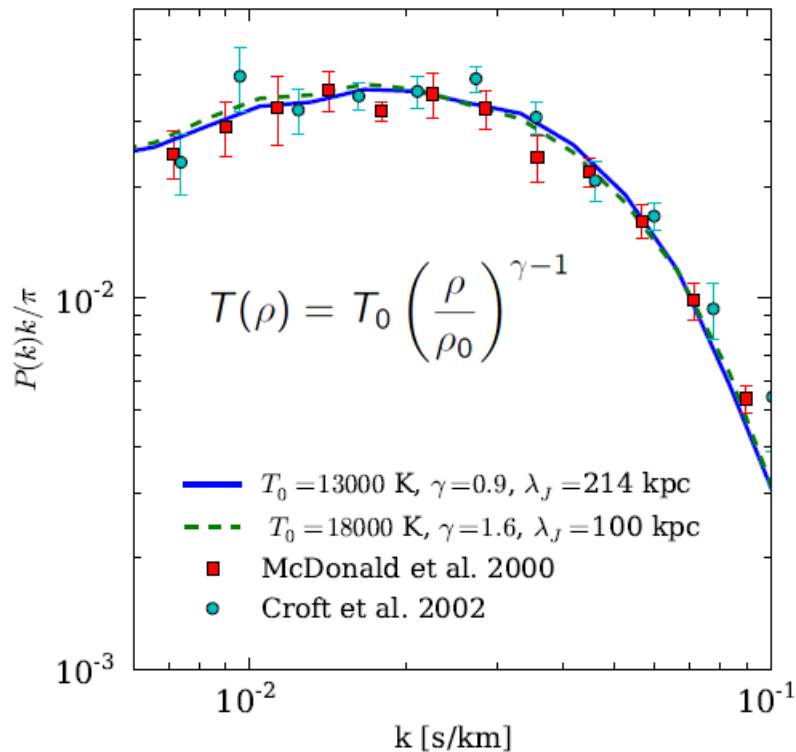


Sensitivity of the LOS Power Spectrum to 3D Smoothing

$z = 3$

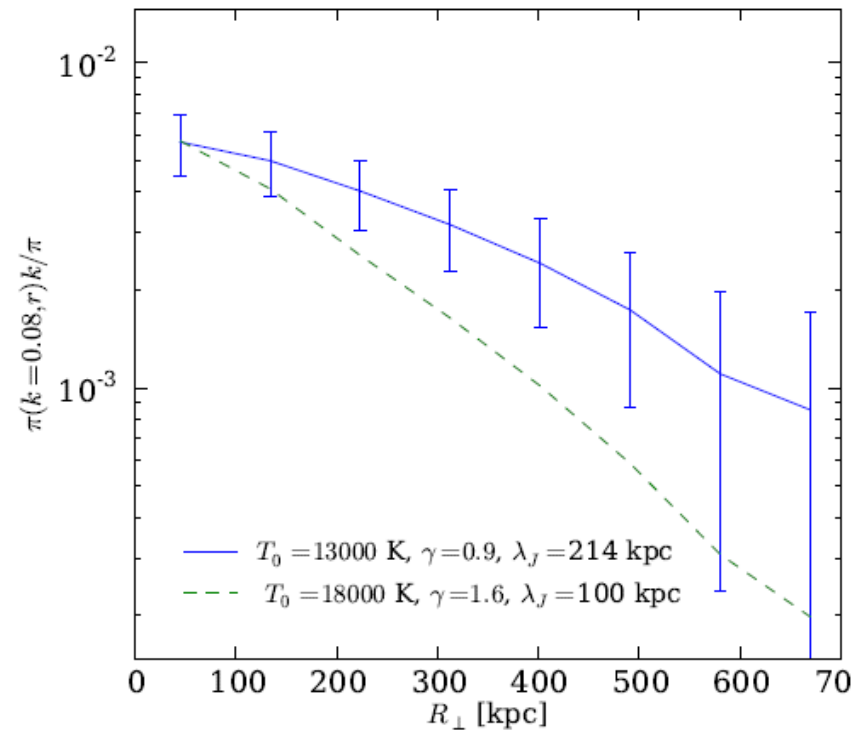


1D Degeneracy Broken by Pairs



Line-of-sight power spectrum

$$P(k) \propto |F(k)|^2$$



Cross power spectrum

$$\pi_{1,2}(k, r_{1,2}) \propto \Re\{\tilde{F}_1^*(k)\tilde{F}_2(k)\}$$

Cross Power: Definition

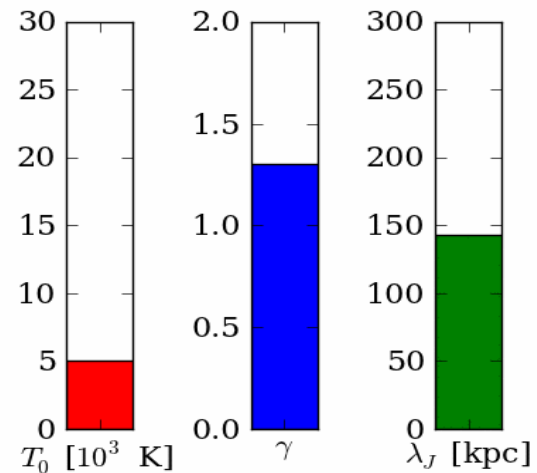
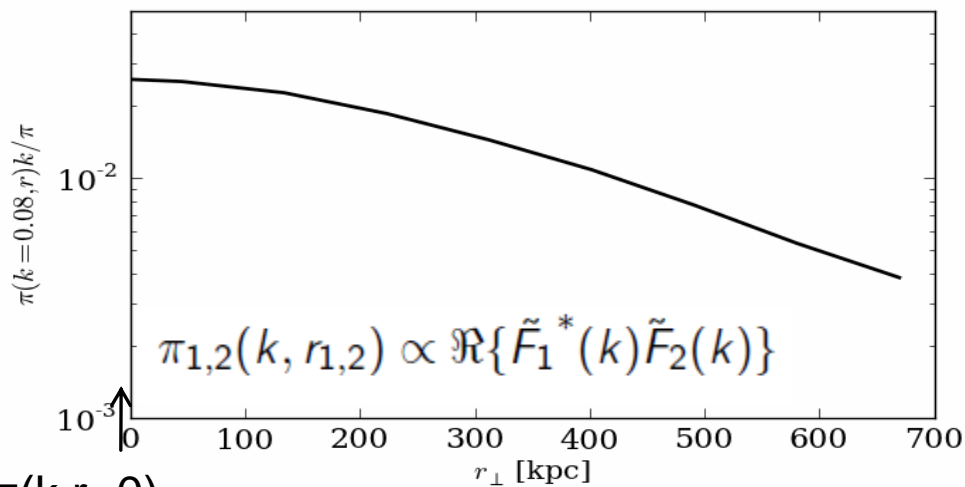
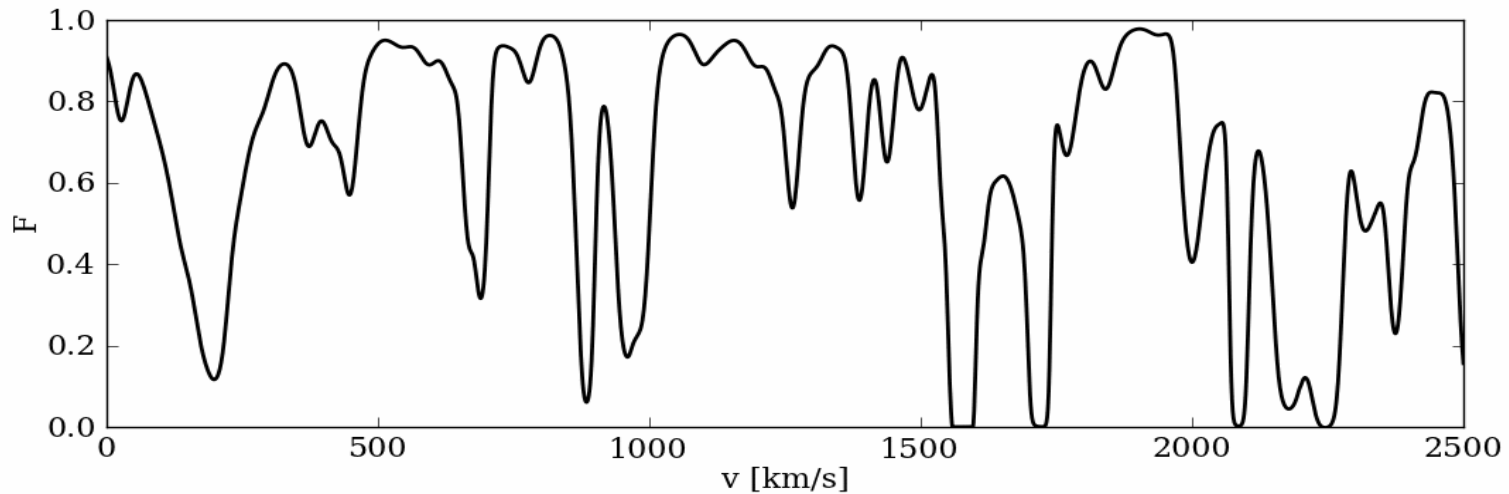
- Fourier transformed of the cross-correlation
- Cross-product of Fourier modes
- Can be expressed in terms of moduli and phases:

$$F_j(k) = a_j(k) \exp[i\Theta_j(k)]$$

$$\pi_{1,2}(k) = a_1(k) a_2(k) \cos(\Theta_1(k) - \Theta_2(k))$$

Cross Power: Sensitivity to 1D smoothing

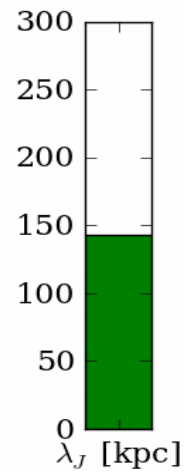
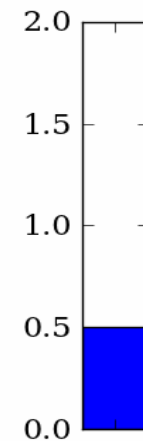
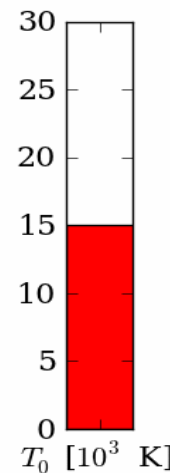
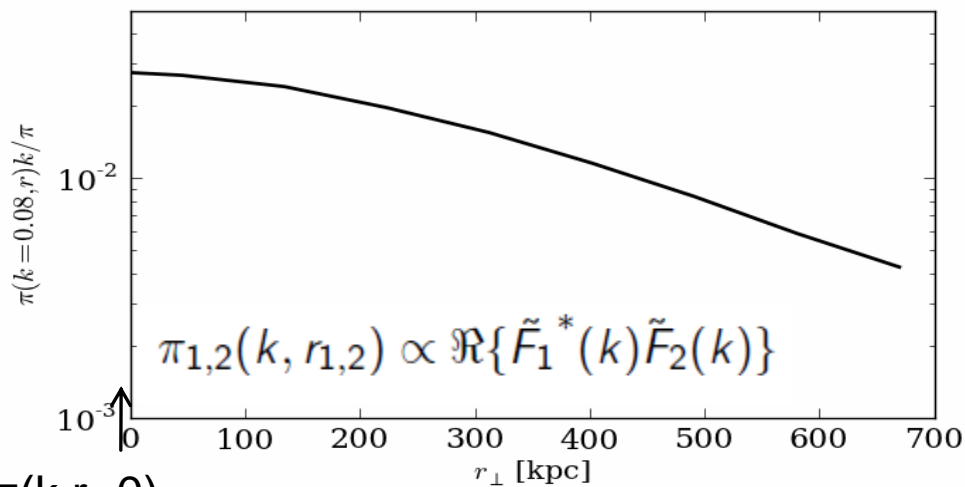
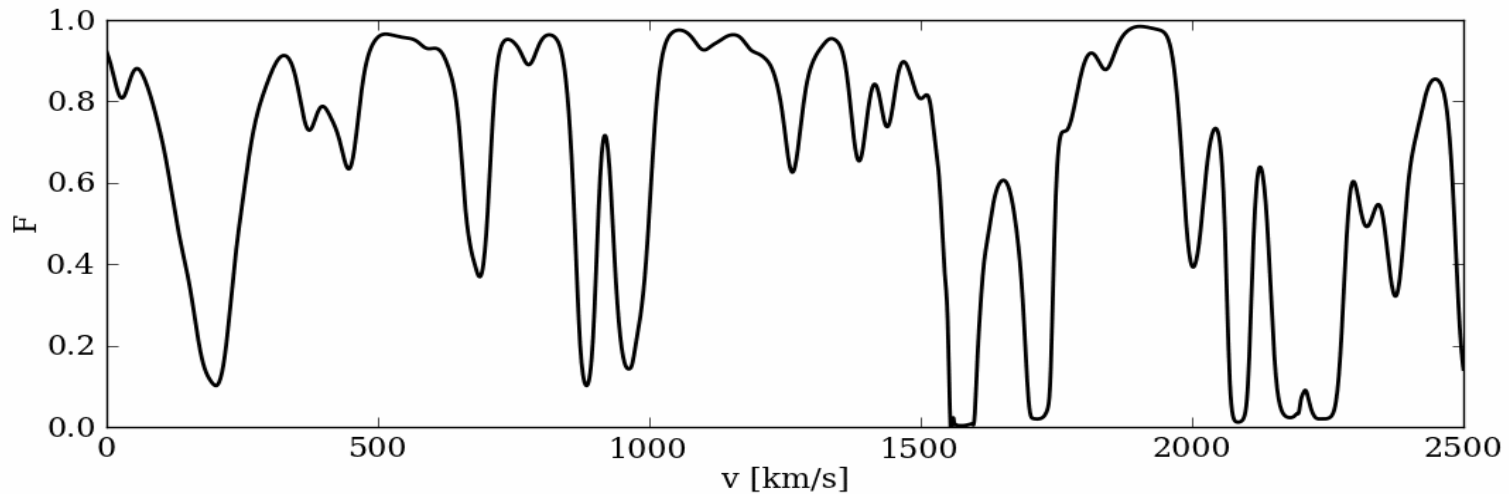
$z = 3$



$P(k) = \pi(k, r=0)$

Cross Power: Sensitivity to 1D smoothing

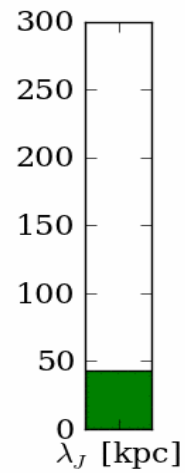
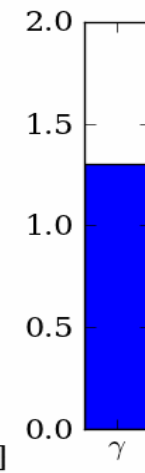
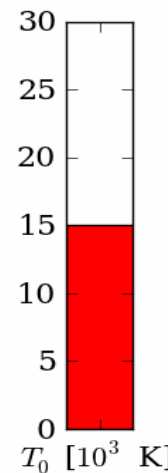
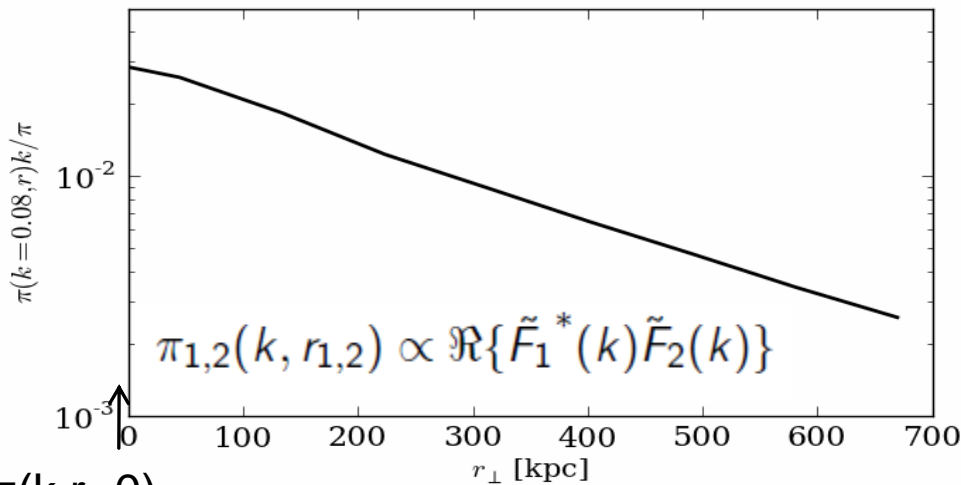
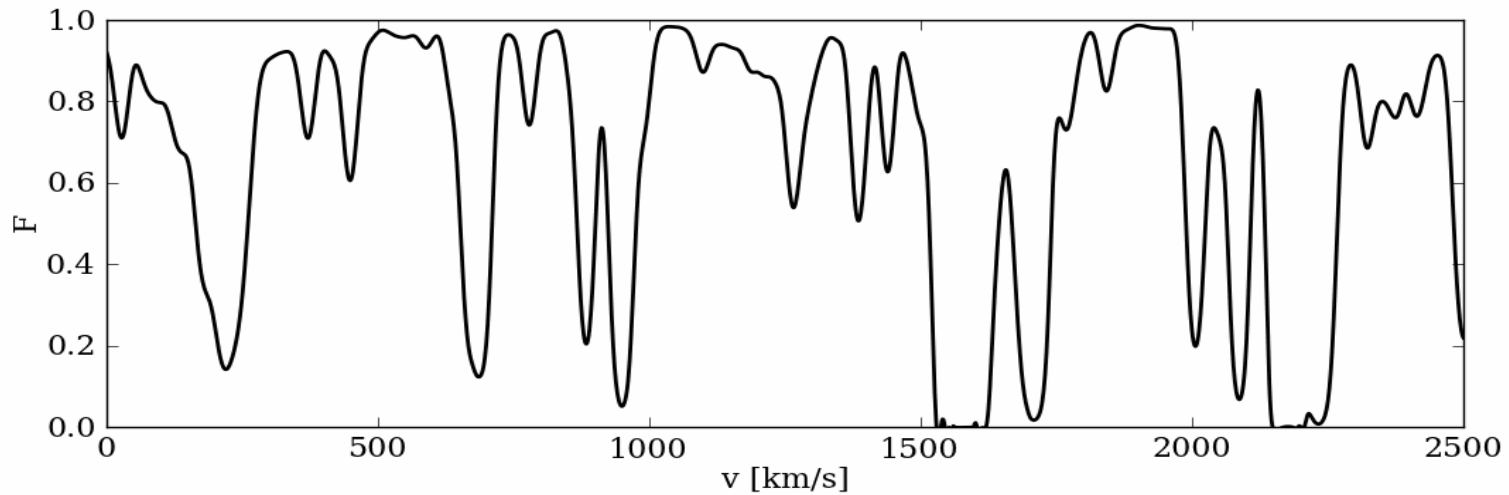
$z = 3$



$P(k) = \pi(k, r=0)$

Cross Power: Sensitivity to 3D smoothing

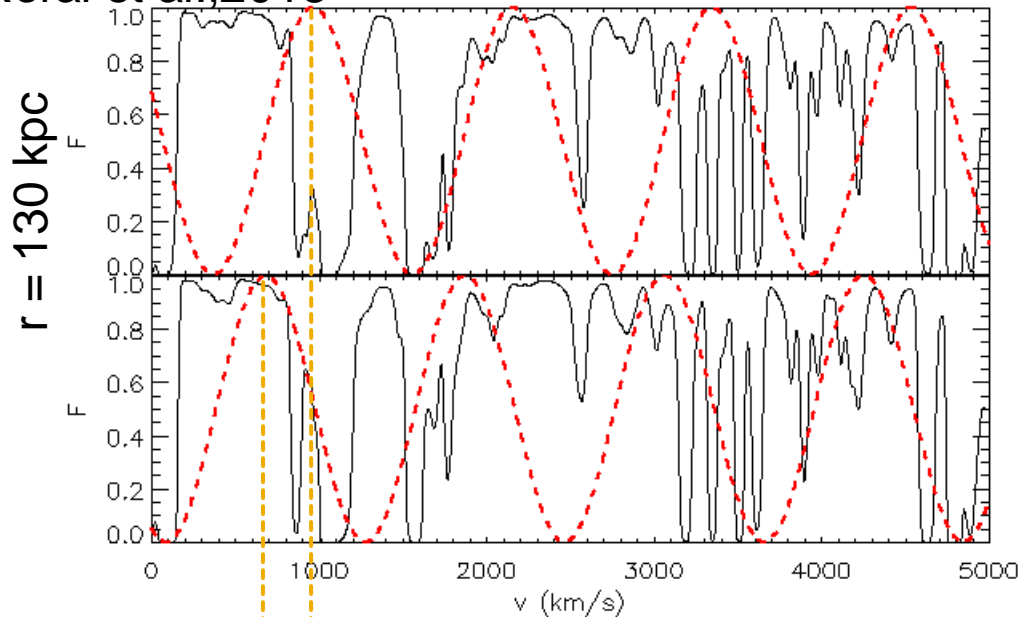
$z = 3$



$P(k)=\pi(k,r=0)$

Isolating the 3D information: Phase Differences

Rorai et al., 2013



Phase difference

$$\pi_{1,2}(k, r_{1,2}) \propto \Re\{\tilde{F}_1^*(k)\tilde{F}_2(k)\}$$

In terms of moduli and phases:

$$\tilde{F}(k) = \rho(k)e^{i\theta(k)}$$

$$\pi(k, r_{1,2}) \propto \underbrace{\rho_1(k)\rho_2(k)}_{\text{amplitudes}} \underbrace{\cos\theta_{12}(k)}_{\text{angular part}}$$

The amplitudes of the modes depend on the LOS power:

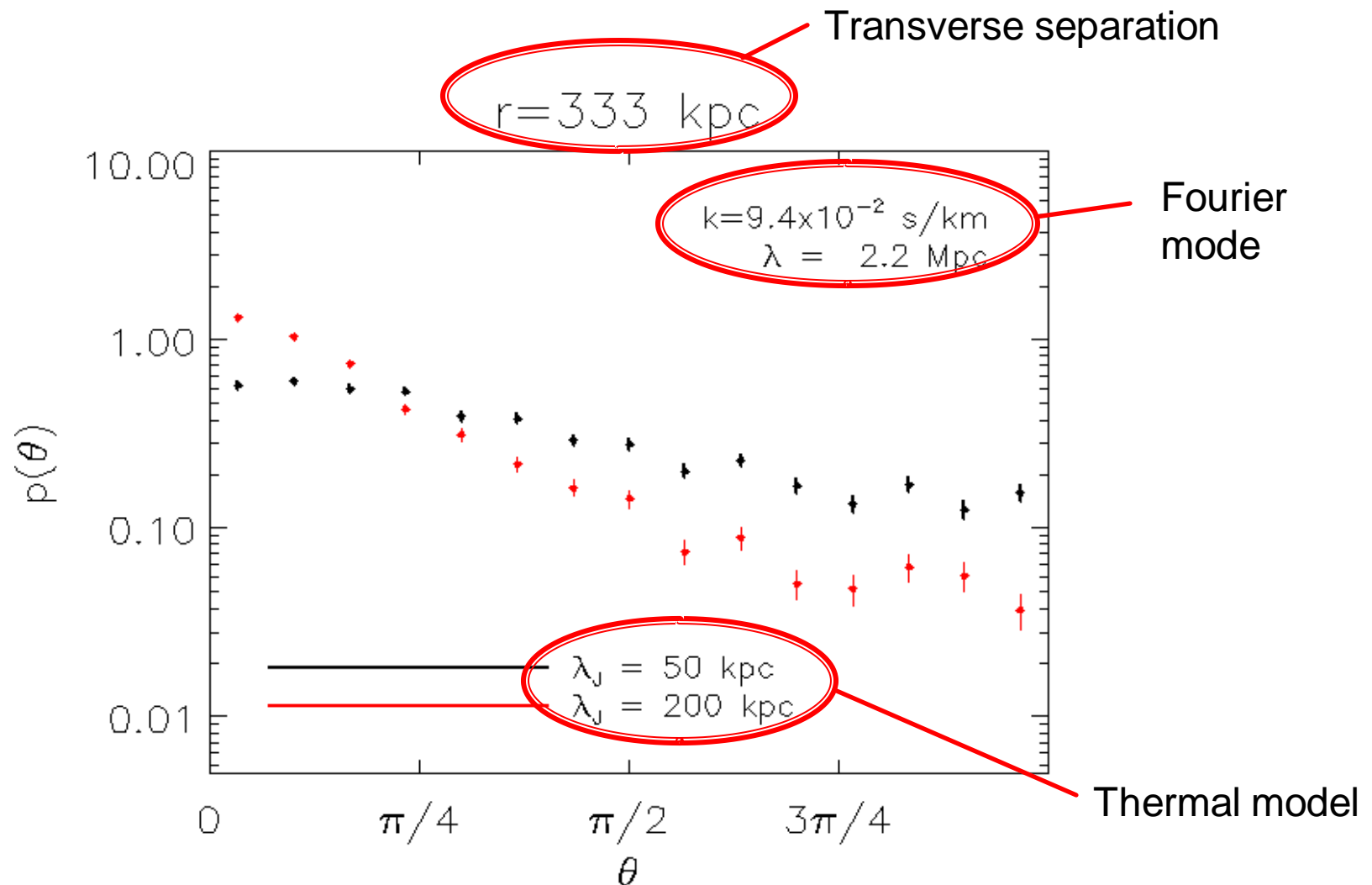
$$P(k) \sim \langle \rho_1\rho_2 \rangle$$

Angular part retains the new 3D information of pairs

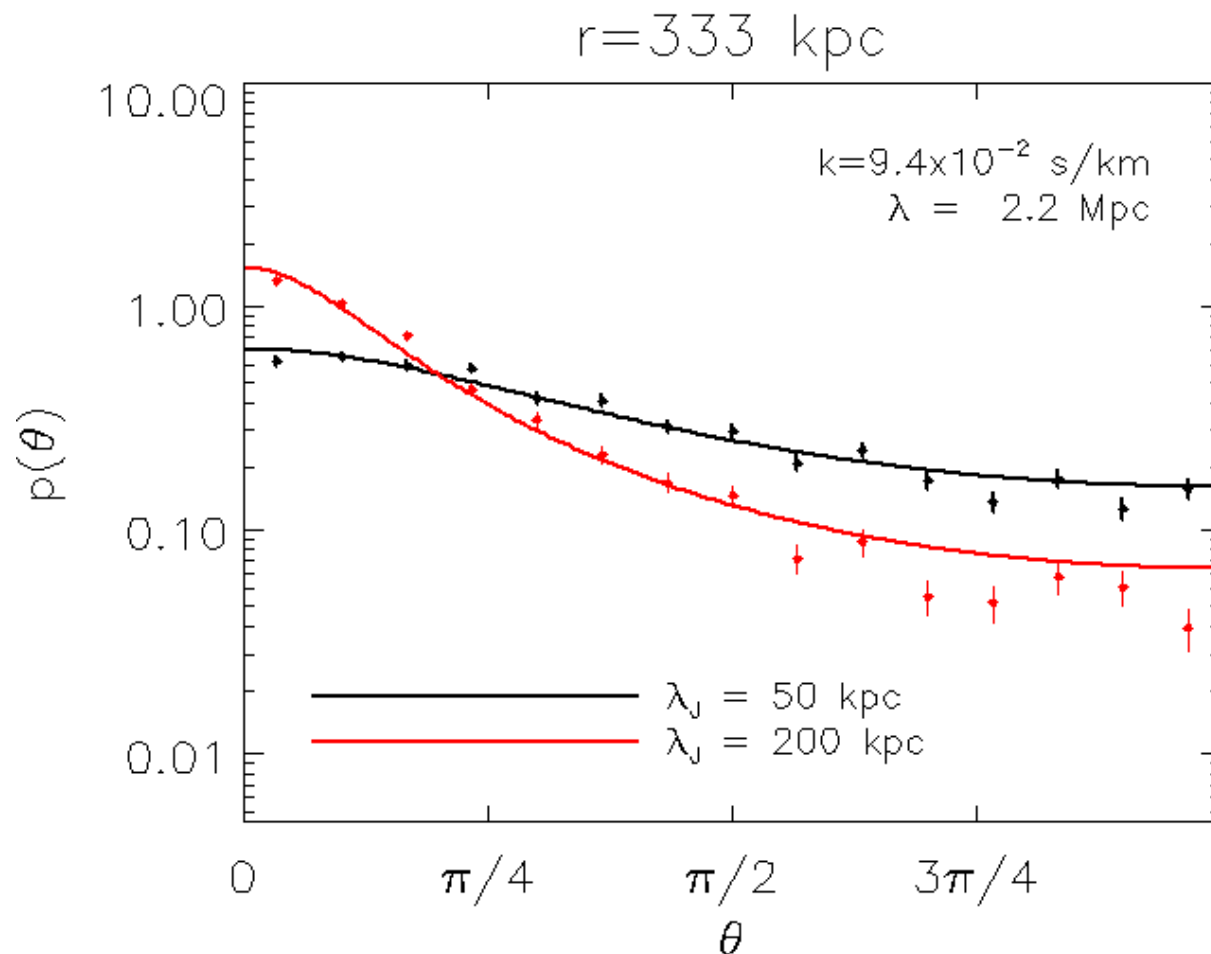
$$\theta_{12}(k) = \arccos\left(\frac{\Re[\tilde{F}_1^*(k)\tilde{F}_2(k)]}{\sqrt{|\tilde{F}_1(k)|^2|\tilde{F}_2(k)|^2}}\right)$$

Problem: what is the best estimator for the angular component?

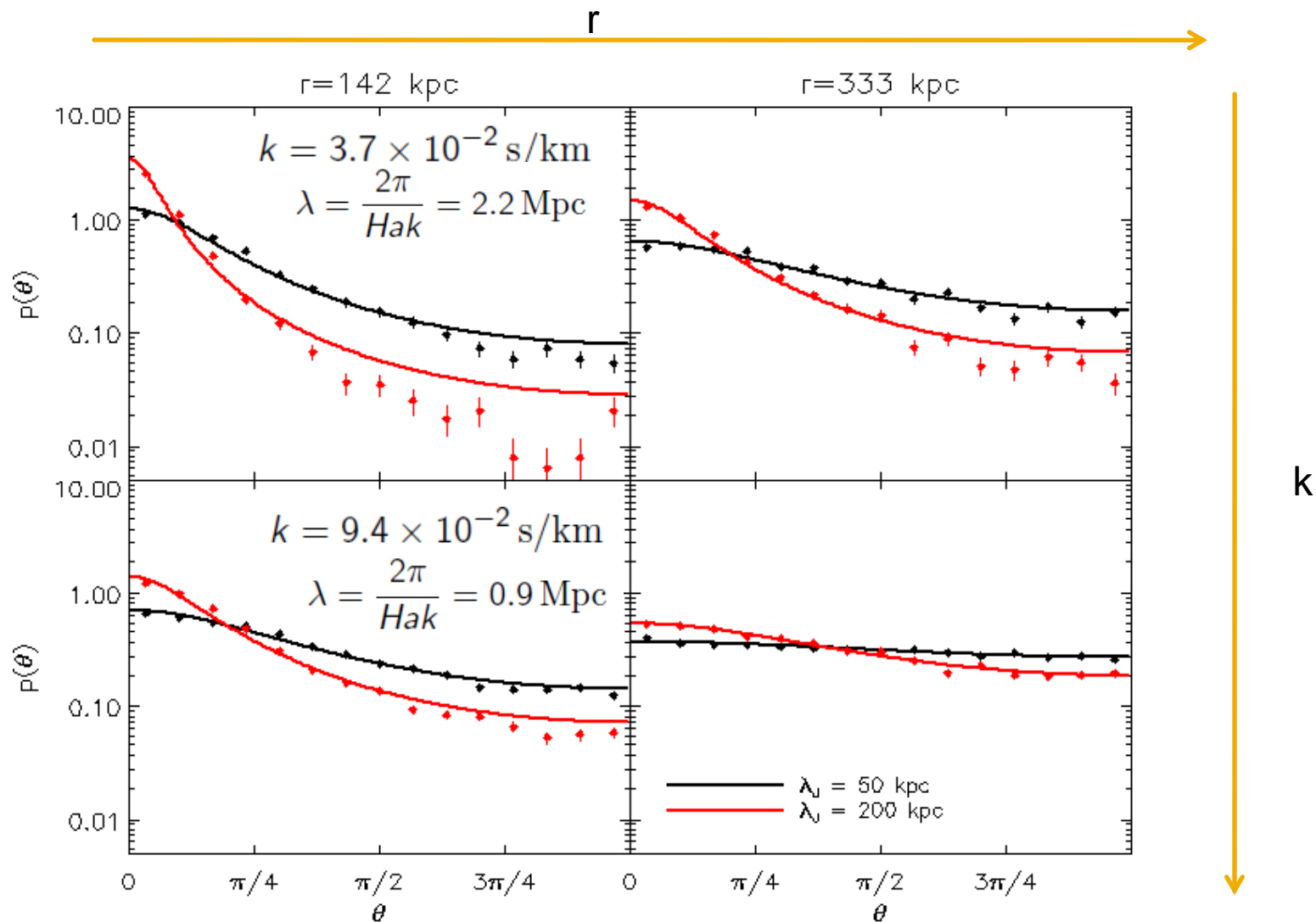
Phase Probability Distributions



Phase PDF: Fit with Wrapped Cauchy



Phase PDF: Dependencies



Bayesian Inference from Phase PDF

1 - Quasar pair data set \rightarrow ensemble of phase differences $\{\Theta(k_i, r_j)\}$

2 – Grid of simulated thermal models \rightarrow Prediction of phase probability distributions as a function of thermal parameters (*Likelihood function*)

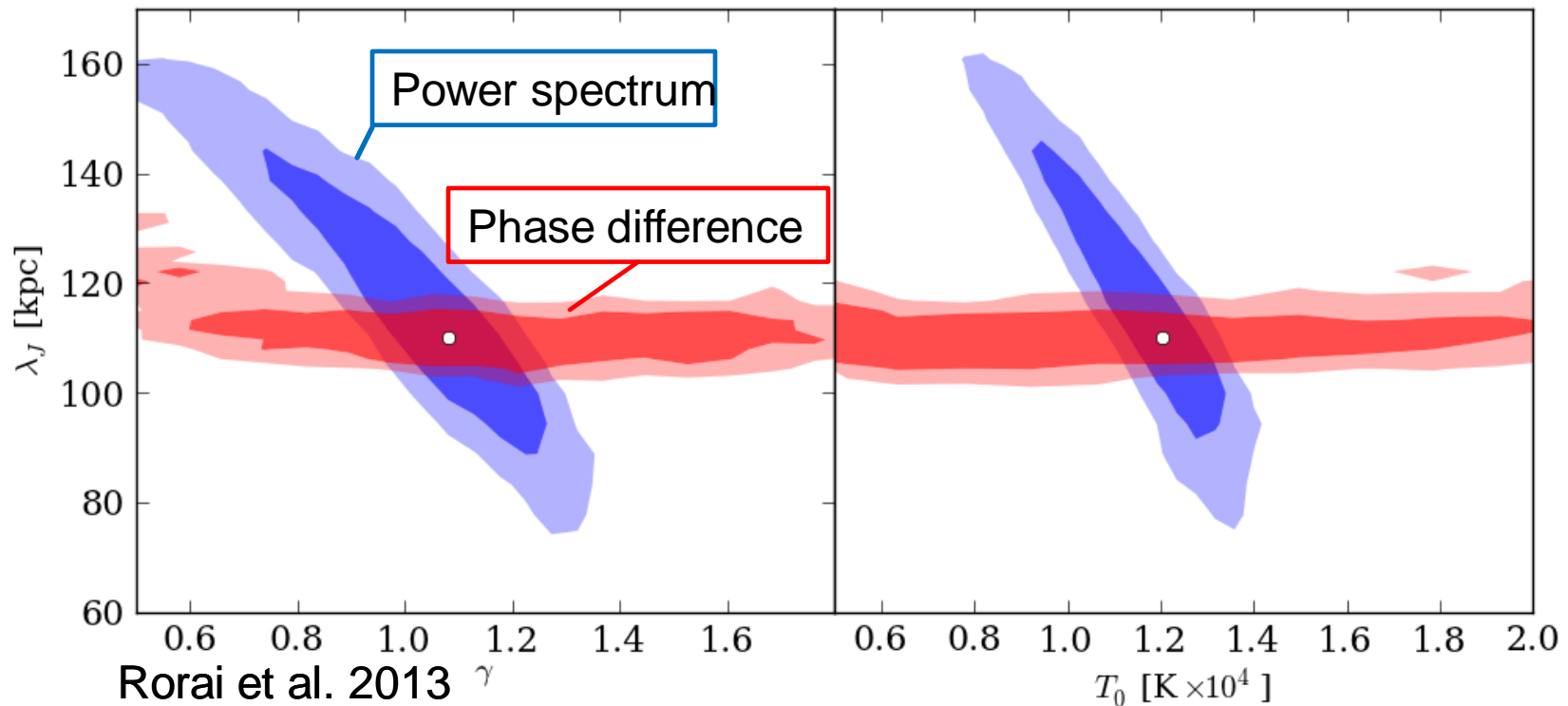
$$L(\{\theta\} | T_0, \gamma, \lambda_J) = \prod_{i,j} P(\theta(k_i, r_j) | T_0, \gamma, \lambda_J)$$

3 - Via MCMC techniques, we do a Bayesian analysis of parameter space, estimating:

- **Degeneracies** between parameters
- **Accuracy** of a measurement

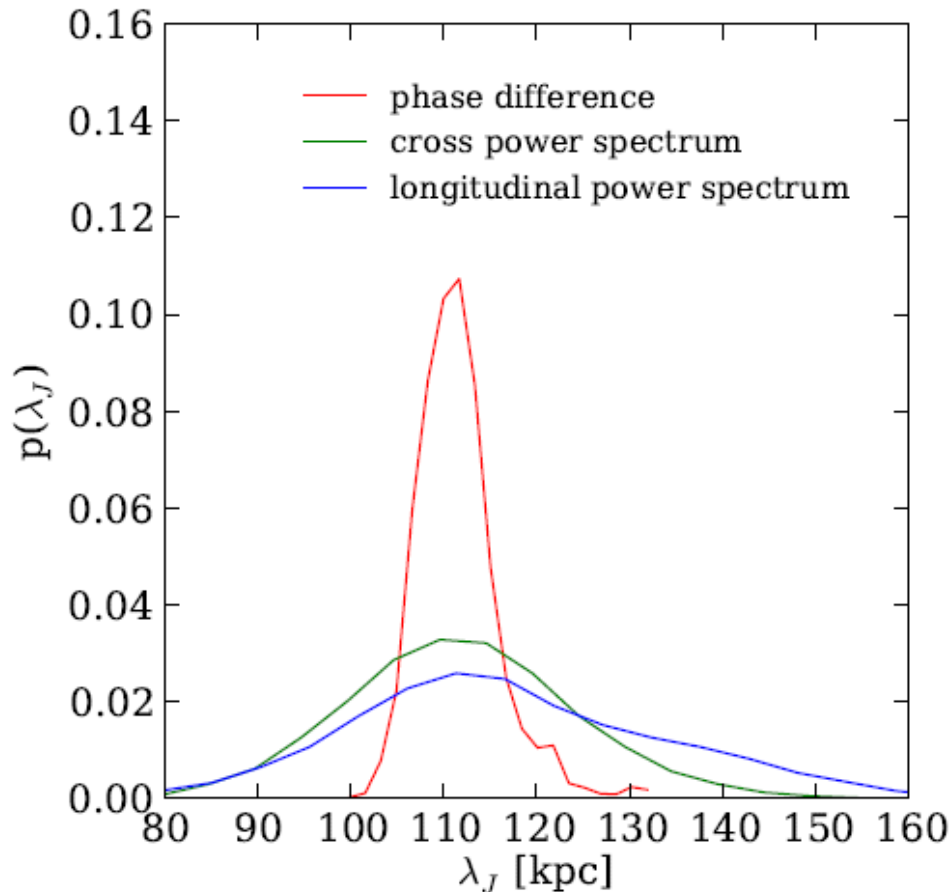
Degeneracies: Study with Mock Data

We assume a sample of 20 spectra (LOS power)
20 pairs (phase difference)



Jeans scale measurements are independent on the equation of state

Predicted Accuracy of the Jeans Scale Measurement



Rorai et al.,2013

We can achieve a precision up to 5% with only 20 pairs

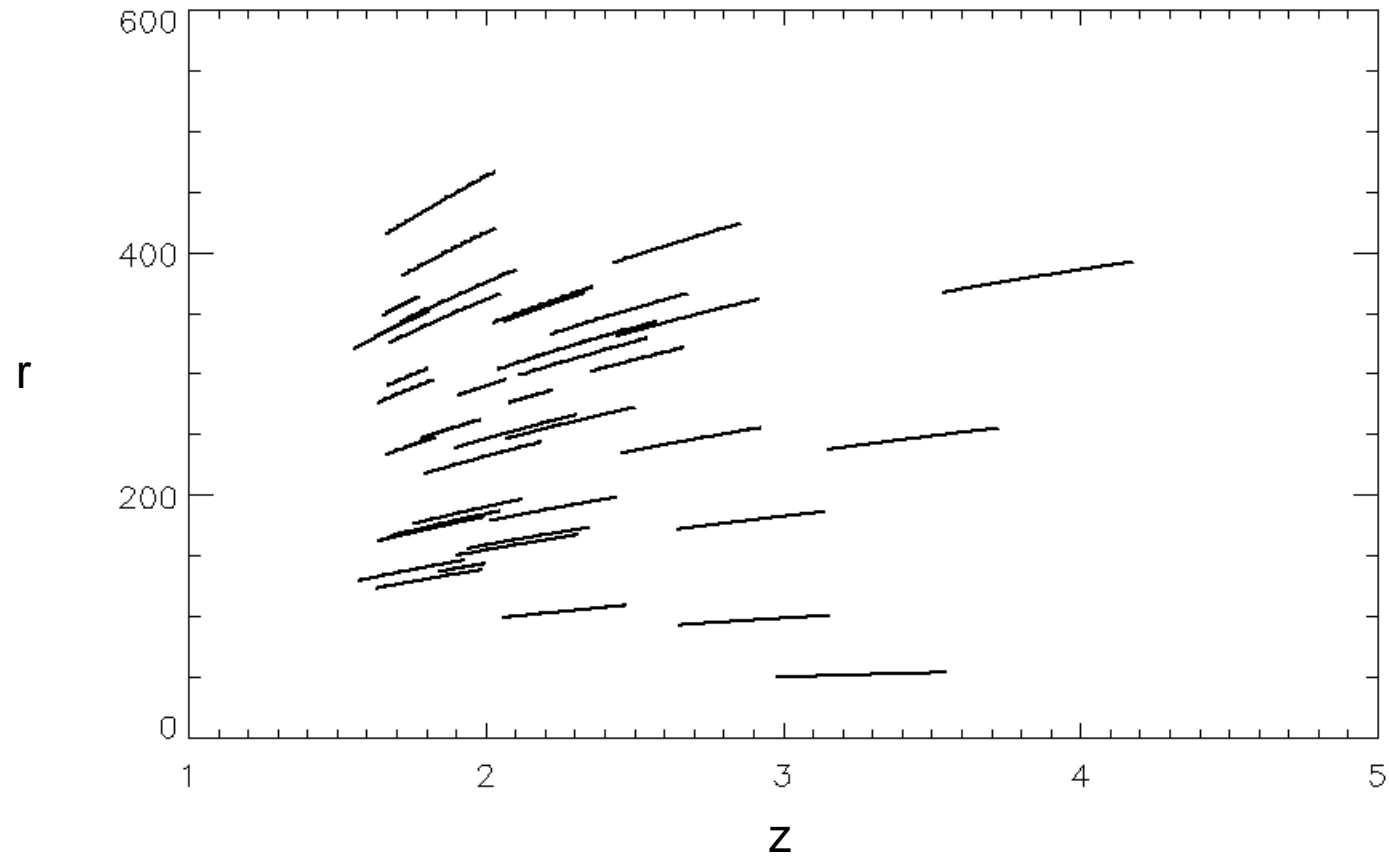
Phase difference statistic is by far more efficient than LOS power and cross-power

Noise decreases the precision by few %

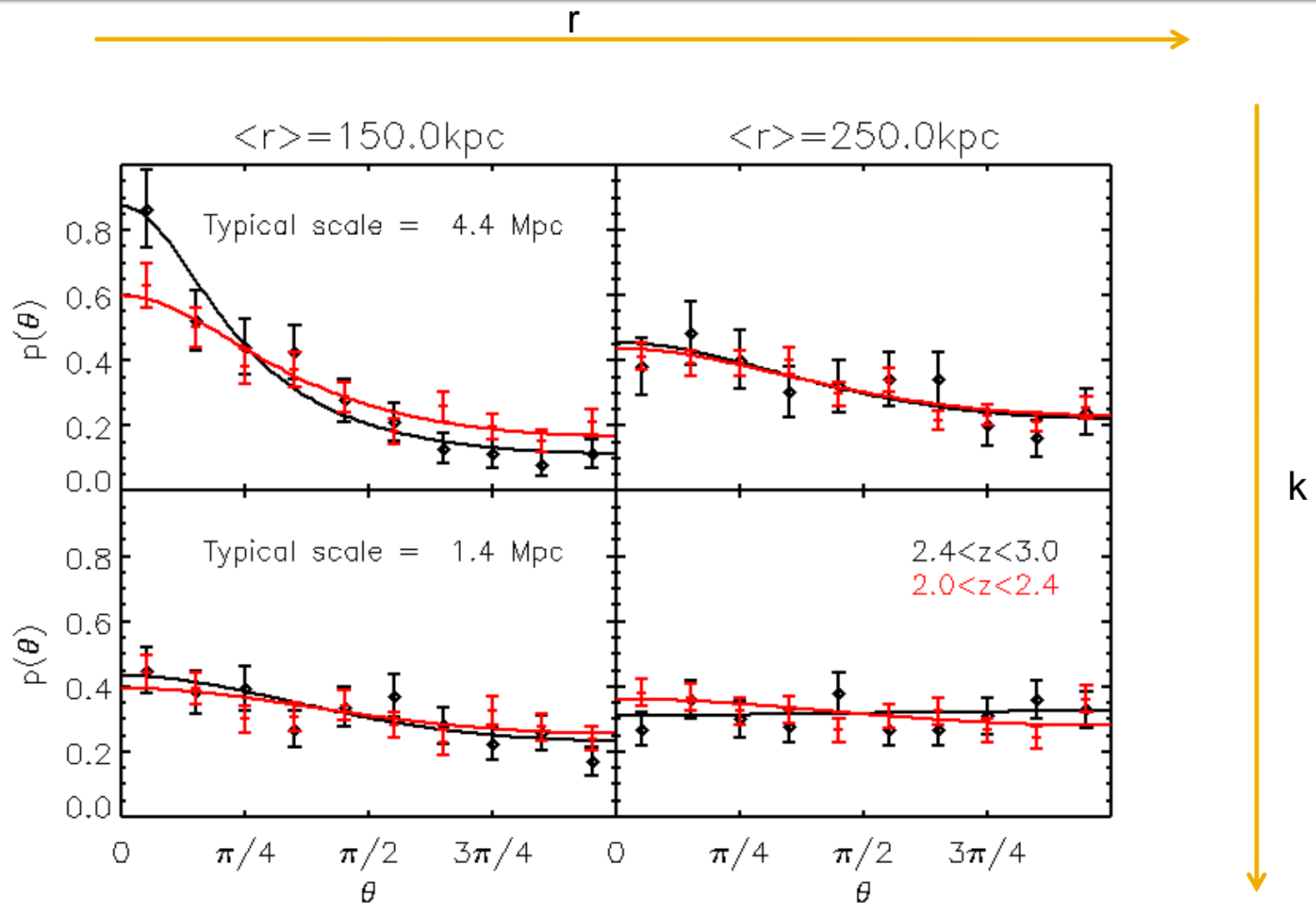
Phase Difference Statistic: Summary

- Cross-power spectrum (and thus cross-correlation) contains line-of-sight information
- Phase differences represent genuine 3D information
- Phase Probability functions follow with good approximation a general functional form
- Phase statistics is very sensitive to the Jeans scale and independent on the EOS

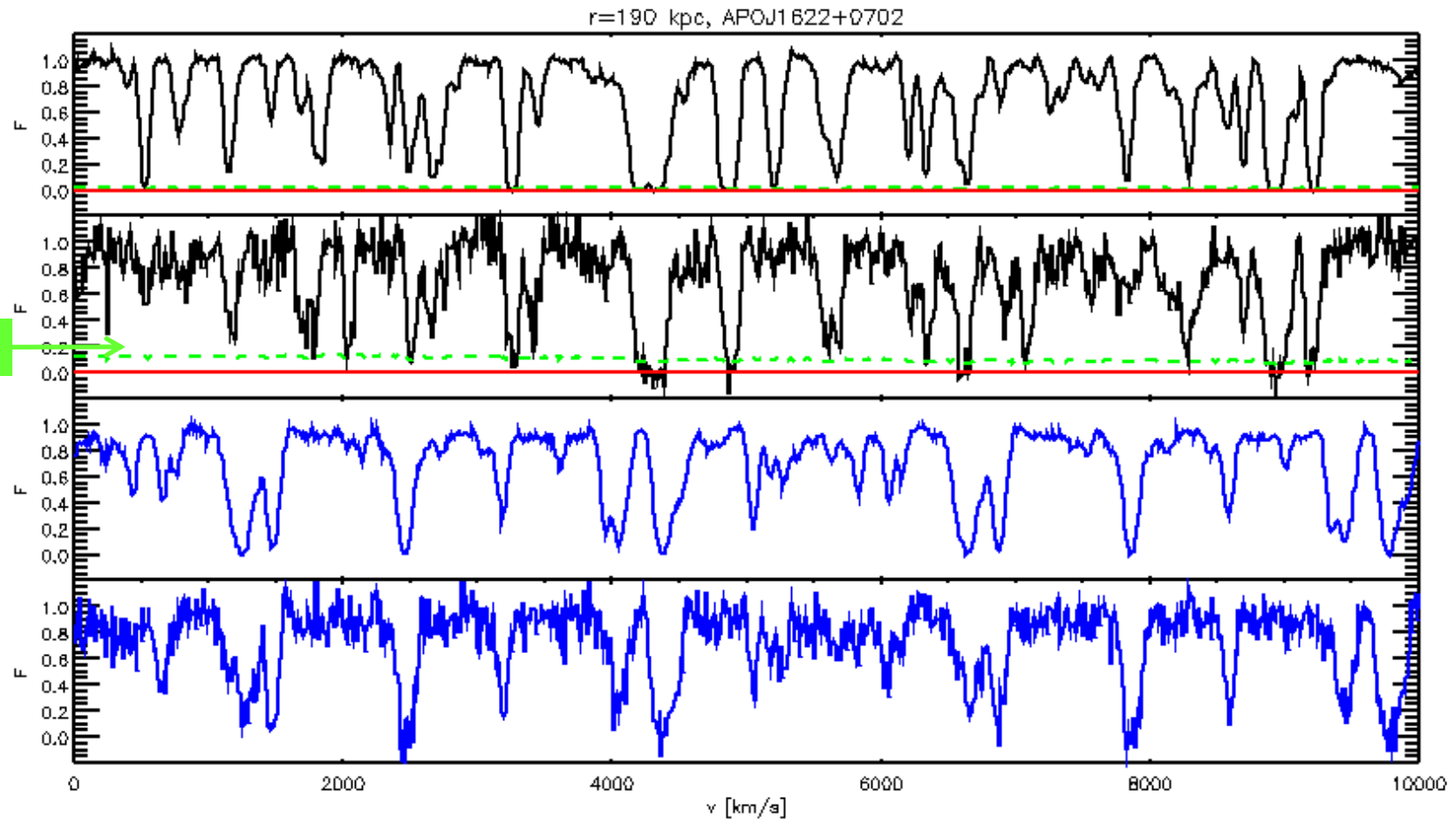
Pair Sample: Overlapping Forest Distribution in z and r



Real Phases: PDF and Dependencies



Forward-Modeling of Simulation

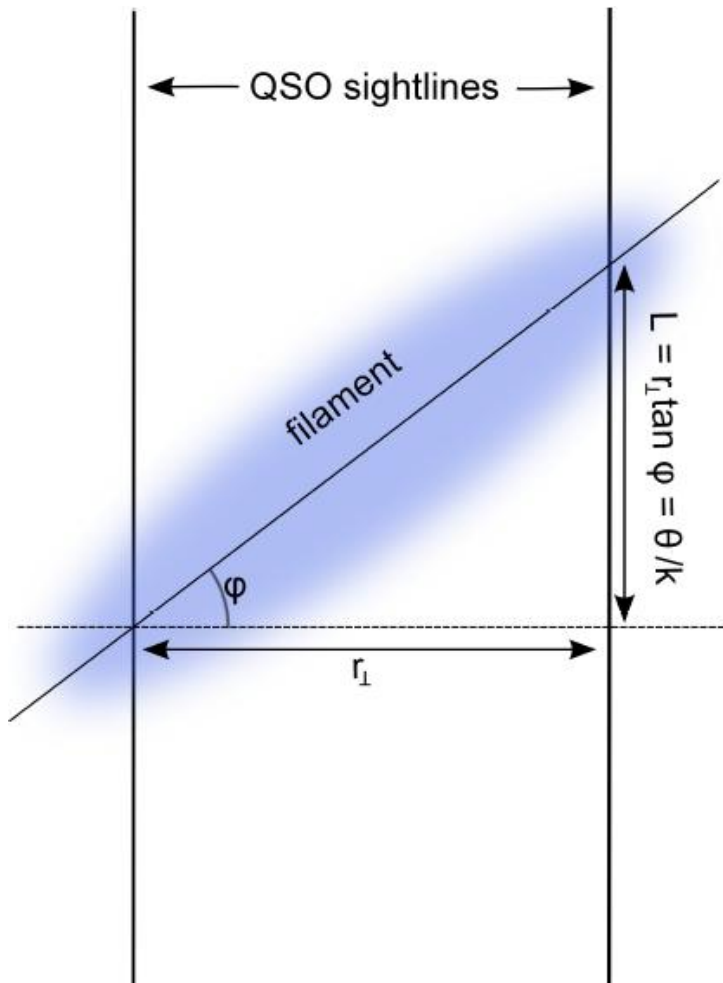


Conclusions

- The Jeans scale of the IGM is of key importance in cosmology:
 - Threshold for **galaxy formation**
 - Record of the **thermal history**
- Close quasar pairs open new possibilities to probe the small-scale structures of the universe
 - Phase analysis will provide a **first precise measurement of the Jeans Scale**
 - Phase statistic **insensitive to the equation of state** of the IGM
 - Combinations with other Ly- α forest statistics can **break degeneracies** and improve constraints on the thermal history
- Small Jeans scale?
 - Preliminary results point toward $\lambda_J < 60\text{-}40$ kpc at $z=2\text{-}3$
 - This would imply an **abundance of small-scale structures**
 - Comparison with hydro simulations will be made in the future

Thank you!

Statistical Distribution of Phases: Ansatz



- A filament produce absorption features in two close spectra

- Phase differences are driven by the orientation of these filaments

- If the orientation ϕ is uniformly distributed, then the phase Θ follows a *wrapped Cauchy distribution*

$$P_{WC}(\theta; \zeta) = \frac{1}{2\pi} \frac{1 - \zeta^2}{1 + \zeta^2 - 2\zeta \cos(\theta)}$$

Defines a 1-parameter family of functions

Nice Properties of Phase Differences

- Invariant under convolution with symmetric Kernel W

$$W(k) = \int W(x) \cos(kx) + i \int \underbrace{W(x)}_{\text{even}} \underbrace{\sin(kx)}_{\text{odd}} = \text{real} \longrightarrow \text{No change in phases}$$

- Thermal Broadening has very small effect
- Resolution does not need to be precisely modeled

- Sensitive to the Jeans scale also at low k

- High-resolution spectra are not required if we resolve the Jeans scale in the transverse dimension

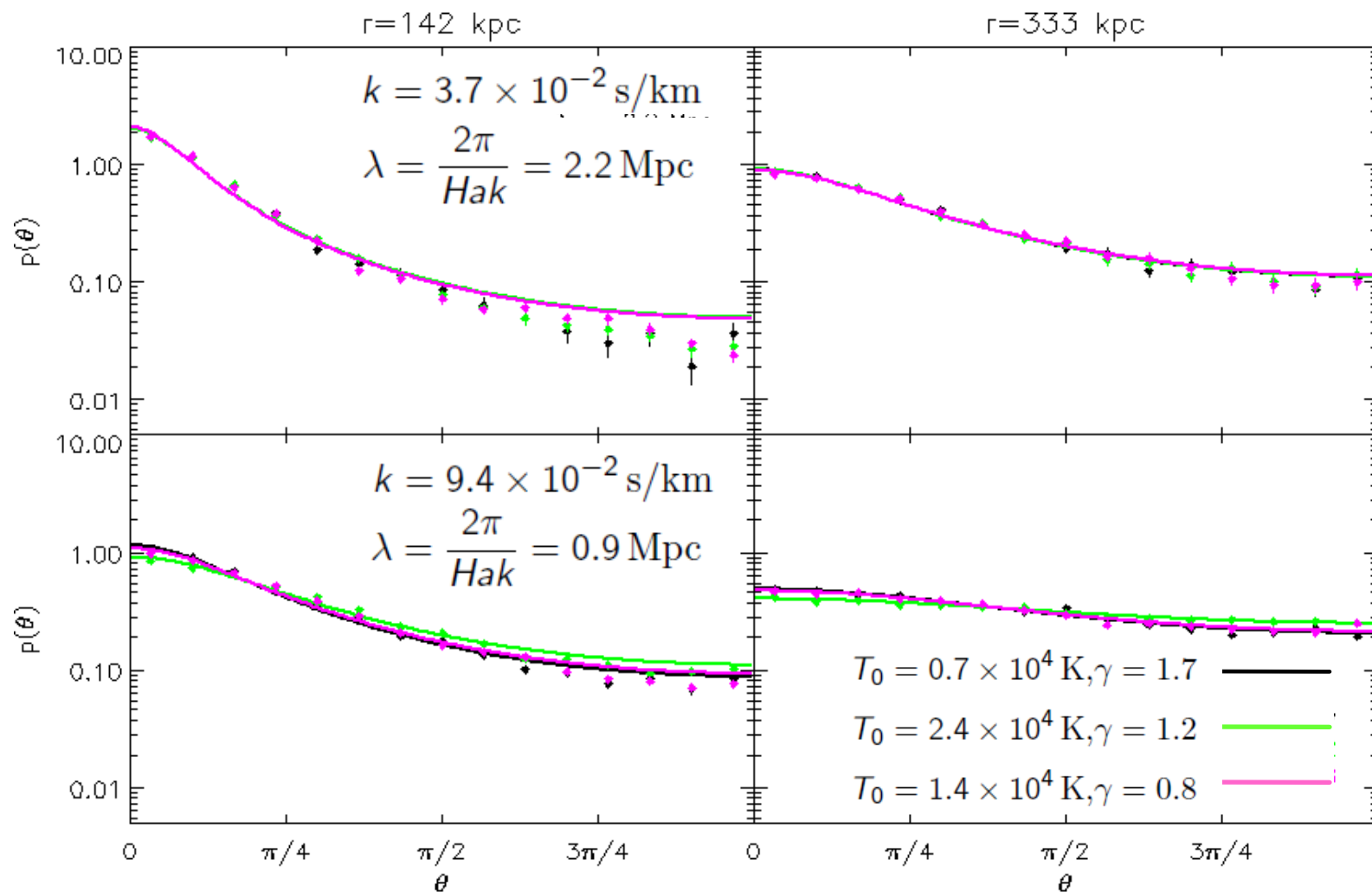
$$P(k) = \int_k^\infty 2\pi k' P_{3D}(k') dk'$$

- Invariant under rescaling of the *normalized* flux

$$\theta_{12}(k) = \arccos \left(\frac{\Re[\tilde{F}_1^*(k) \tilde{F}_2(k)]}{\sqrt{|\tilde{F}_1(k)|^2 |\tilde{F}_2(k)|^2}} \right)$$

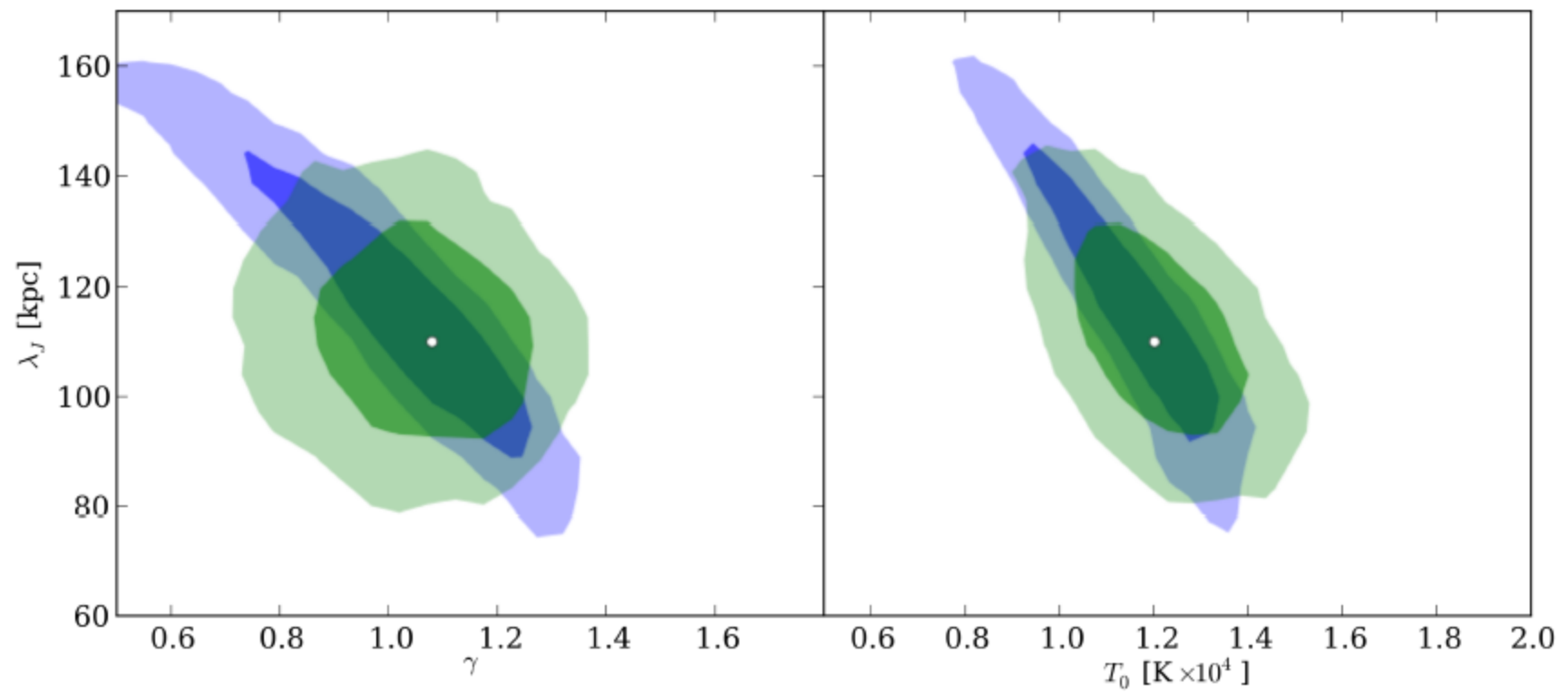
- Robustness to uncertainties in continuum fitting

Sensitivity to the Equation-of-State Parameters



Phases are almost independent of the equation of state

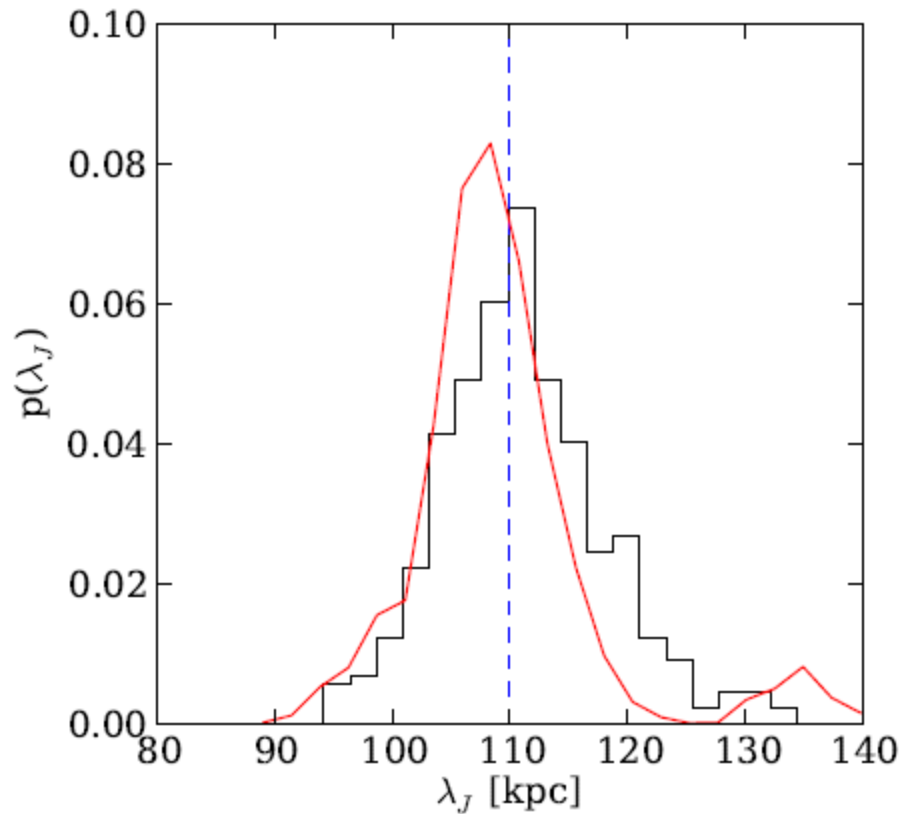
Sensitivity of the Cross-Power Spectrum



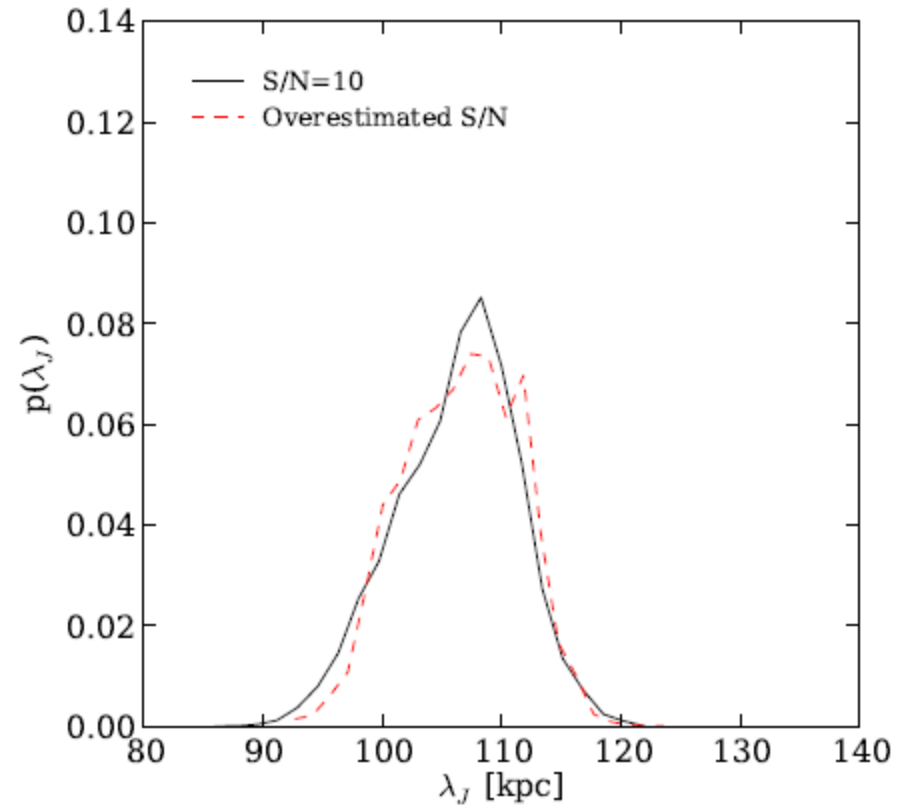
The Cross-Power is partially affected by the line-of-sight degeneracies

Bias tests

Repeating our “experiment” for 400 different mock dataset shows that our method is not biased

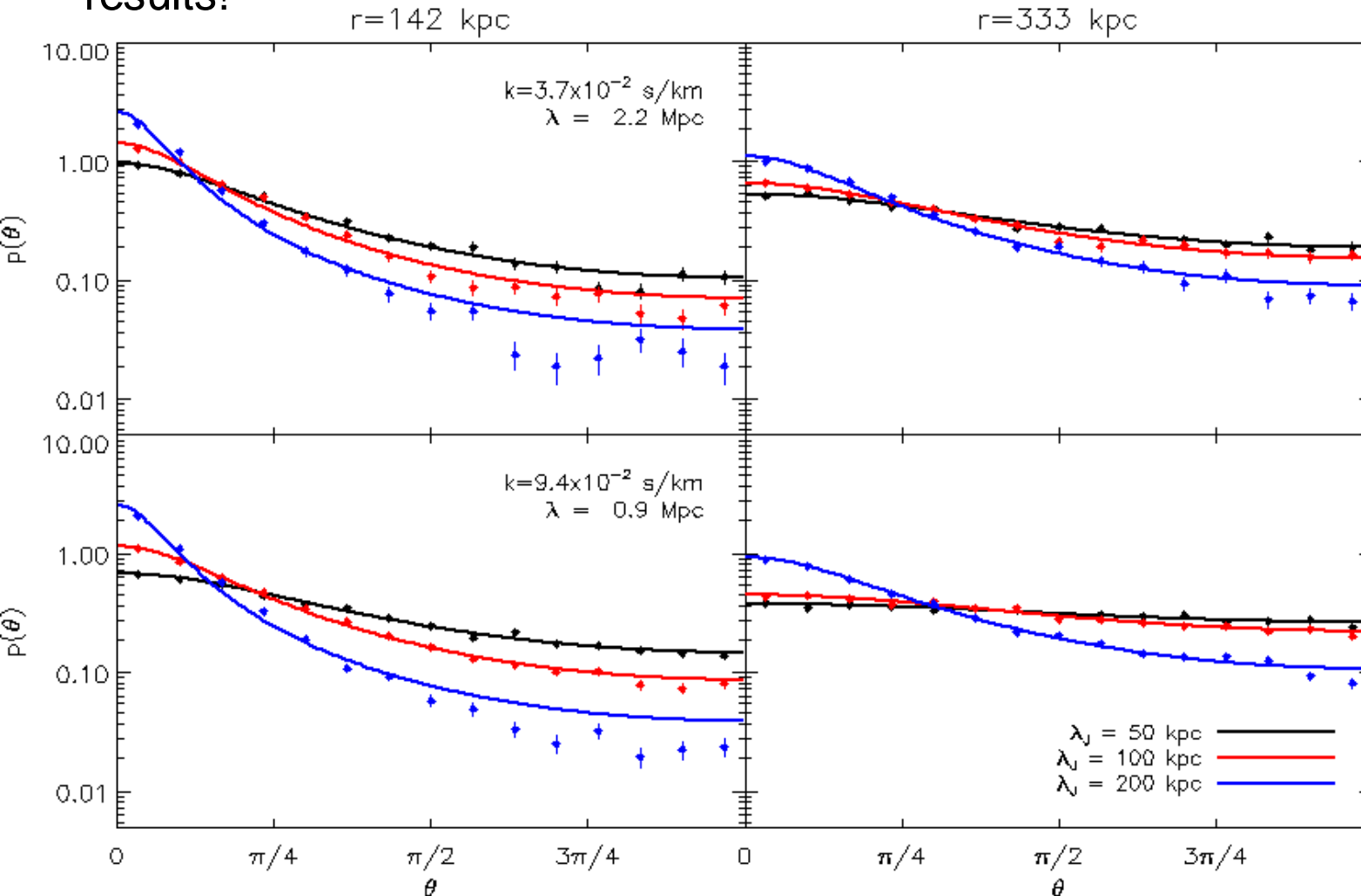


Overestimating the S/N by 20% does not significantly modify the outcome of the measurement.



Sensitivity to the Jeans scale

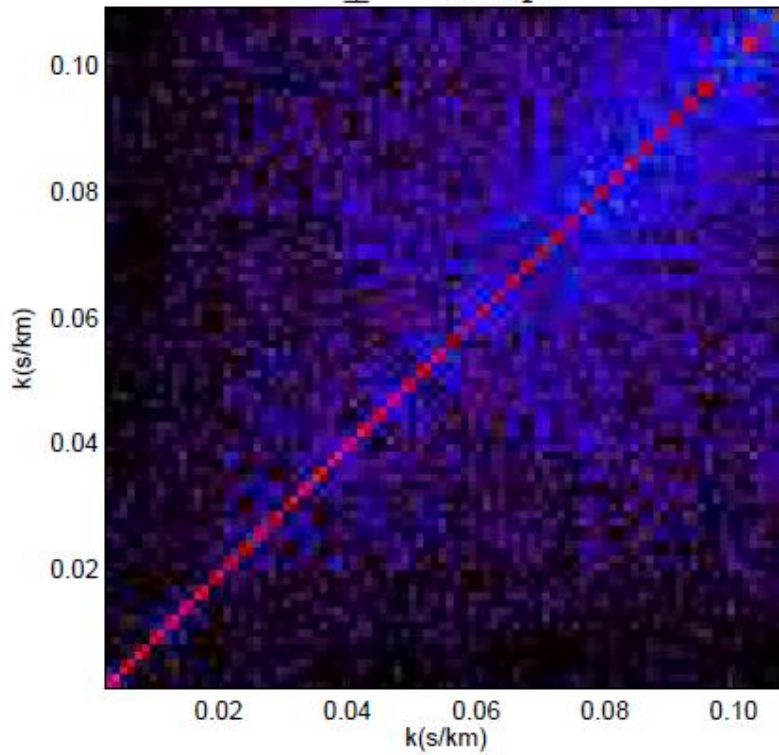
The wrapped Cauchy function provides an excellent fit to the simulation results!



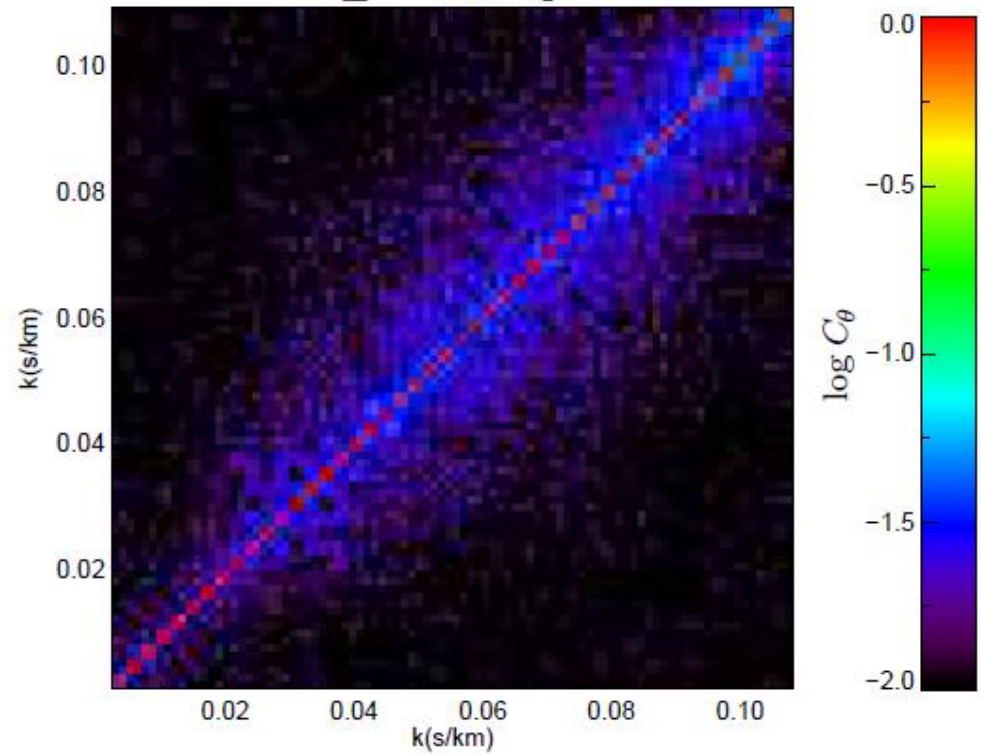
Phase difference distribution for DENSITY, in real space

Are Phases Independent?

$r_{\perp} = 70$ kpc

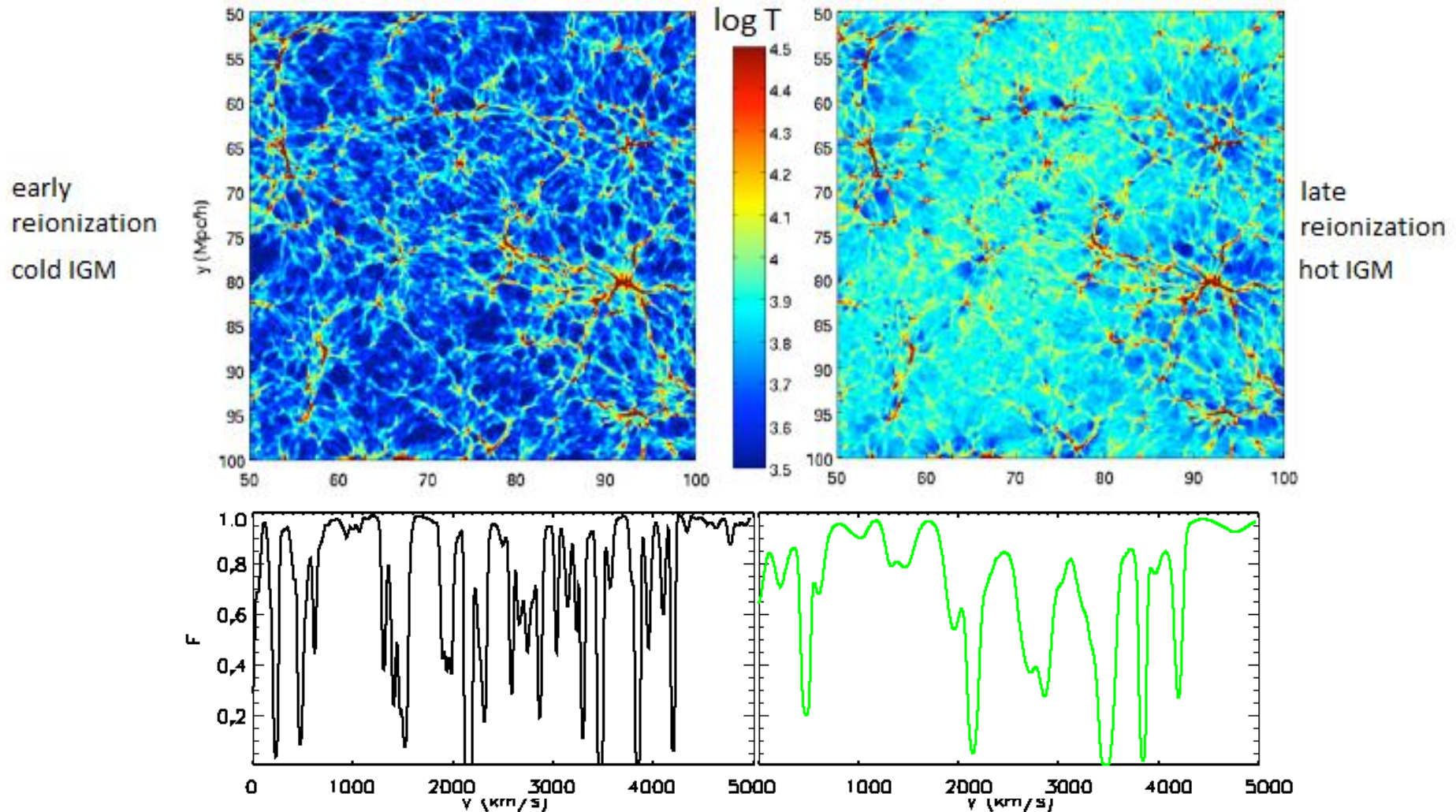


$r_{\perp} = 430$ kpc



Correlation factors are at most of the order of 3%

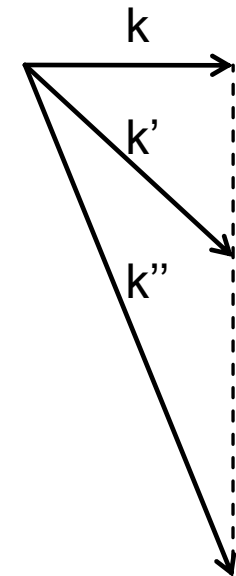
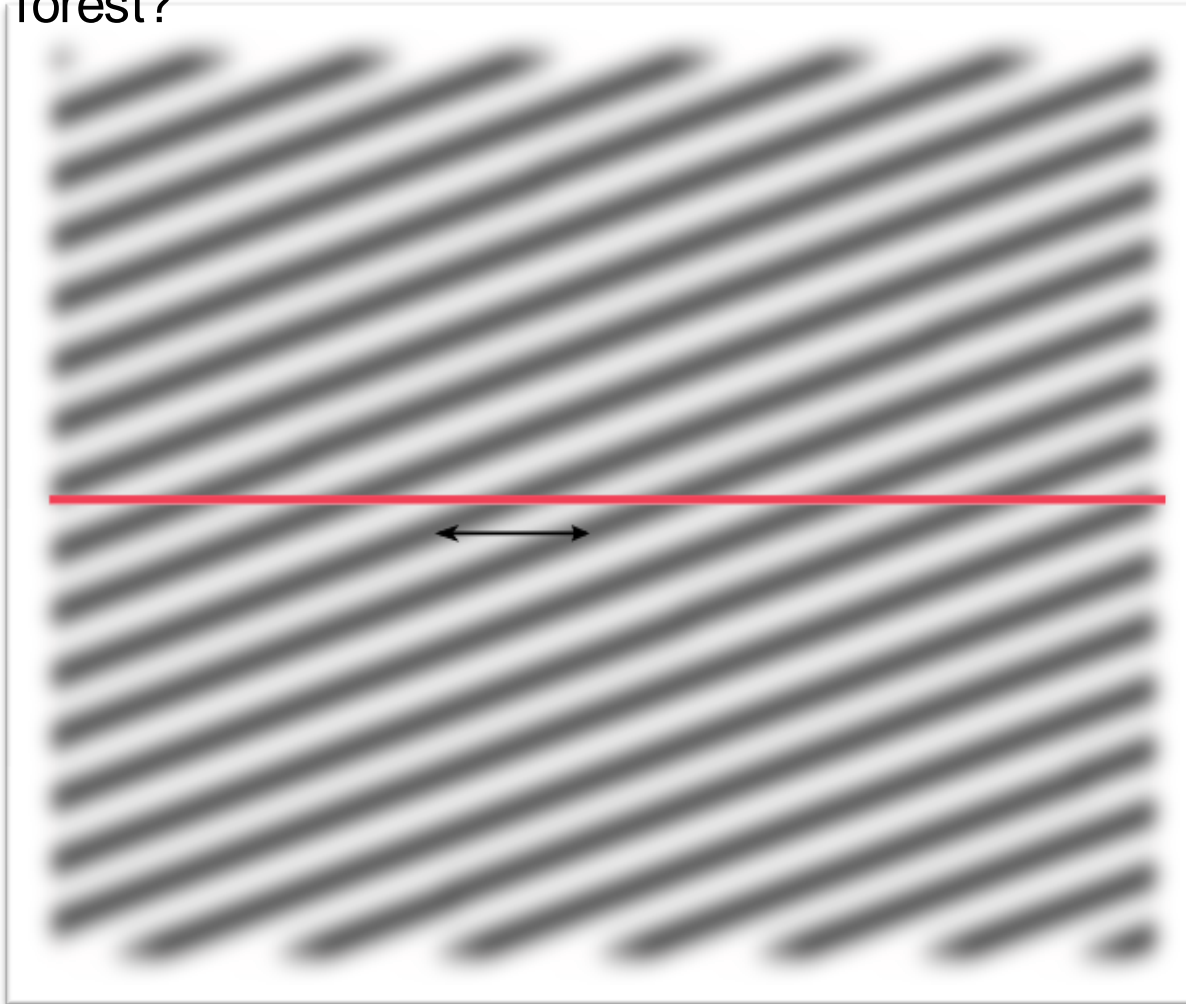
Thermal state of the IGM and reionization history



The **Lyman α forest** can be used to constrain the thermal state of the IGM

Sensitivity to the Jeans scale

Why does pressure smoothing affects large-scale modes of the forest?



High-k 3D modes projects into low-k modes in 1D spectra