

Lecture 21

Big Bang nucleosynthesis

Borusmetto, chap VI + A3

WHY BBN?

Stars like the sun have 25% of their mass in Helium:

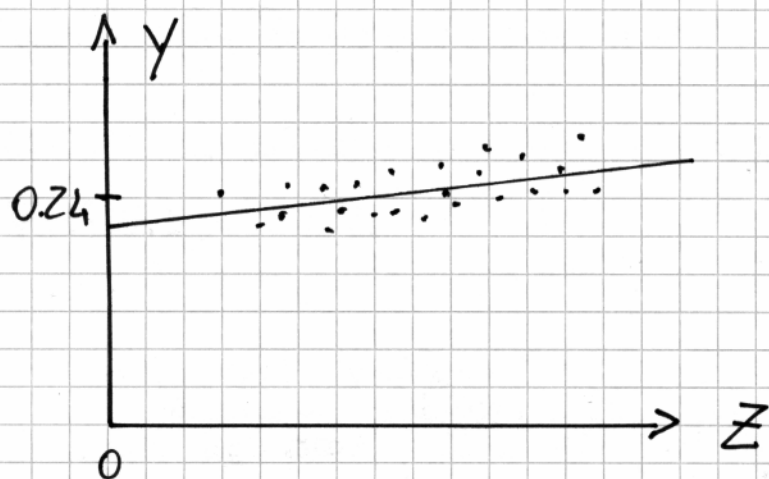
$$X \equiv \frac{M_H}{M_{TOT}} \approx 0.73, \quad Y \equiv \frac{M_{He}}{M_{TOT}} \approx 0.25, \quad Z \equiv \frac{M_{met}}{M_{TOT}} \approx 0.02$$

Here we call "metals" all elements more massive than He

It is believed that metals were synthesised in stars

But stars are ~5% of baryons, it's hard to have a ~25% of mass in He with stellar nucleosynthesis.

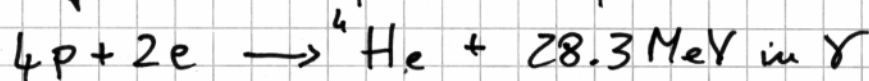
Helium is present also in stars with very low metallicity, Z



So even a "primordial" star, with $Z=0$, would have Helium:

$$Y_p = 0.238$$

Fusion of H to He produces photons:



so $\frac{\Delta m}{m} = 0.75\%$. To produce $Y = 0.24$ one would transform 0.18% of baryons to photons

$$\Omega_{\gamma, He} = \Omega_b \times 0.0018 \approx 9 \times 10^{-5} \approx \Omega_{\gamma 0}$$

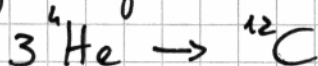
this radiation is not observed.

So 23.8% of baryons should be processed to He in the early universe.

WHY ONLY He (AND LIGHT ELEMENTS)?

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To go beyond He one has to trigger the 3α reaction:

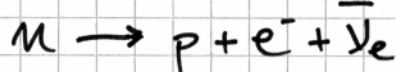
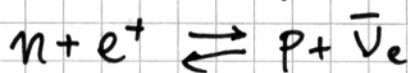
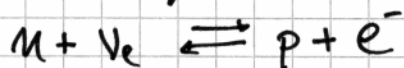


but this requires a higher temperature, while the temperature in cosmology keeps decreasing.

Heavier elements can only be synthesised inside stars, where gravitational contraction makes temperature increase

FREEZING OF n/p RATIO

p and n appear after the Q-H transition, and are kept in equilibrium by:



Their number densities require chemical potentials to be known.

But their number ratio is (demonstrated in Vittorio):

$$\frac{n_n}{n_p} = \exp\left(-\frac{\Delta mc^2}{k_B T}\right) \left(\frac{m_n}{m_p}\right)^{3/2}, \quad \Delta mc^2 = 1.3 \text{ MeV}$$

When ν 's decouple at $T_{\nu, \text{dec}} = 300 \text{ keV}$, the first two channels become ineffective, while neutron decay has a time scale of 15 min, $\gg 0.9 \text{ s}$ that is the age of the universe then

$\Rightarrow n_n/n_p$ freezes

Let's define: $X_n \equiv \frac{n_n}{n_b}$, $X_{n, \text{fr}} = \frac{1}{1 + \frac{n_p}{n_n}|_{\text{fr}}}$

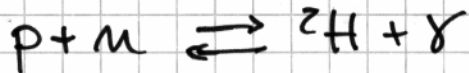
$$\left.\frac{n_n}{n_p}\right|_{\text{fr}} = \exp\left(-\frac{1.3 \text{ MeV}}{0.9 \text{ MeV}}\right) \approx 0.23$$

$$X_{n, \text{fr}} \approx 0.18, \quad X_n = X_{n, \text{fr}} \exp\left(-\frac{t - t_{\nu, \text{dec}}}{\tau_\beta}\right)$$

where $\tau_\beta = 887 \text{ s} \sim 1000 t_{\nu, \text{dec}}$

THE DEUTERIUM BOTTLENECK

The first reaction that takes place is:



This reaction destroys ${}^2\text{H}$ if γ has an energy $E_\gamma \geq 2.2 \text{ MeV}$

This energy corresponds to a temperature of $\sim 2.5 \times 10^9 \text{ K}$

Being $> 0.5 \text{ MeV}$, this should happen in the lepton era.

But $n_\gamma \gg n_b$ because $\eta \ll 1$

Note: after e^-e^+ annihilation n_γ and n_b both evolve $\propto a^{-3}$, so η is constant.

If the reaction is in balance:

$$n_p \times n_n \simeq n_d \times n_\gamma$$

As an order of magnitude: $n_p \sim n_n \sim n_b$

$$\Rightarrow \frac{n_d}{n_b} \sim \frac{n_b}{n_\gamma} \sim \eta \ll 1$$

so as long as many photons can participate to the reaction destroying ${}^2\text{H}$, deuterium cannot accumulate

To overcome this deuterium bottleneck we need to wait that

$$\frac{n_\gamma (E \geq 2.2 \text{ MeV})}{n_\gamma} \lesssim \eta$$

SAHA EQUATION

p, n and d ($\equiv {}^2\text{H}$) are non-relativistic, so:

$$n_i = \int_{s_i} \left(\frac{m_i k_B T}{2\pi \hbar^2} \right)^{3/2} e^{-\frac{m_i c^2 - \mu_i}{k_B T}} \quad i = p, n, d$$

$$n_p + n_n = n_d, \quad \mu_\gamma = 0$$

$$N_{sm} = N_{sp} = N_{sr} = 2, \quad N_{sd} = 4$$

$$m_p c^2 + m_n c^2 = m_d c^2 + B_d, \quad B_d = 2.2 \text{ MeV}$$

Then:

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$$\frac{n_d}{n_p n_n} = \left(\frac{M_d}{m_p m_n} \right)^{3/2} \left(\frac{k_B T}{2\pi \hbar^2} \right)^{-3/2} e^{-\frac{B_d}{k_B T}}$$

Using

$$n_\gamma = \frac{2 \zeta(3)}{\pi^2} \left(\frac{k_B T}{\hbar c} \right)^3 \quad \text{and} \quad m_p \sim m_n \sim \frac{1}{2} m_d$$

and calling:

$$R(T) \equiv \frac{m_p m_n}{m_d m_\gamma}$$

We obtain Saha equation:

$$R(T) = \frac{\sqrt{\pi}}{2 \zeta(3)} \left(\frac{m_p c^2}{4 k_B T} \right)^{3/2} e^{-\frac{B_d}{k_B T}}$$

The reaction is at its half when $m_d \sim m_n$:

$$\Rightarrow \frac{m_p m_n}{m_d} \sim m_p \sim m_b$$

$$\Rightarrow m_\gamma R(T) \sim m_b \Rightarrow R(T) \sim \eta$$

such a low value is driven by $\exp(-B_d/k_B T)$

$$\Rightarrow T_{\text{bottleneck}} \approx 100 \text{ keV}, \quad t_{\text{bot}} \approx 130 \text{ s}$$

So nuclear reactions have to wait ~ 2 min to start, and are over in ~ 5 min

Clearly $t_{\text{bot}} \sim \tau_\beta$, so some neutrons will have decayed by then.

So the important parameters for BBN are

$$T_{\nu, \text{dec}} \rightarrow N_0$$

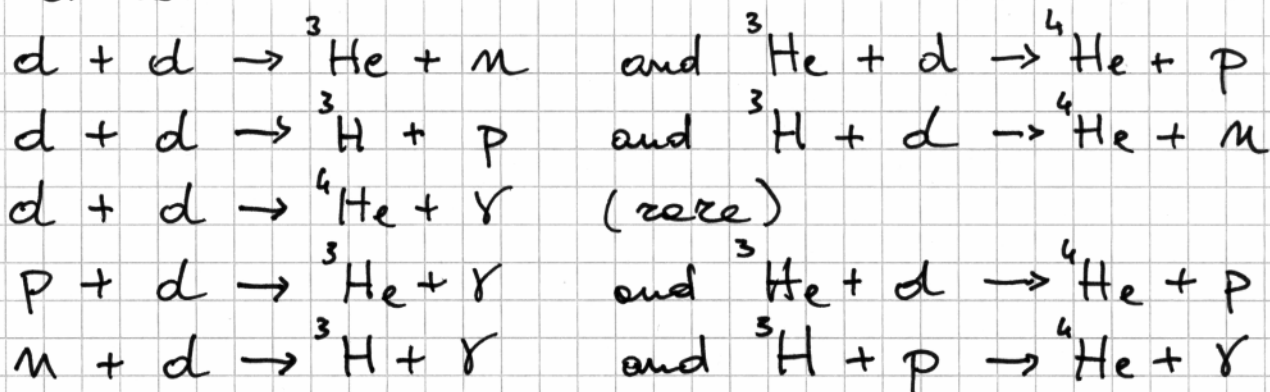
$$\eta \rightarrow \Omega_b h^2$$

FROM ${}^2\text{H}$ TO ${}^4\text{He}$

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${}^4\text{He}$ is the most stable nucleus, so most neutrons will end up in ${}^4\text{He}$ nuclei

Reactions:



The resulting abundance of ${}^4\text{He}$ depends only weakly on η , but depends on N_b

$\Rightarrow N_b = 3$, well before CERN results

The abundance of intermediate products, d and ${}^3\text{He}$, depends strongly on η

$\Rightarrow \Omega_b h^2 \approx 0.022$, in agreement with CMB

But $\Omega_m \approx 0.3$ and $h \approx 0.7$, so $\Omega_b < \Omega_m$

This is one of the strongest pieces of evidence of non-baryonic dark matter

This thermal history was laid down even before detecting the CMB, see the paper by Dicke, Peebles, Roll and Wilkinson of 1963, linked in the web page

END OF BBN

Reactions stop at

$$T \approx 50 \text{ keV}$$

$$t \approx 600 \text{ s}$$

$$ct \approx 2 \times 10^{13} \text{ cm} \approx 2 \text{ AU}$$

$$a \approx 4.7 \times 10^{-9}$$

$$\text{acto} \approx 6.1 \times 10^{19} \approx 20 \text{ pc}$$