



Cosmology I

Second part: FRW models

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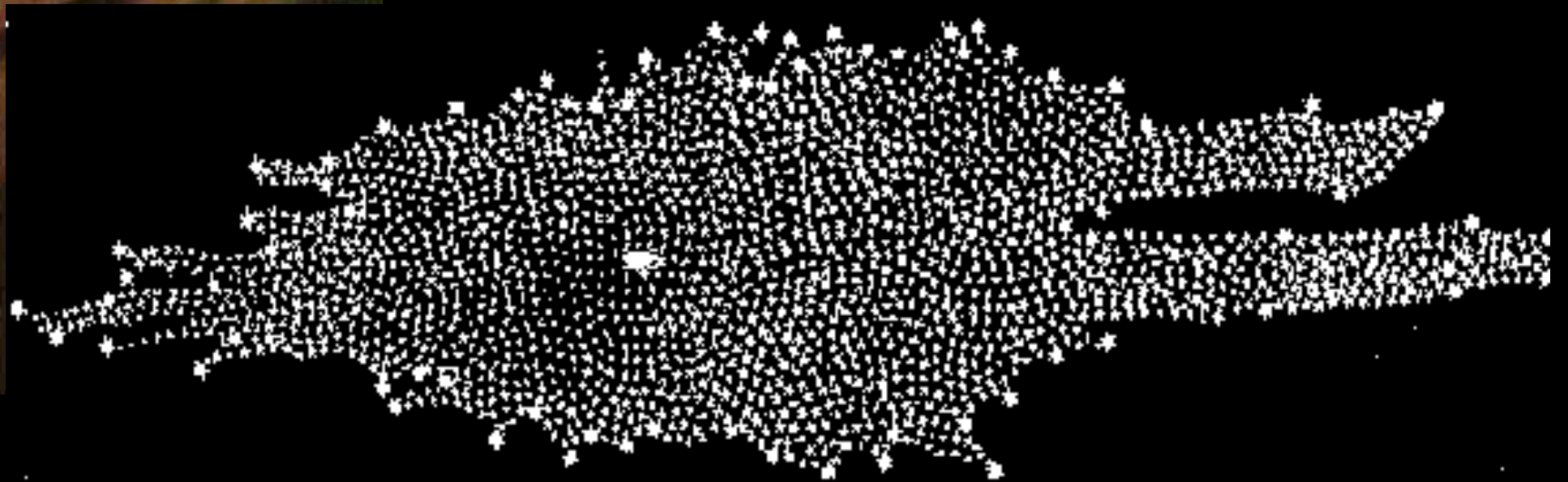
Epistemological facts:

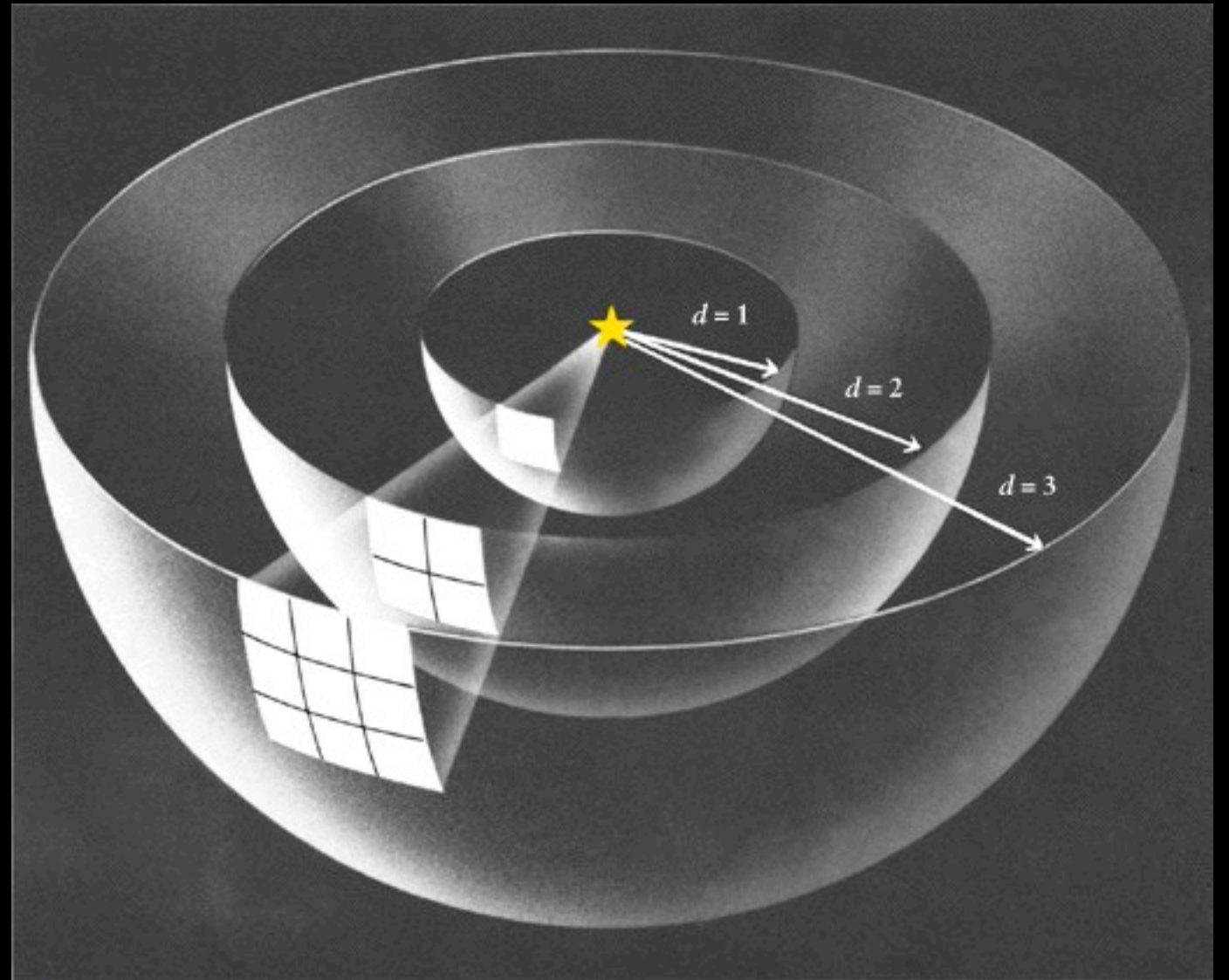
- ★ The Universe is everything we can observe or experiment.
- ★ But we cannot observe to infinity, we have a horizon.
- ★ Under the hypothesis of cosmic homogeneity, we can follow the evolution of the Universe in time.
- ★ We need to take profit of all evidence.



William Herschel
1738-1822

1785







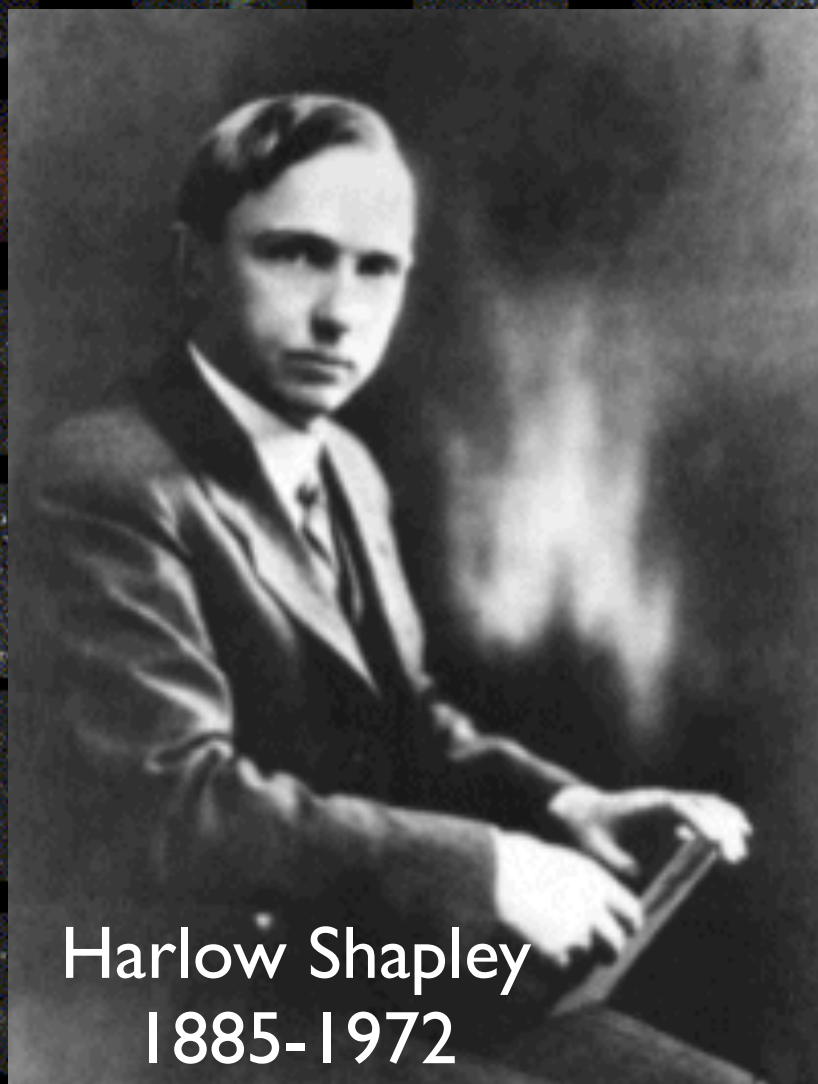
Harlow Shapley
1885-1972

1916

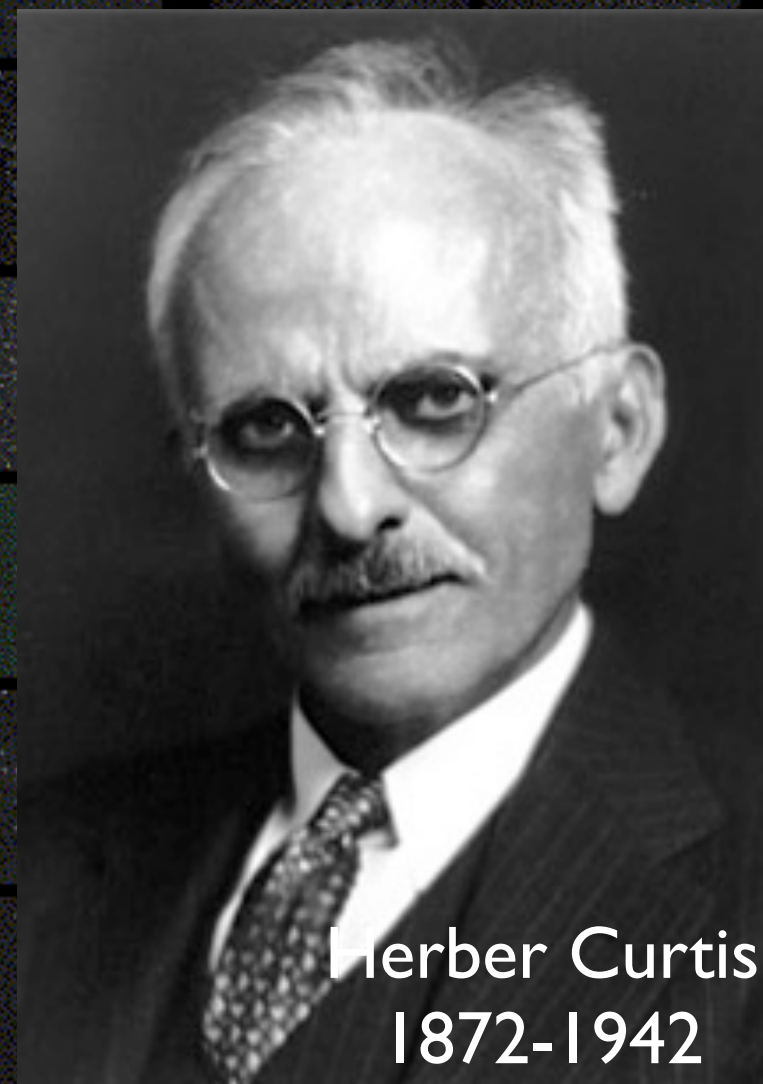


The great Shapley-Curtis debate

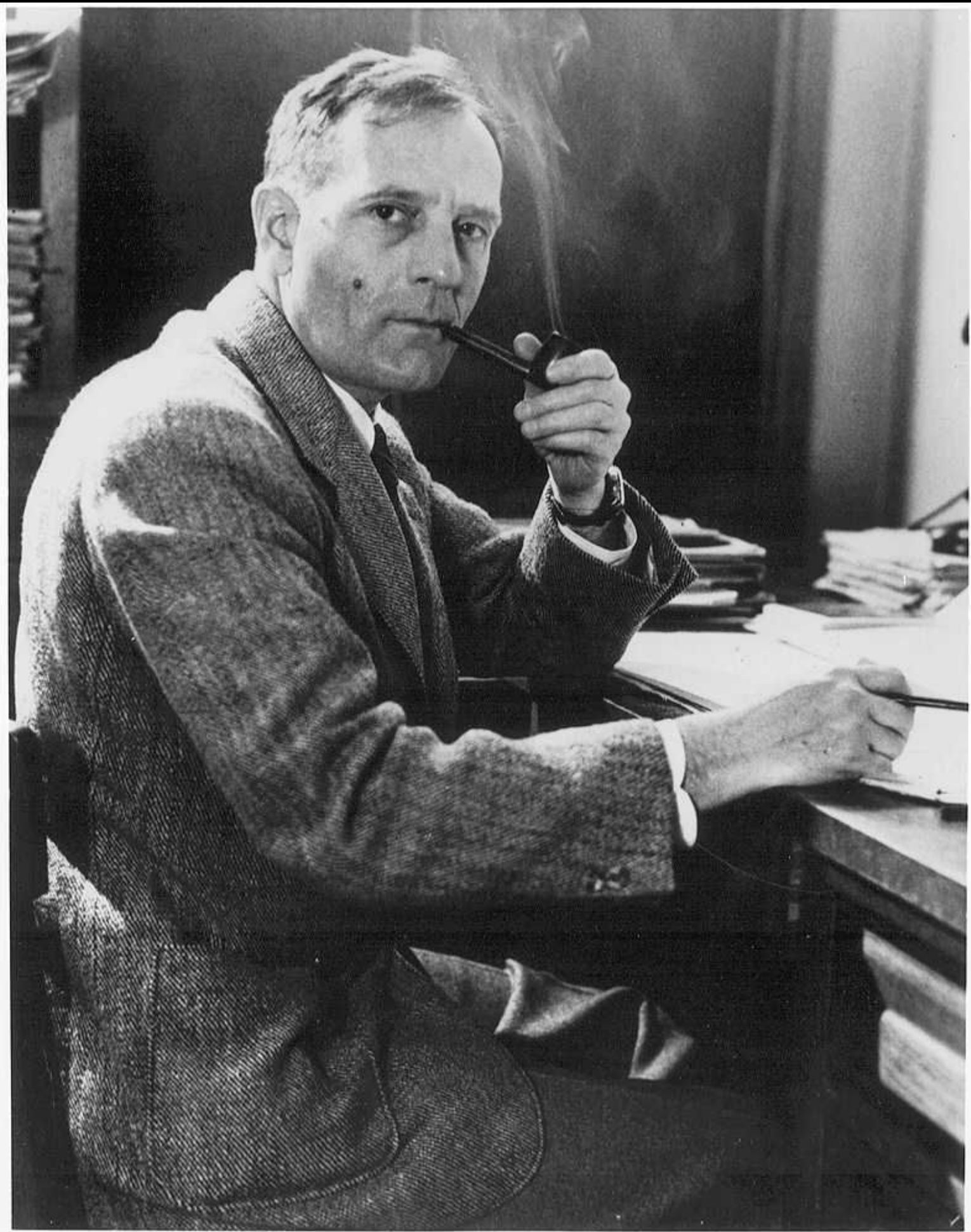
26 April 1920



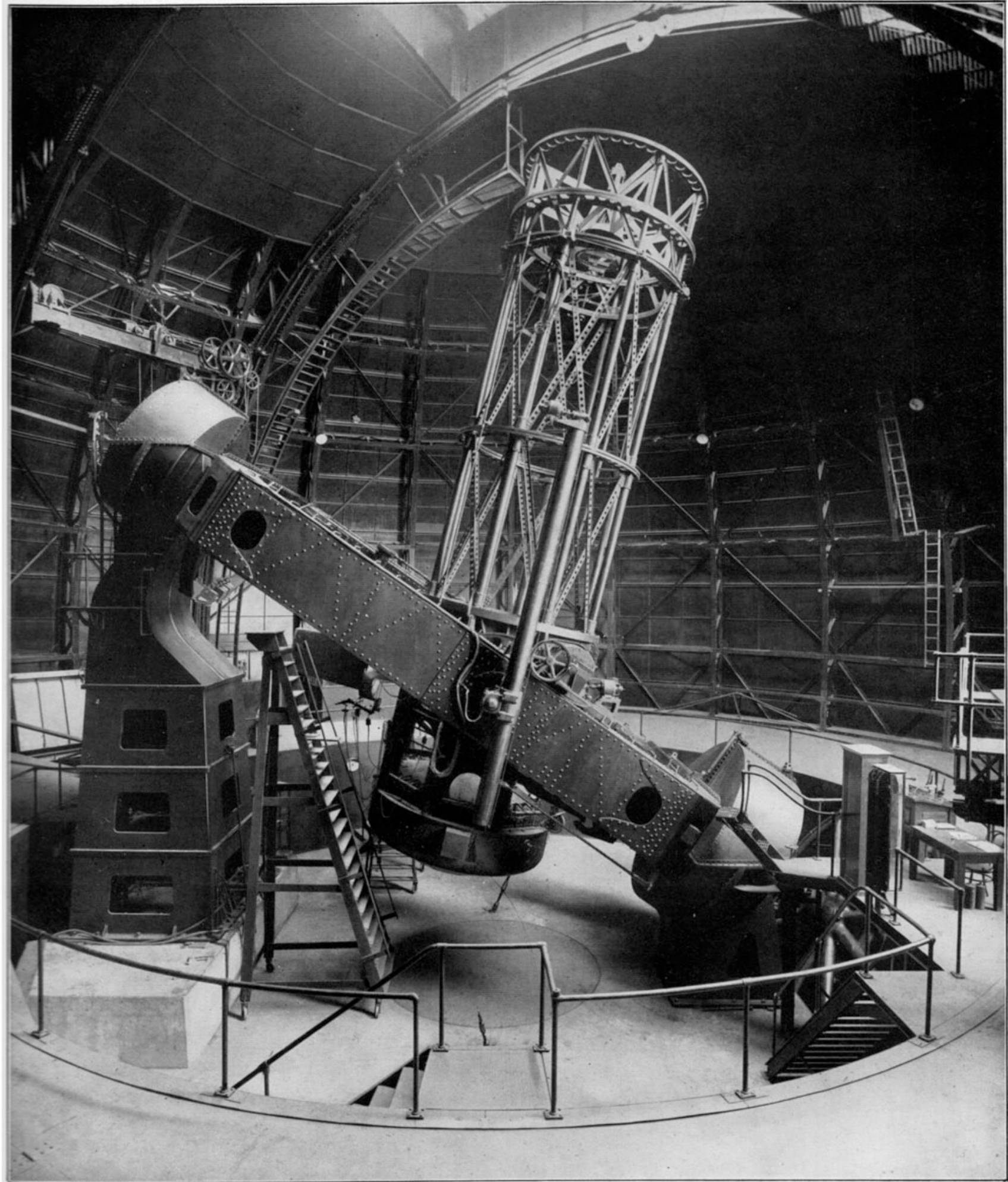
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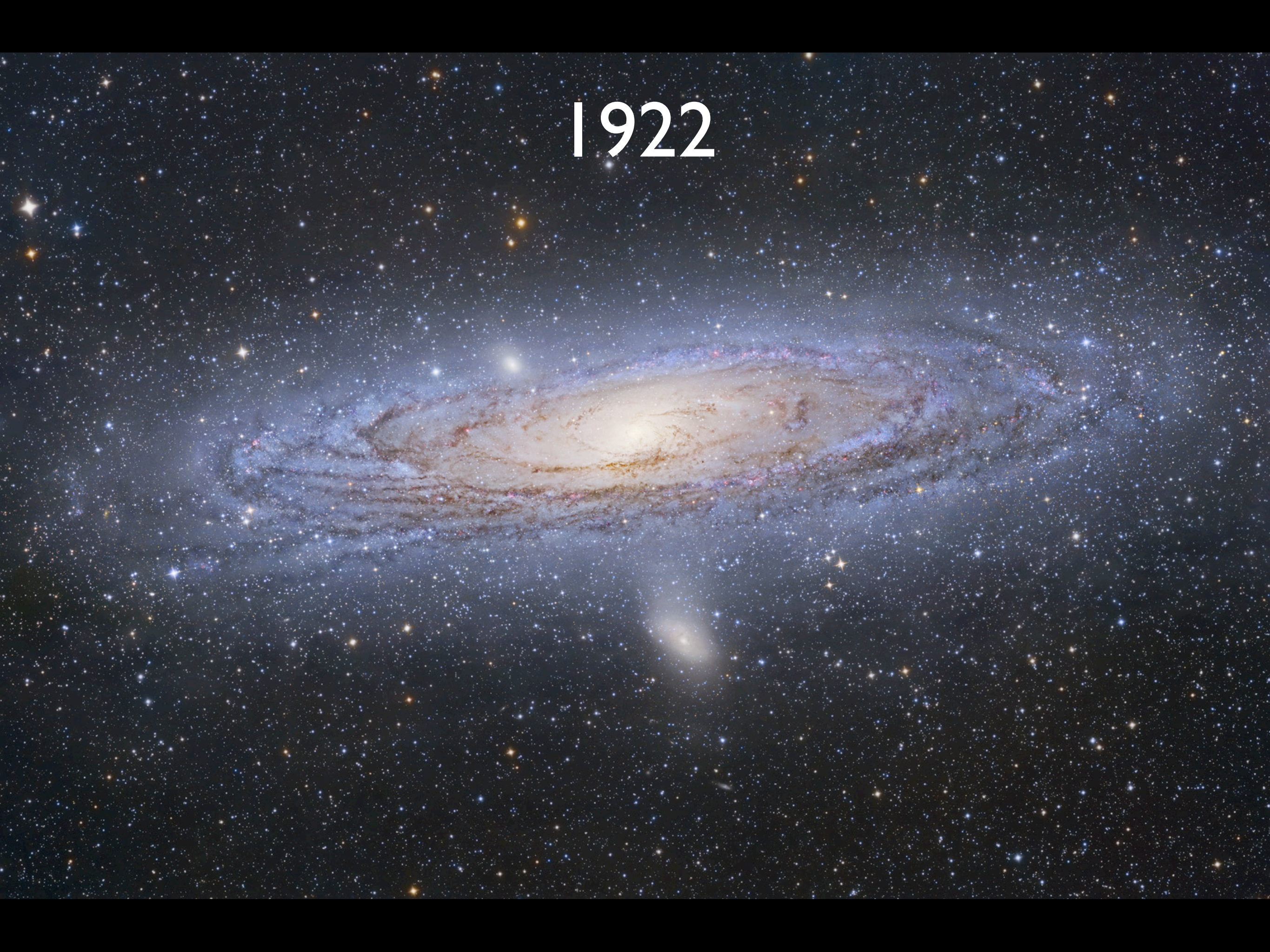
Herber Curtis
1872-1942



Edwin Hubble
1889-1953

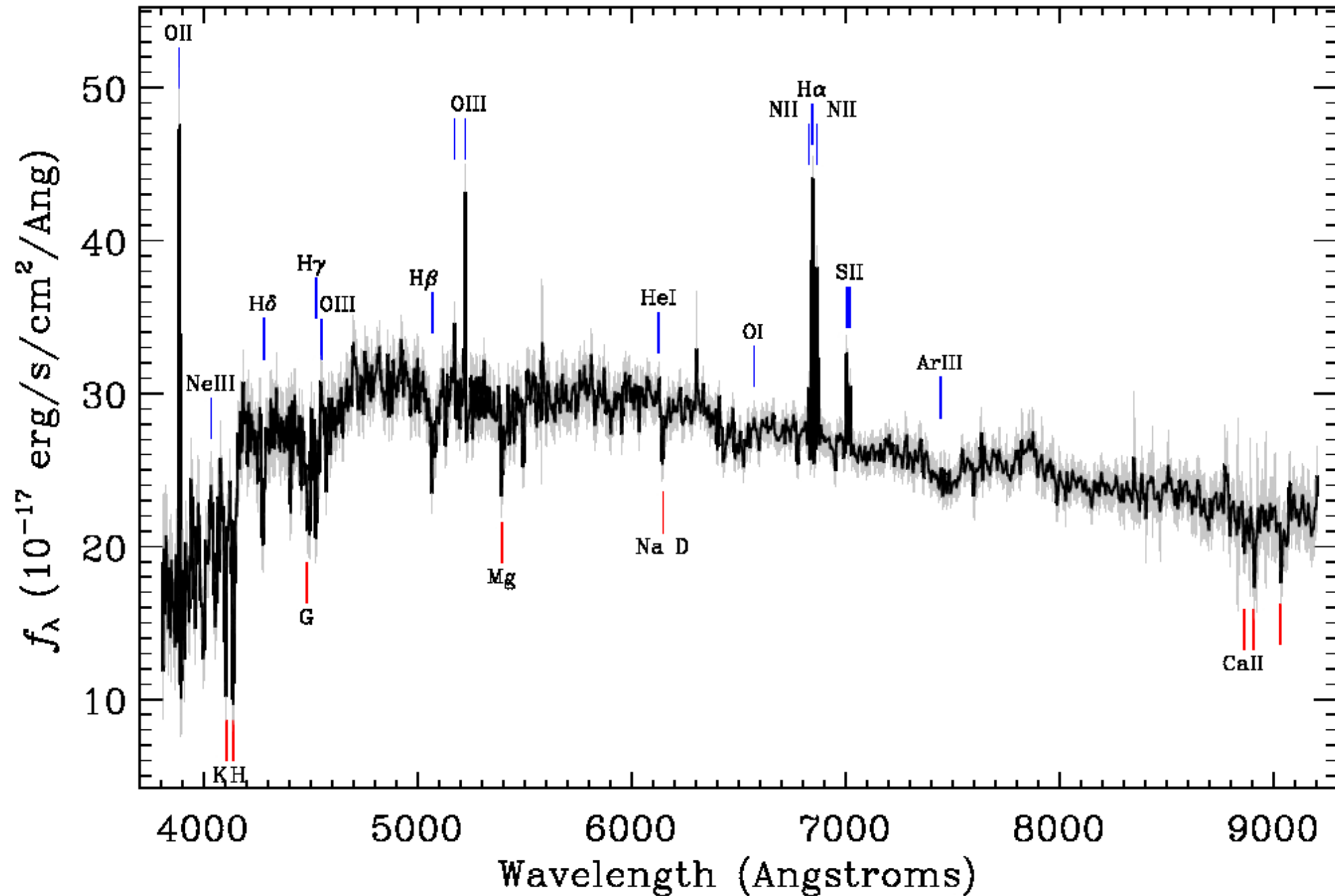


1922



Line-of-sight velocity can be measured with Doppler shifts

Survey: *sdss* Program: *legacy* Target: *GALAXY ROSAT_D ROSAT_E*
RA=25.65806, Dec=-1.22998, Plate=401, Fiber=125, MJD=51788
 $z=0.04263 \pm 0.00002$ Class=GALAXY AGN
No warnings.



1929: Expansion of the Universe, Hubble-Lemaitre law

$$cz = H_0 d$$

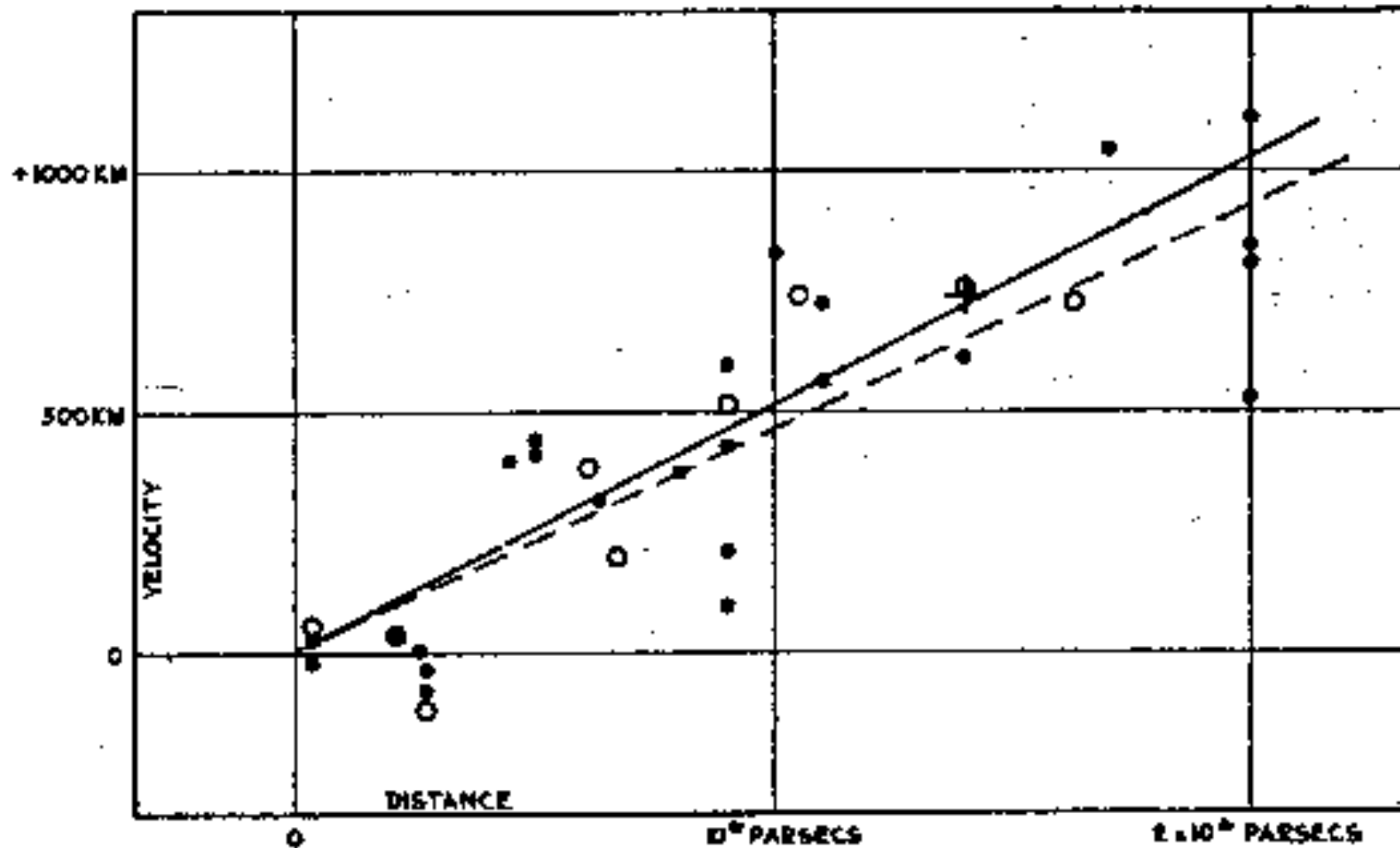
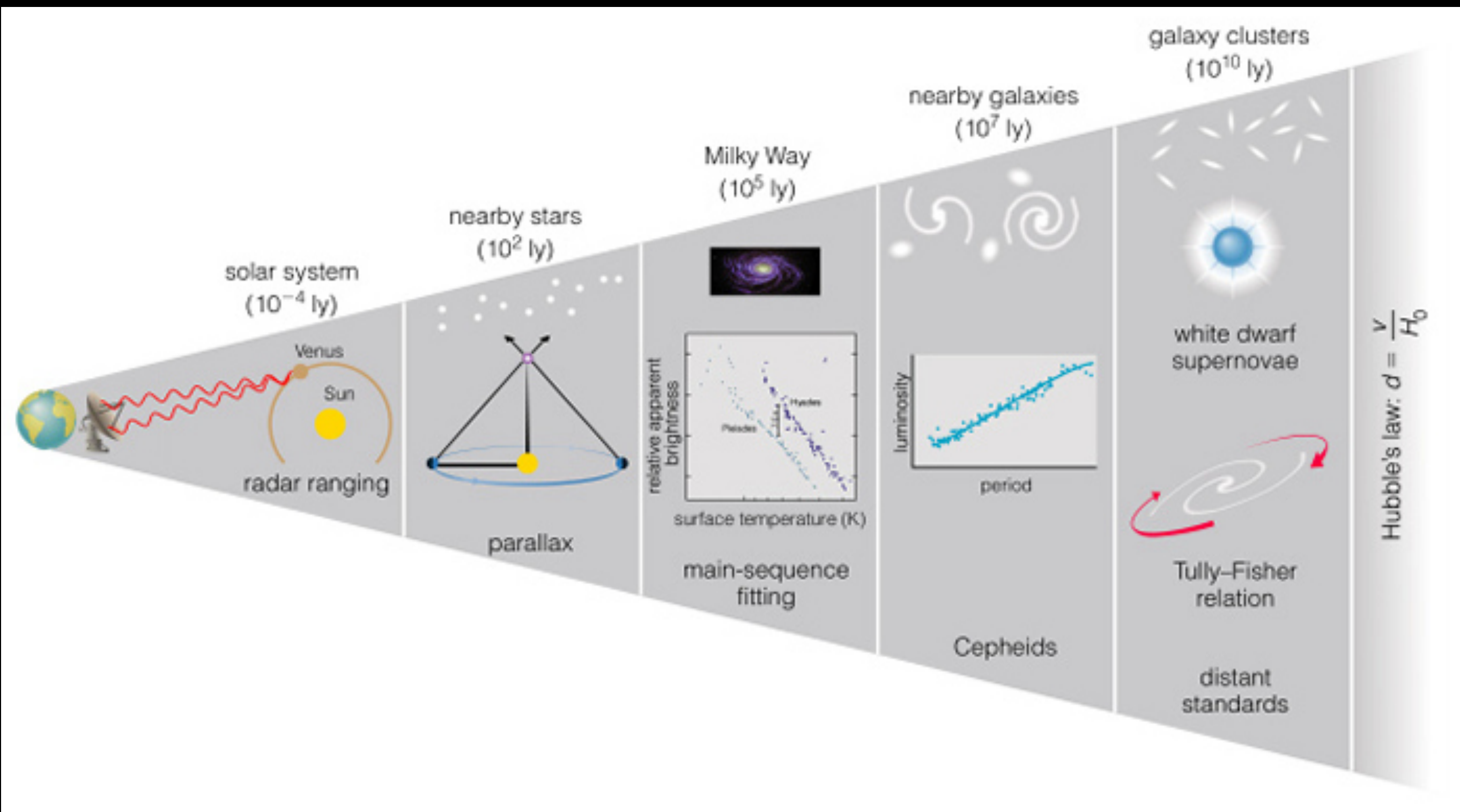
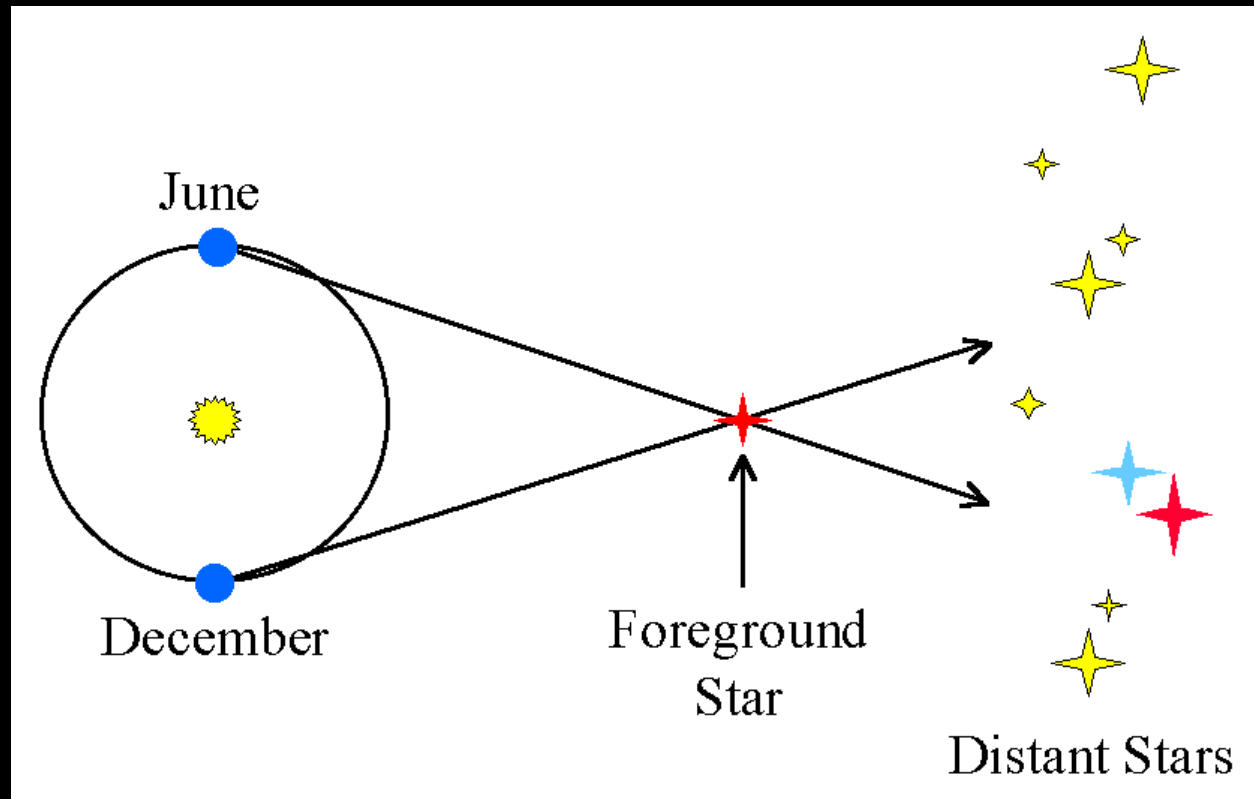


FIGURE 1

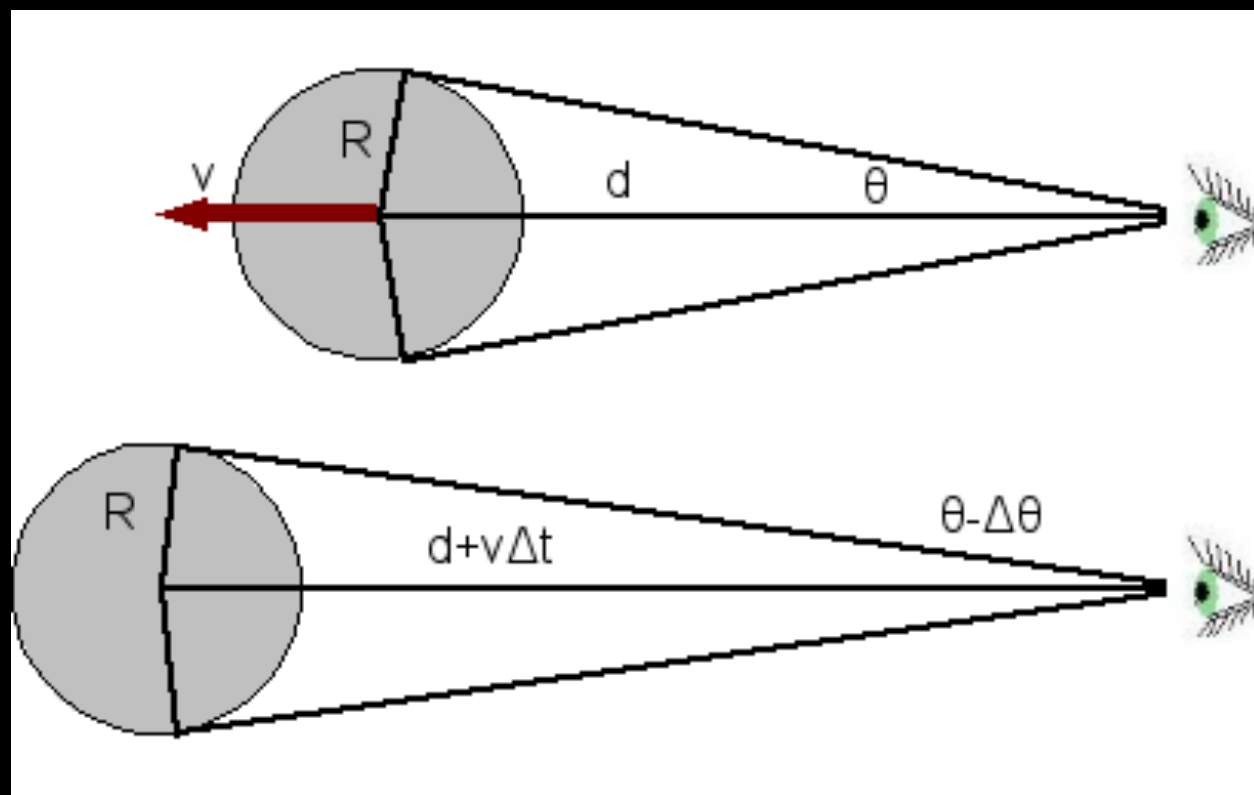
Distances require a “ladder” of measurements



parallax



moving cluster



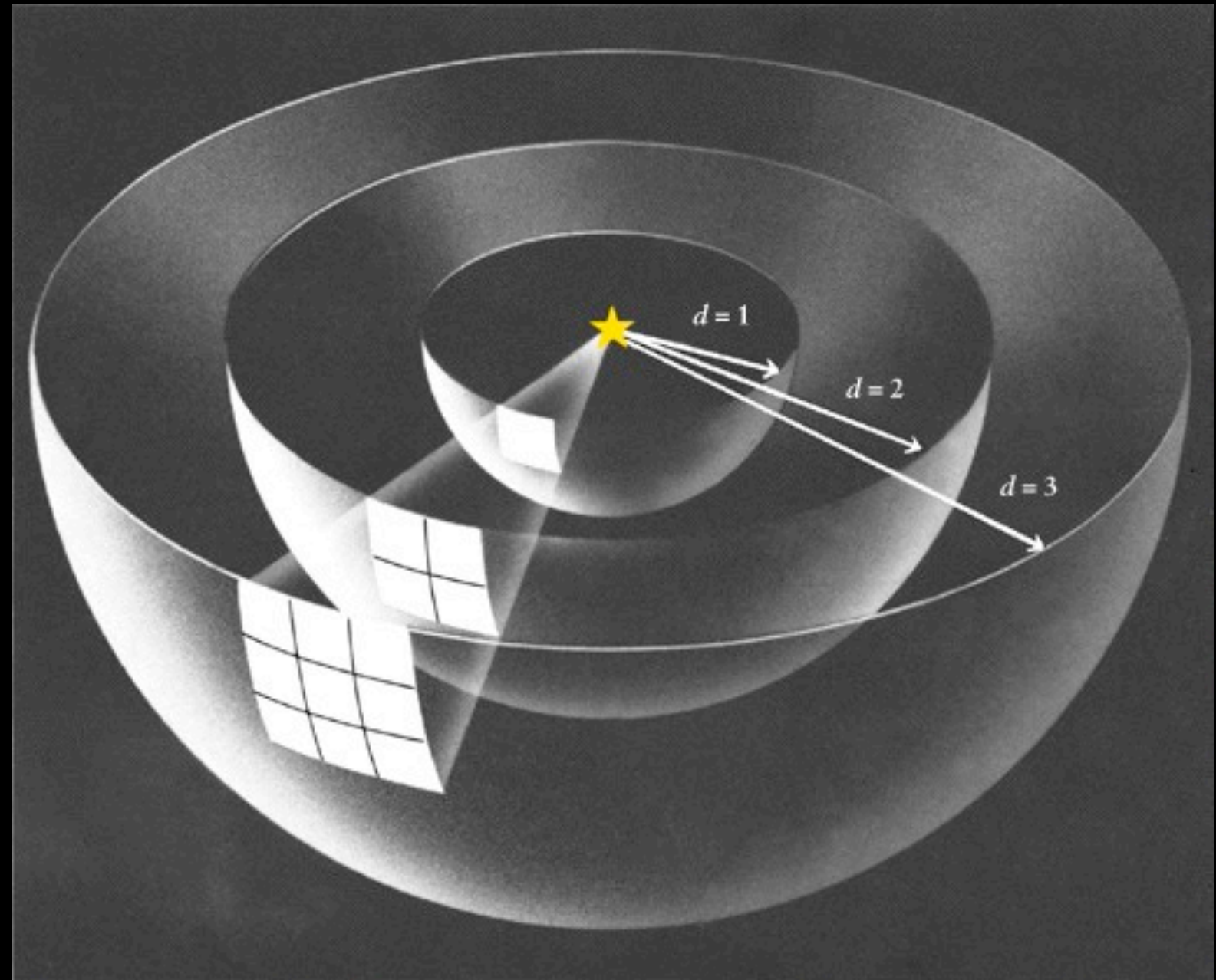
Distance indicators:

standard candle -

$$f = \frac{L}{4\pi d^2}$$

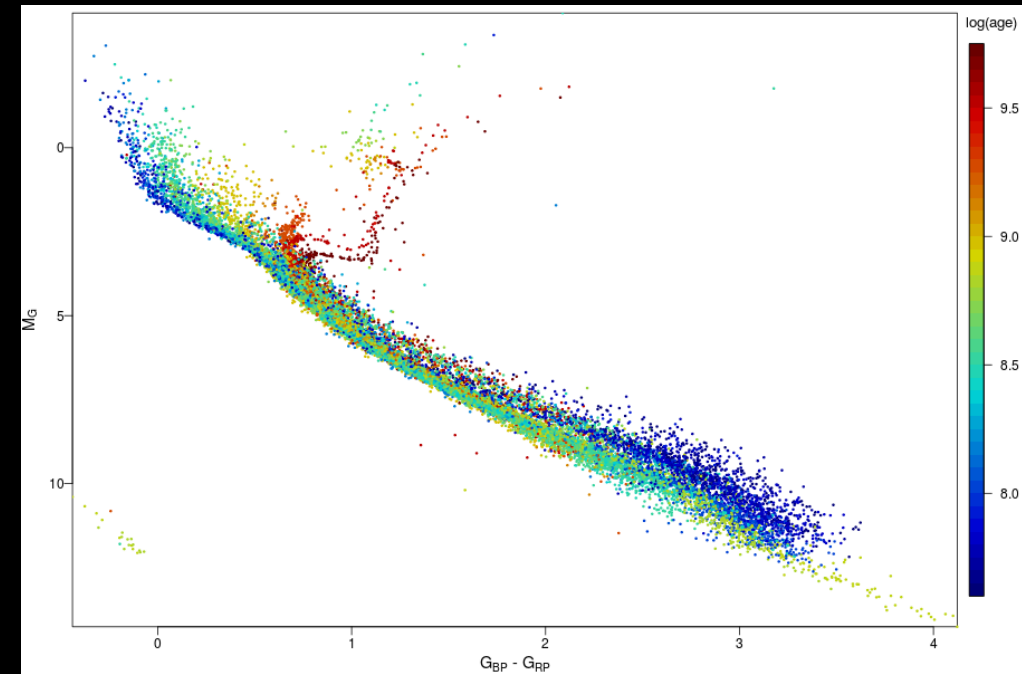
standard ruler -

$$\theta = \frac{D}{d}$$

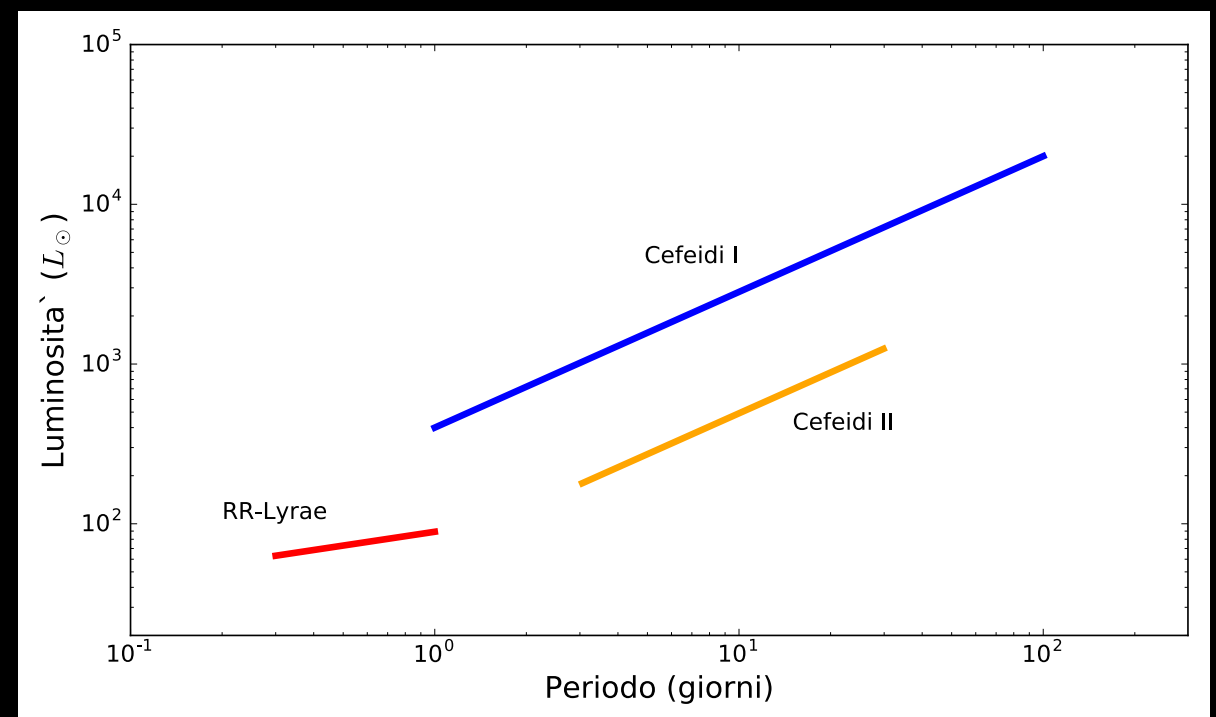


Alternatively, one can use a relation between a distance-independent and a distance-dependent quantity:

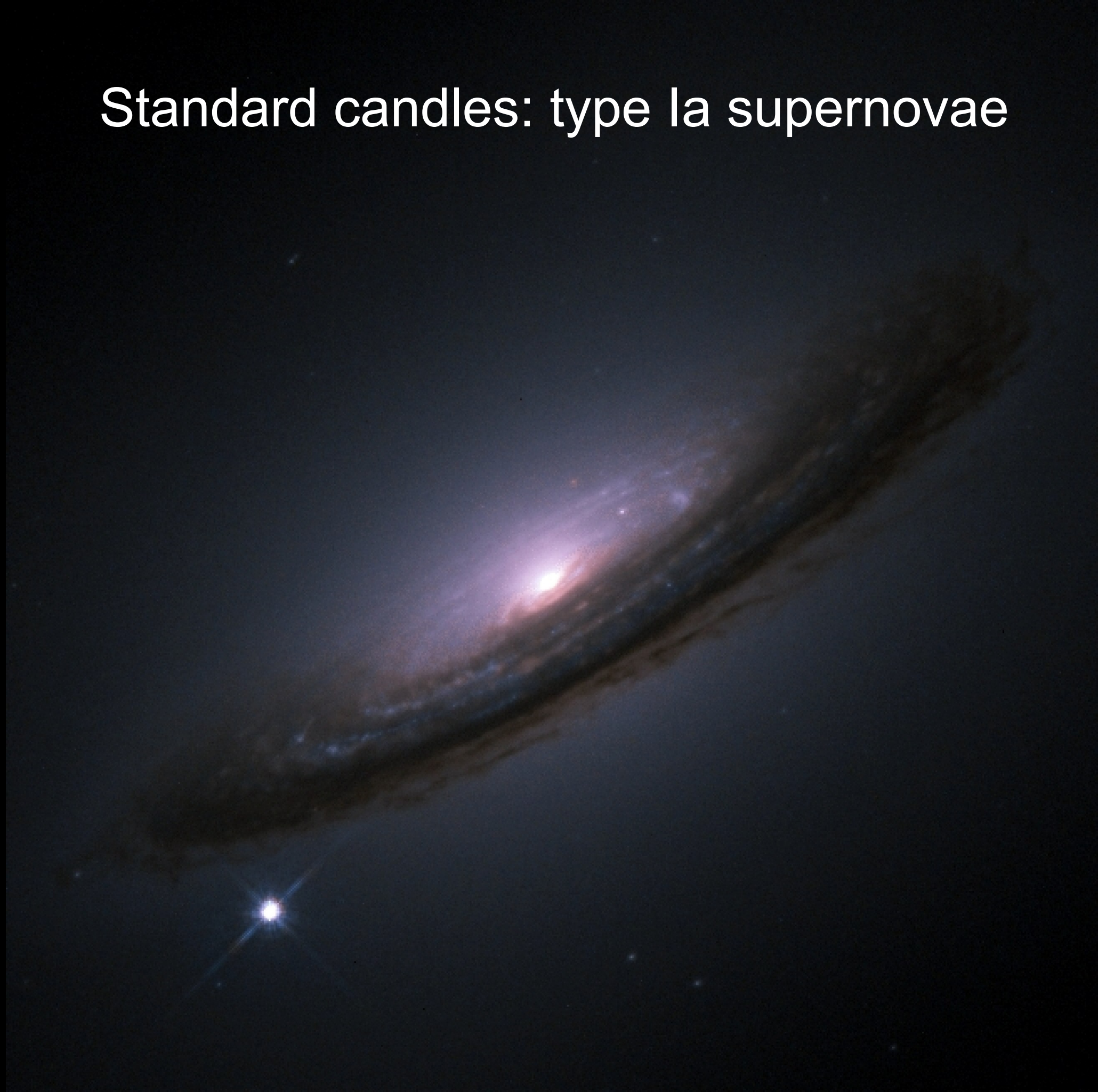
main-sequence fitting:



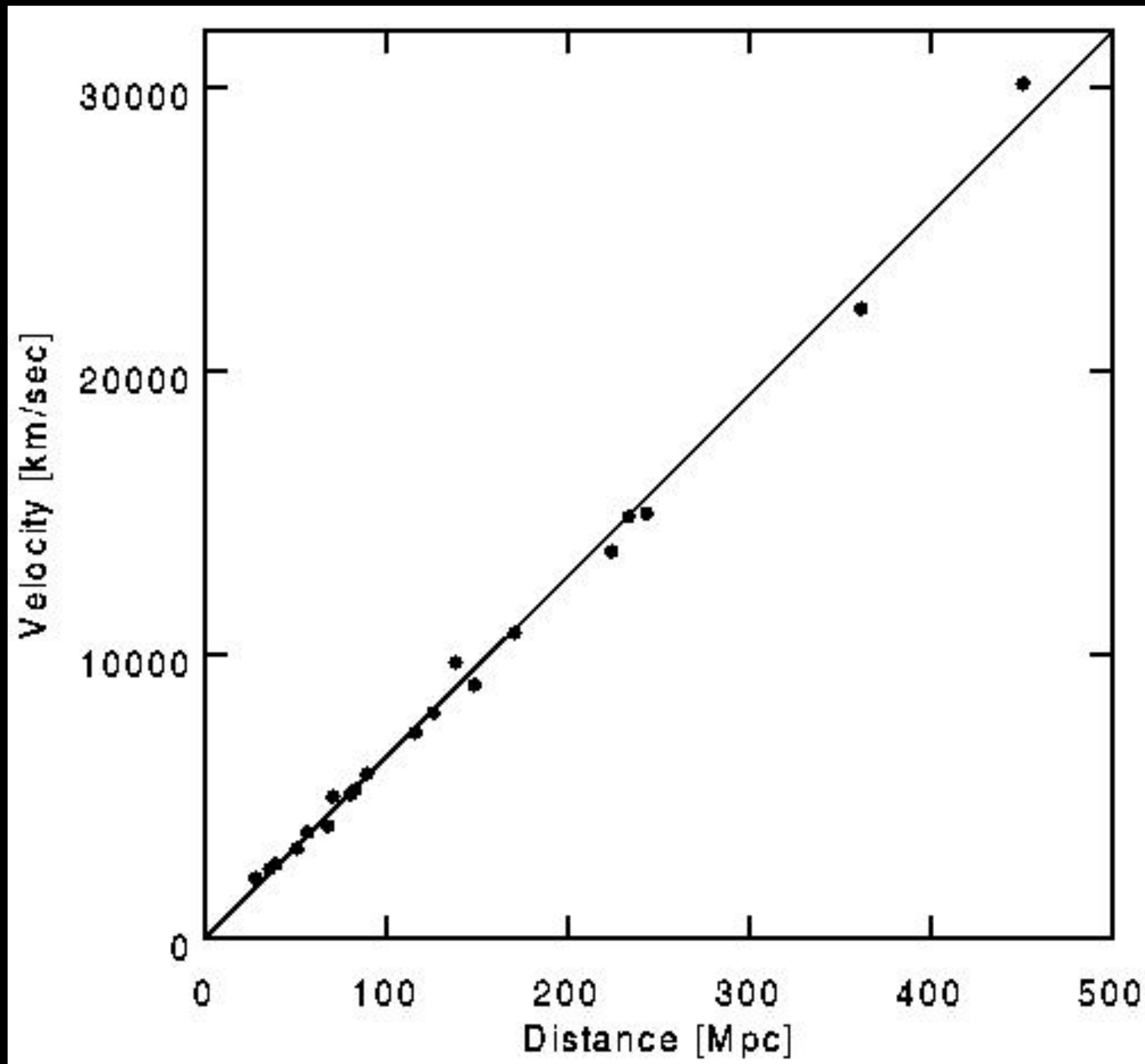
variable stars (cepheids):



Standard candles: type Ia supernovae

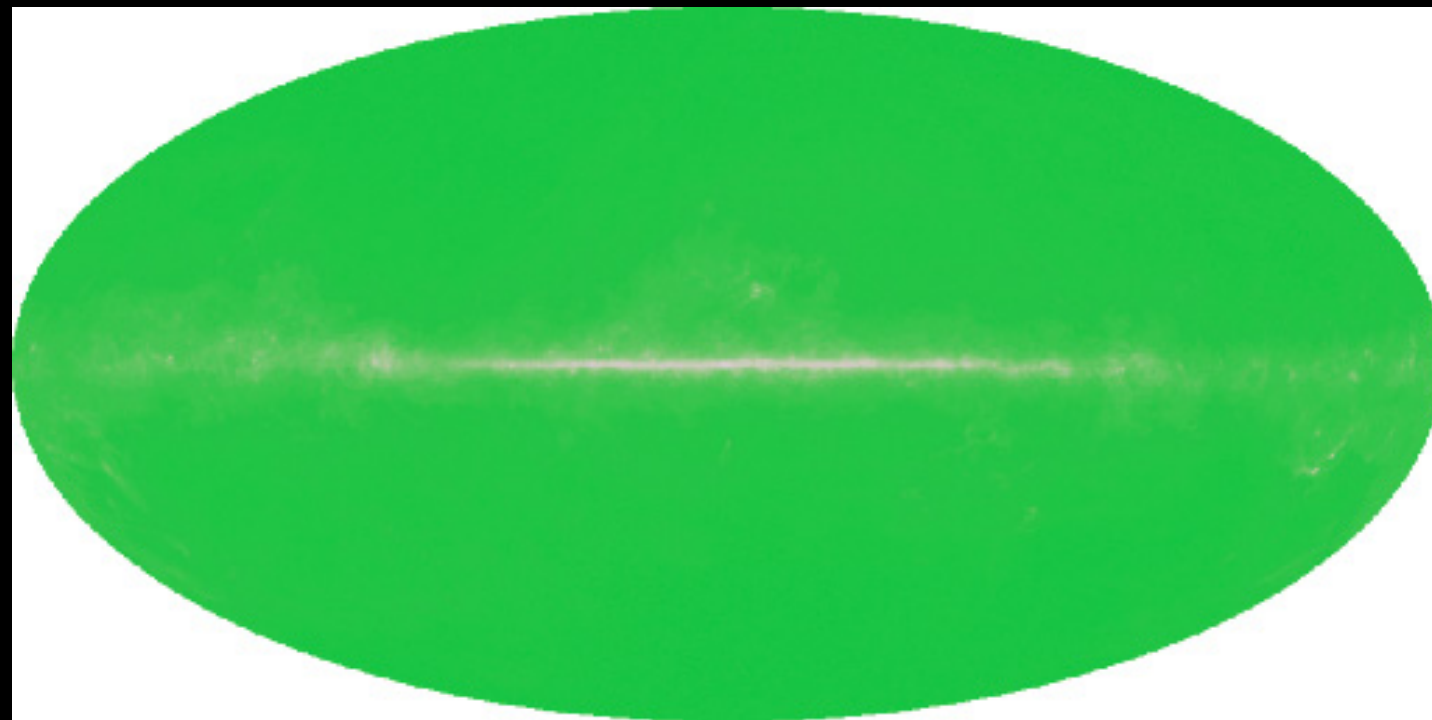


Standard candles: type Ia supernovae



Cosmological principle:

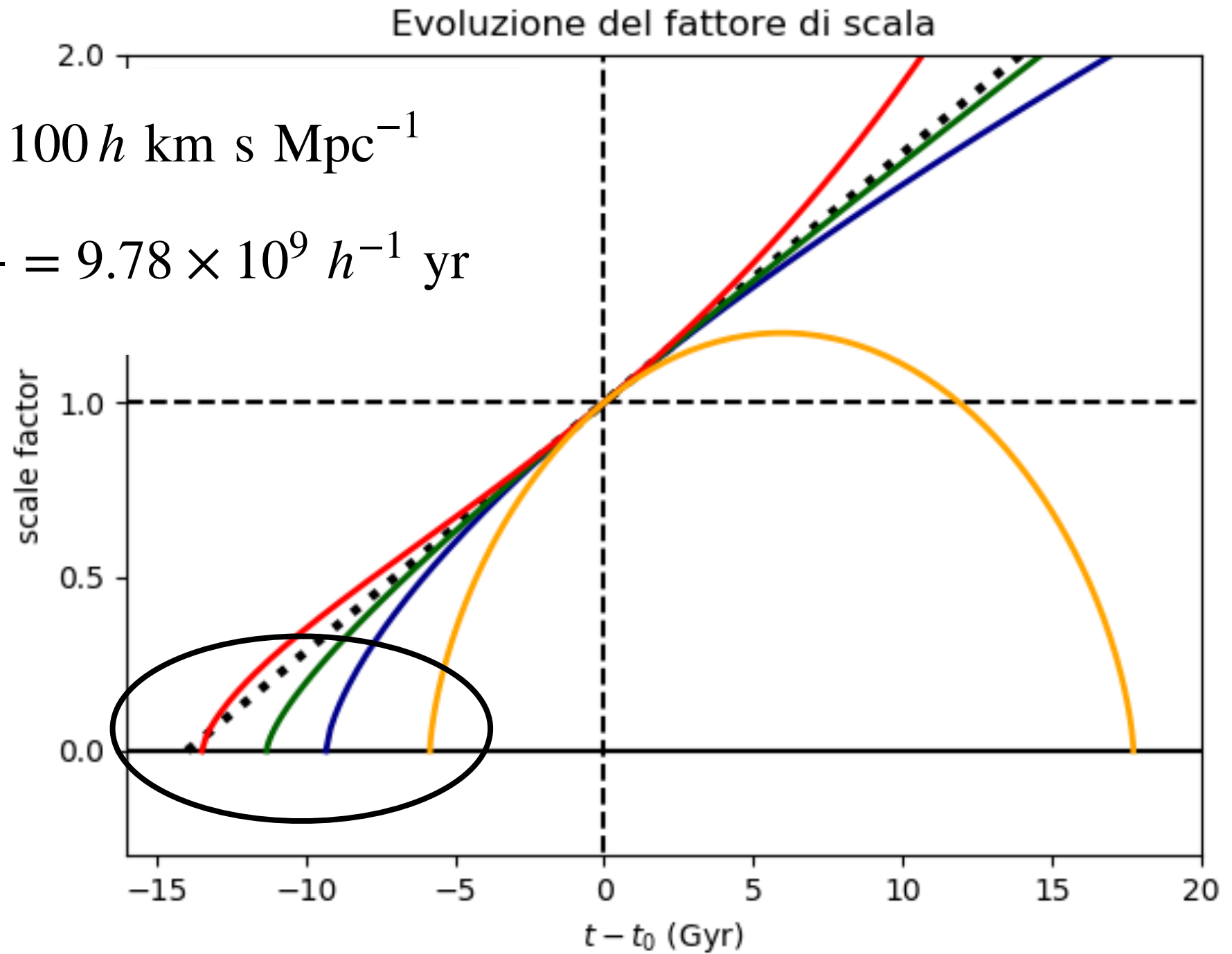
There are no privileged observers in the Universe.
Observed on large scales, the Universe is
homogeneous and **isotropic**.



Observational evidence of isotropy: CMB.
Isotropy around a un-privileged observer implies
homogeneity (if the metric is analytic)

Hubble time

$$H_0 = 100 h \text{ km s Mpc}^{-1}$$
$$t_H = \frac{1}{H_0} = 9.78 \times 10^9 h^{-1} \text{ yr}$$



Big Bang

Hot Big Bang

The universe was much denser and hotter in the distant past, density gets to infinity at $t=0$

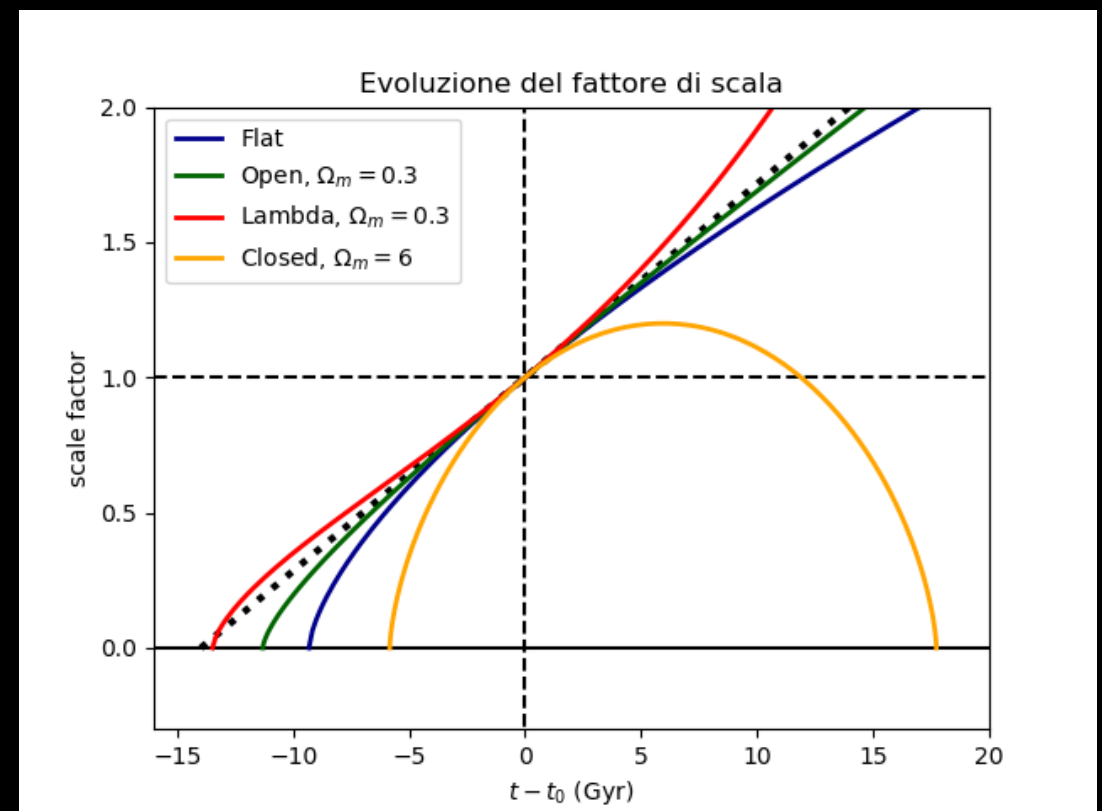
This is a **naked singularity** (no EH)

Hubble law does **not** imply that the observer is at the center

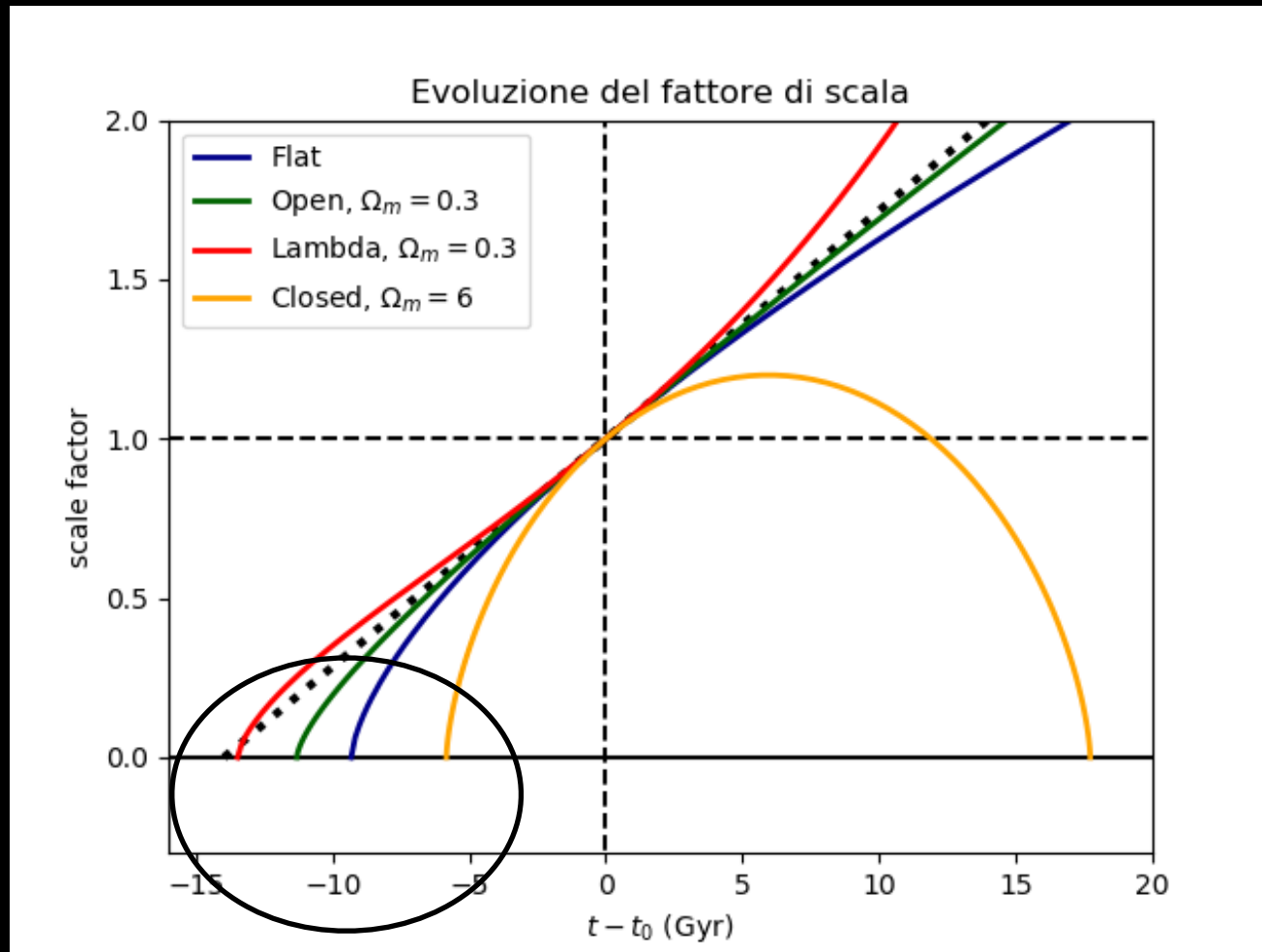
All the space expands, keeping homogeneity and isotropy

If the universe is infinite, it remains infinite at all times $t > 0$

Because we do not understand physics at high energies, we should consider the Big Bang as an **extrapolation** of known physics



A problem with the age of the Universe?



$$H_0 = 100 h \text{ km s Mpc}^{-1}$$

$$t_H = \frac{1}{H_0} = 9.78 \times 10^9 h^{-1} \text{ yr}$$

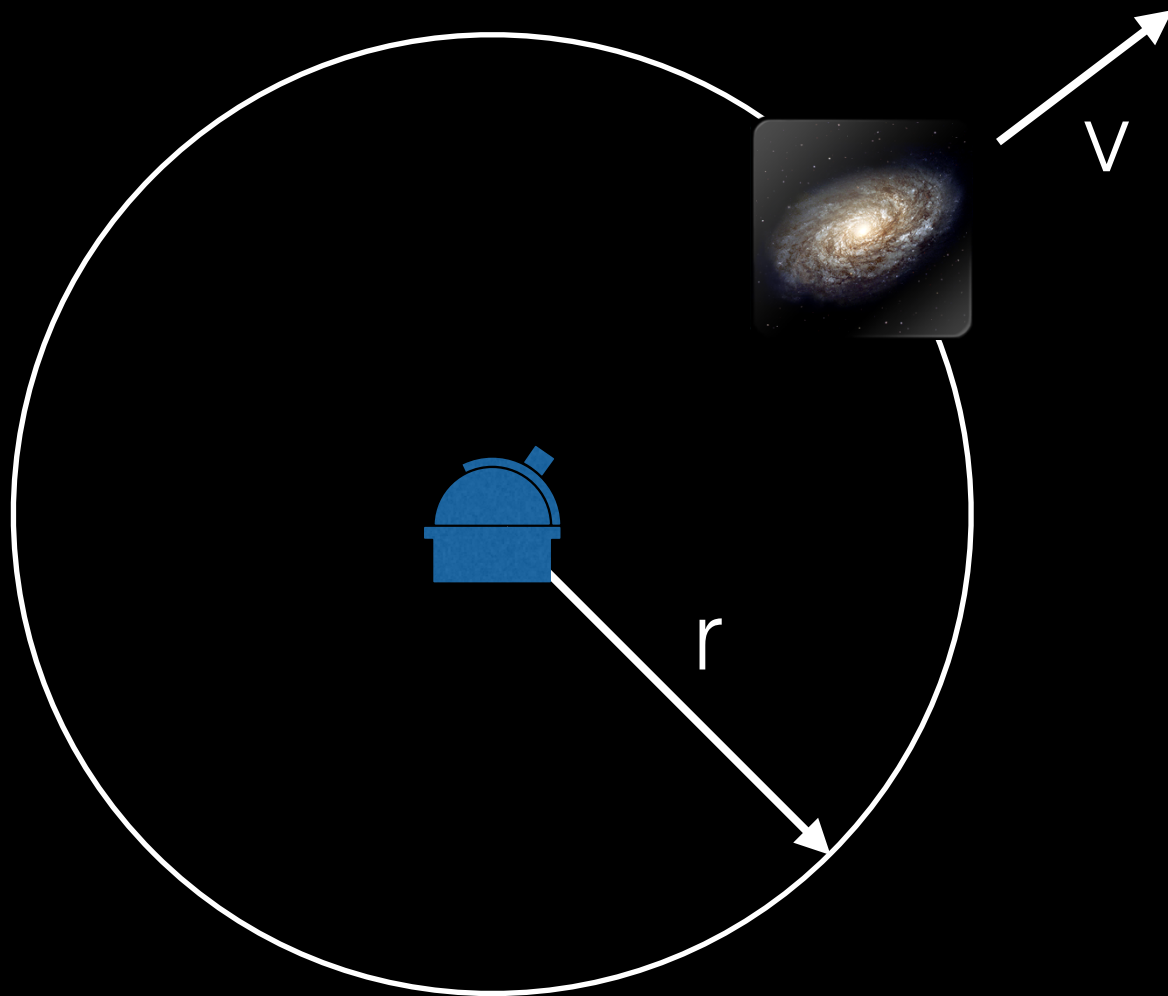
$$t_{GC} \simeq 20 \text{ Gyr}$$



Cosmology with Newtonian gravity

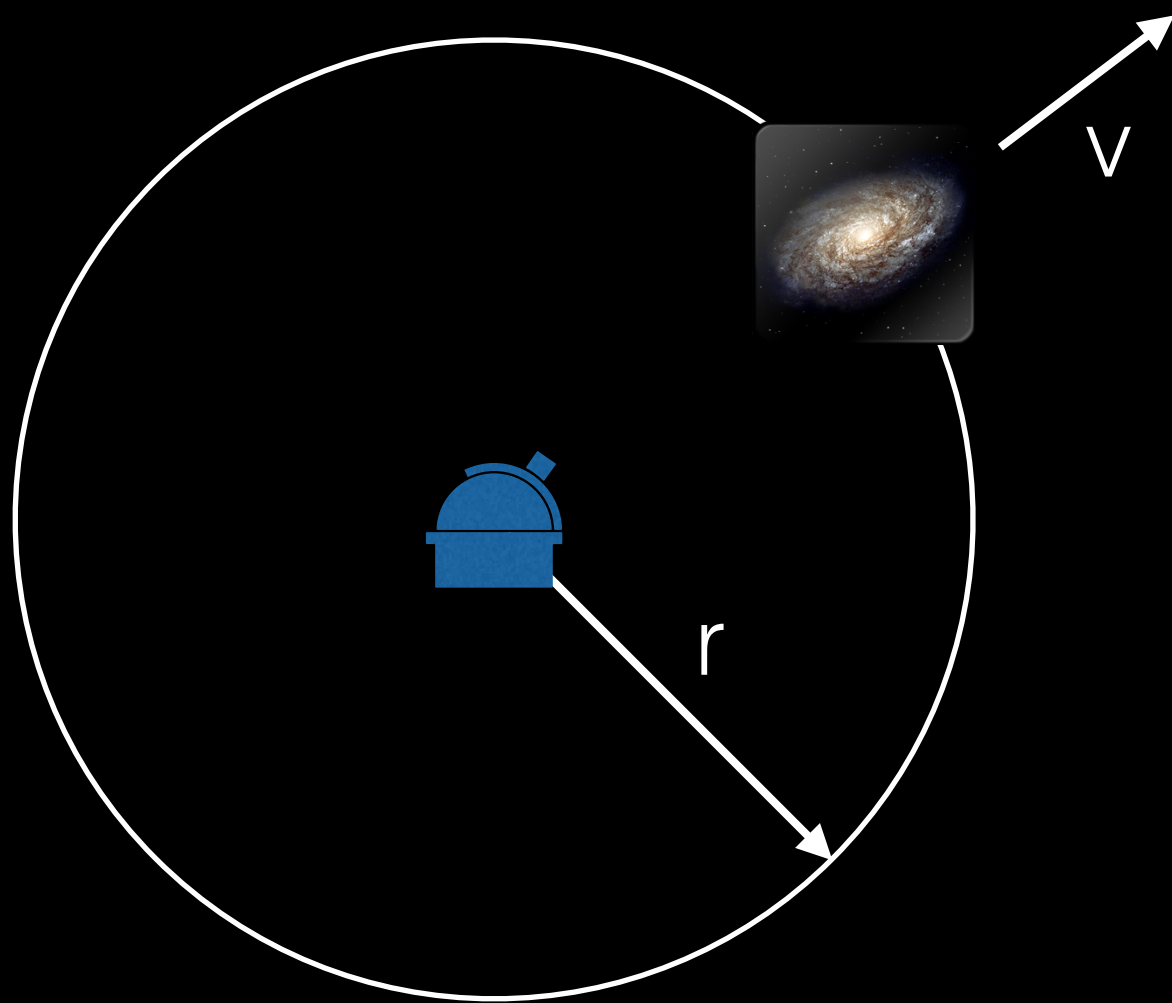
$$\nabla^2\Phi = 4\pi G\rho \quad \Longrightarrow \quad \Phi = \text{const} \times r^2$$

$$\text{cosmological principle} \quad \Longrightarrow \quad \rho = 0$$



Using a result of GR, Birkhoff's theorem, it is possible to demonstrate that in spherical symmetry matter outside a radius r does not influence the metric

Cosmology with Newtonian gravity



$$E = \frac{1}{2}mv^2 - \frac{GmM}{r} = \text{const}$$

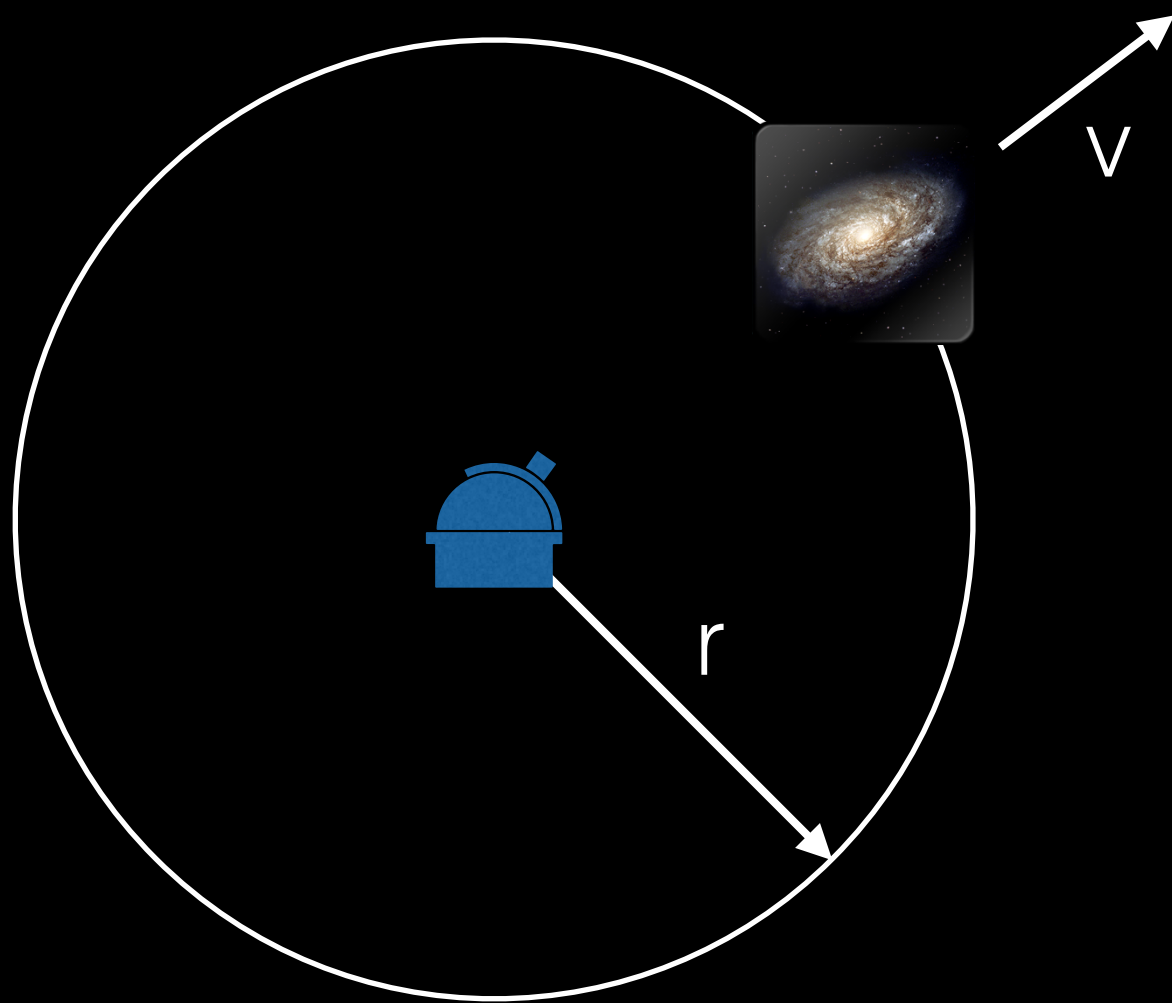
$$v = H_0 r, \quad M = \frac{4\pi}{3}r^3 \rho$$

$$E = 0 \implies \rho = \rho_c = \frac{3H_0^2}{8\pi G}$$

critical density:

$$\rho_c = 2.77 \times 10^{11} h^2 M_\odot \text{Mpc}^{-3} = 1.9 \times 10^{-29} h^2 \text{g cm}^{-3}$$

Cosmology with Newtonian gravity



scale factor : $a(t) = \frac{r}{r_0}$

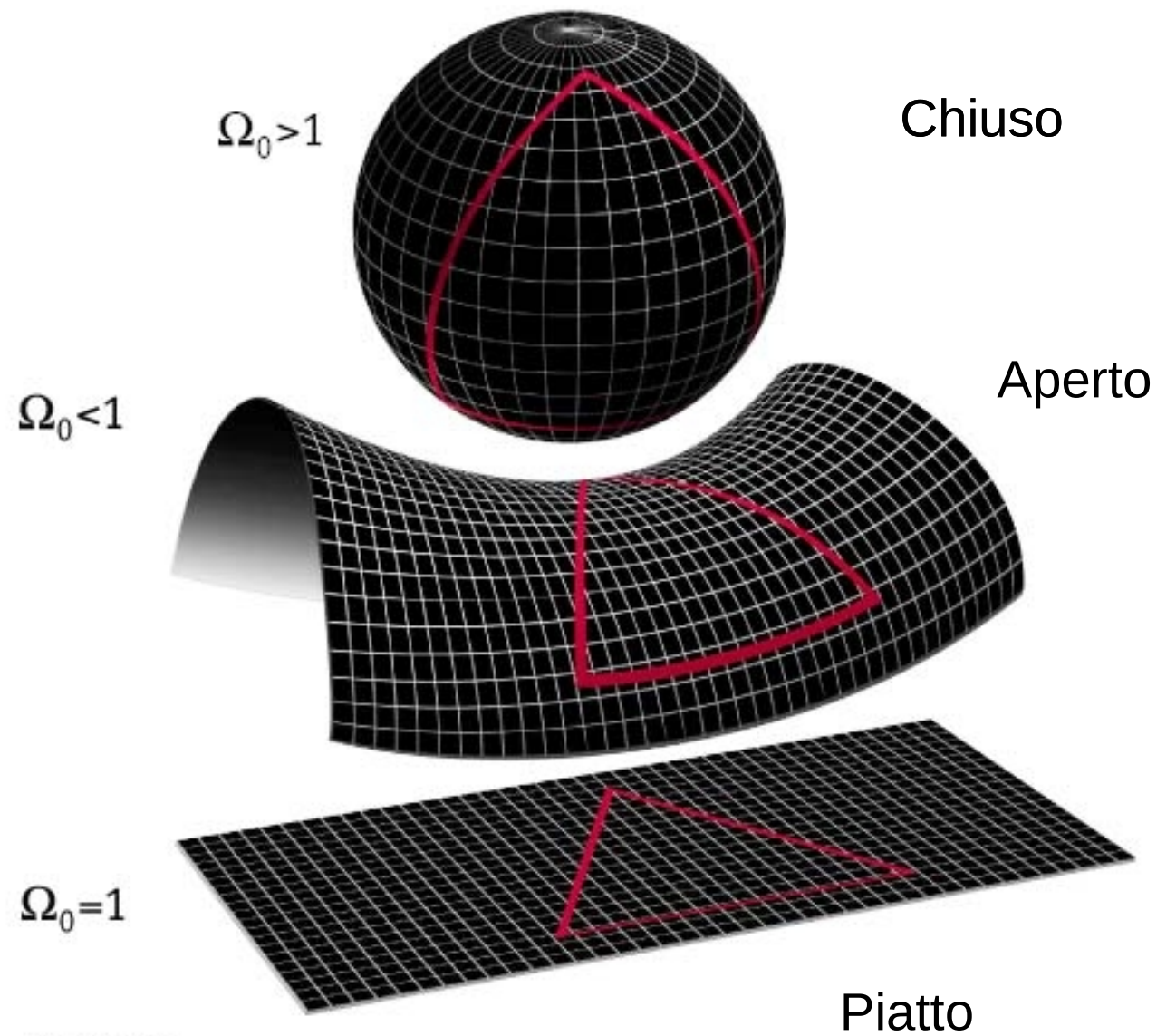
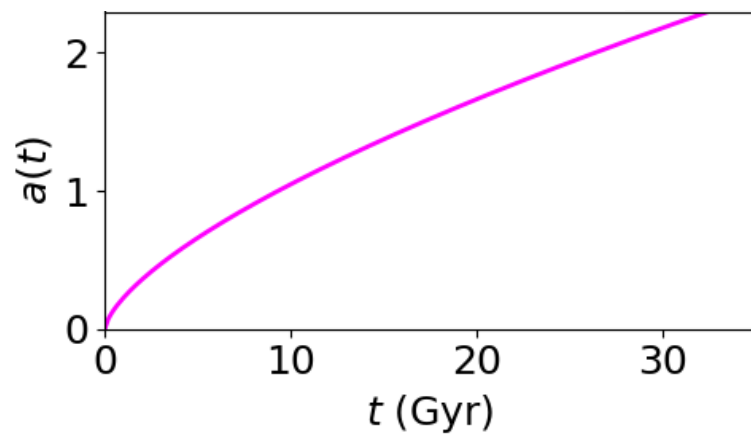
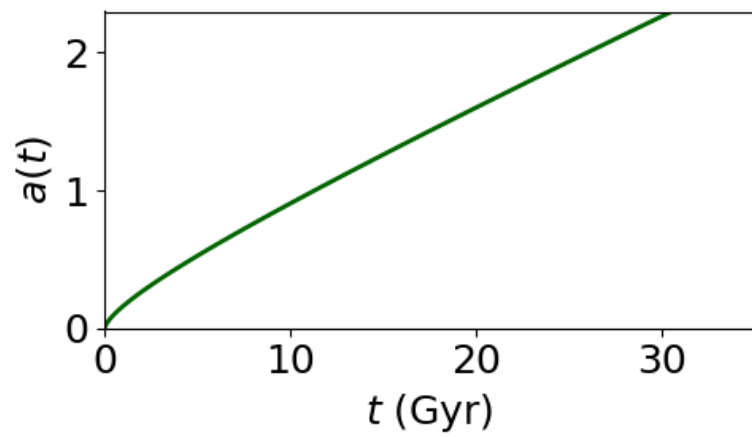
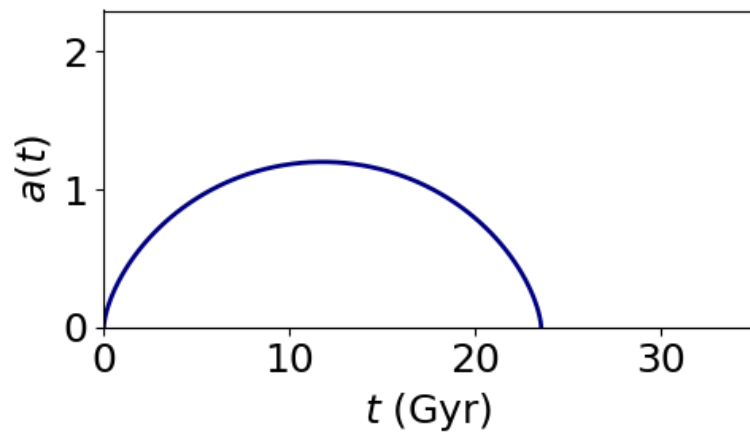
density parameter : $\Omega_0 = \frac{\rho}{\rho_c}$

Hubble parameter : $H = \frac{1}{a} \frac{da}{dt}$

evolution with generic E:

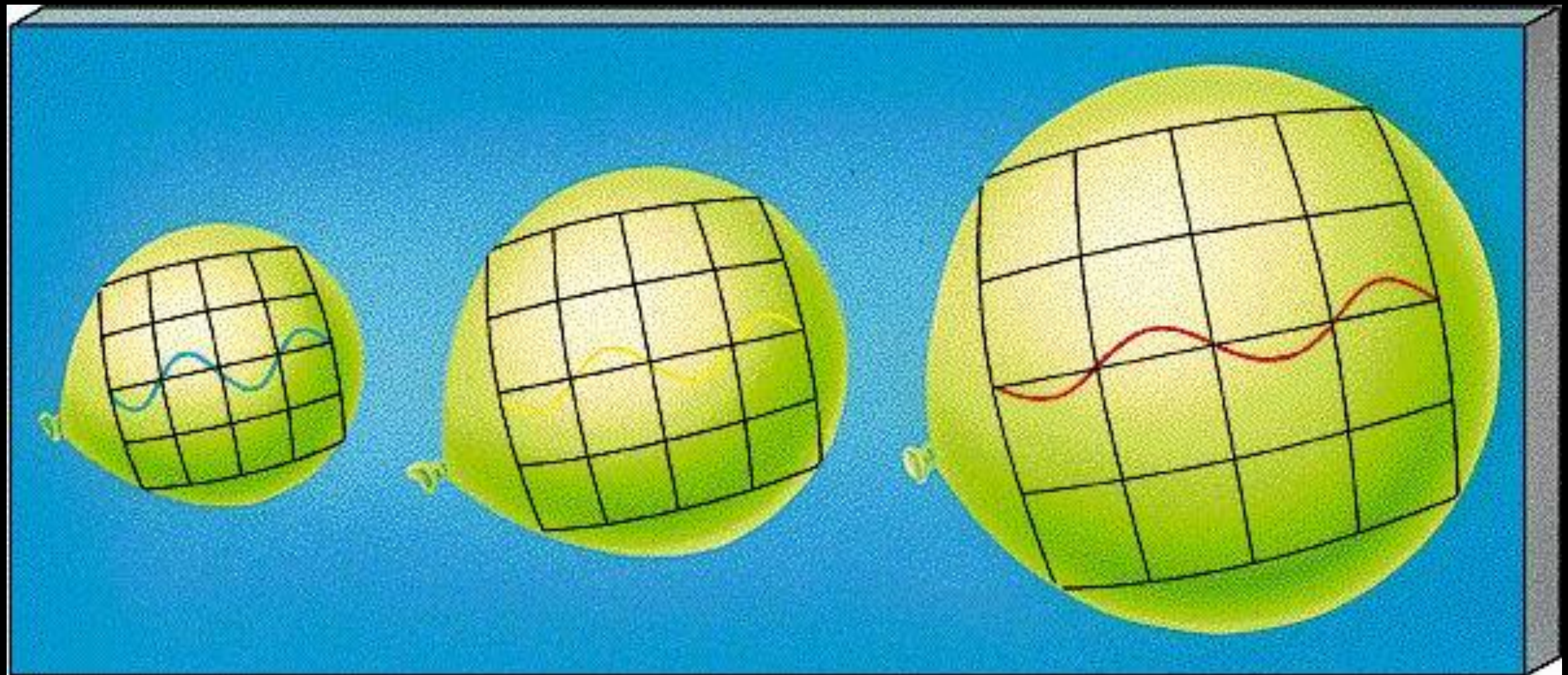
second Friedmann equation $\left(\frac{H}{H_0} \right)^2 = \frac{\Omega_0}{a^3} + \frac{1 - \Omega_0}{a^2}$

Fate of the universe

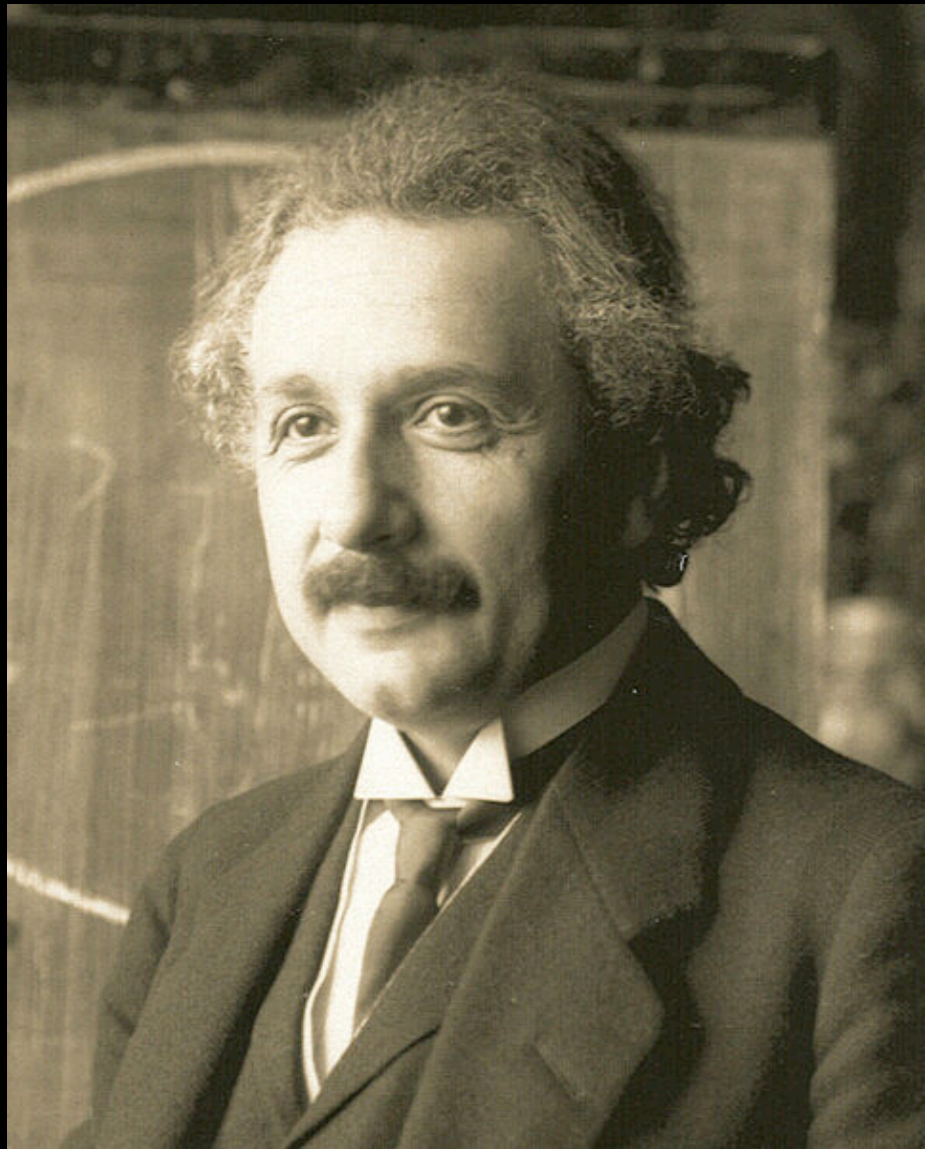


MAP990006

Cosmological redshift



The cosmological constant: 1, Einstein



The universe was thought to be maximally symmetric: homogeneous, isotropic, static

$$1915: \quad R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi GT_{\mu\nu}$$

The only way to obtain an (unstable) static solution is to add a cosmological constant:

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi GT_{\mu\nu}$$

1929: Hubble law, the static solution is ruled out

Lambda: Einsteins' biggest blunder!

The cosmological constant: 2, QFT ('80s)

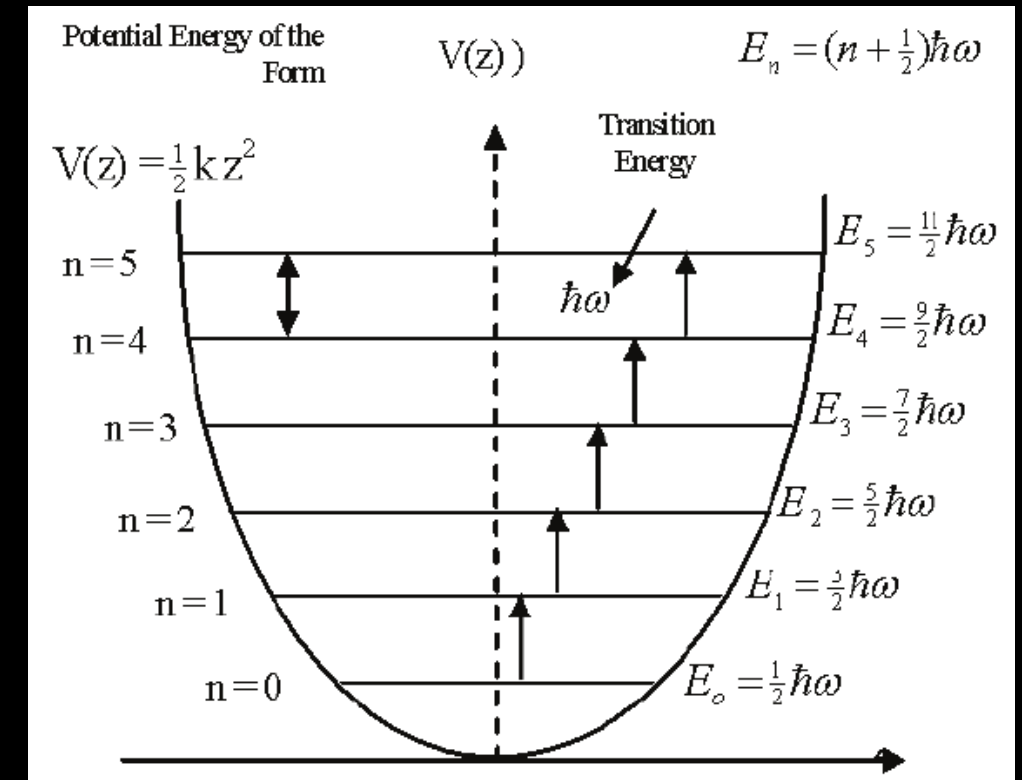
$$p = -\rho$$

Suppose space is filled with a quantum field; void is associated with an energy $E = h\nu/2$ for each oscillation mode

Quantum vacuum exerts a negative pressure, like in the Casimir effect

A constant vacuum energy is not measurable, but it contributes to the energy density

This gives an effective cosmological constant term, resulting in:



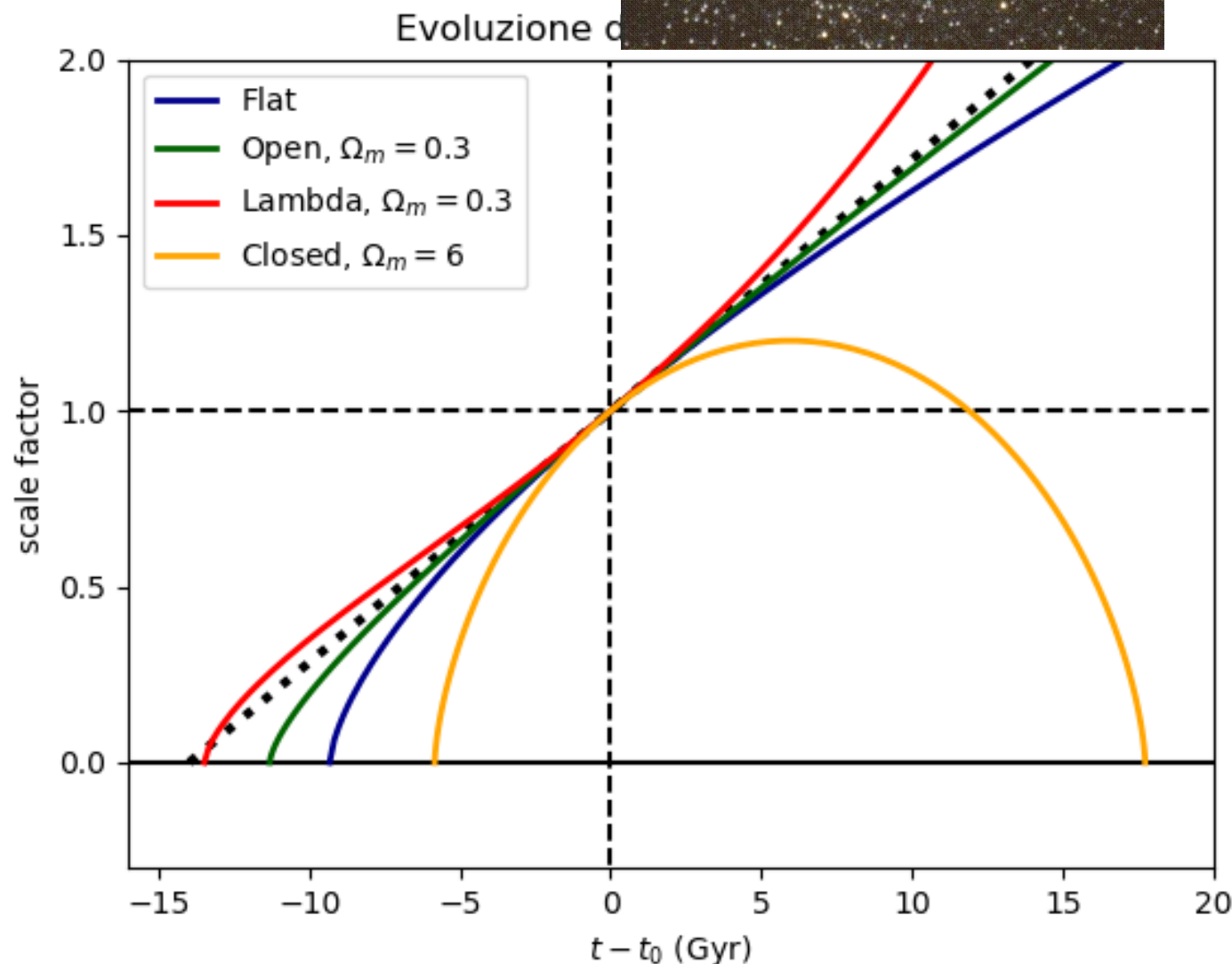
$$\Omega_\Lambda \sim 10^{120}$$

The cosmological constant: 3, FRW models ('90s)



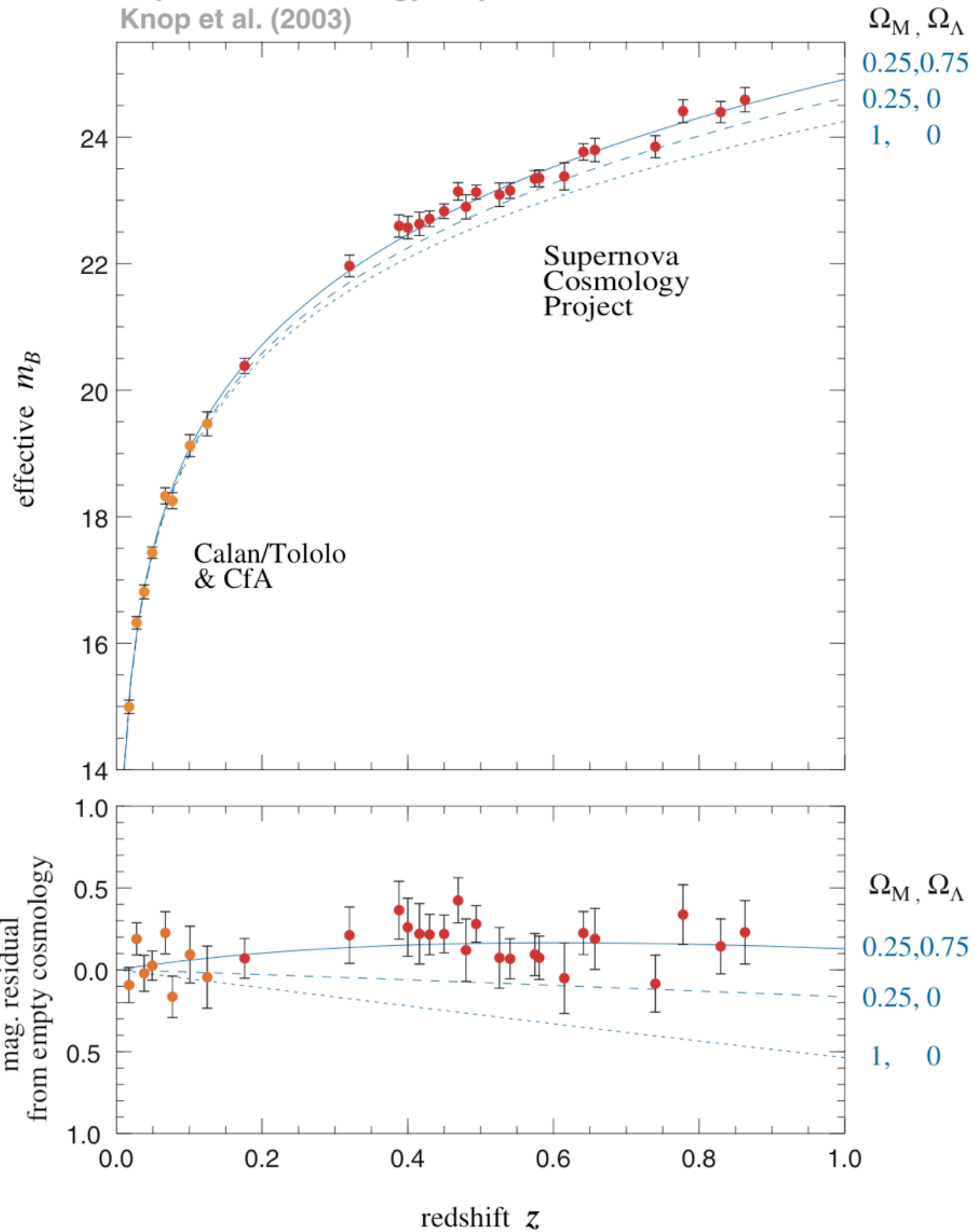
Solution of age problem
(and others):

- open universe
- cosmological constant



But cosmic inflation predicts a flat universe, so theoretical prejudice makes the open universe unattractive

Supernova Cosmology Project
Knop et al. (2003)



The LambdaCDM model

Most matter is made up by a dark matter particle, that does not interact with light and starts with low velocity dispersion ("cold")

Most energy is in the form of a cosmological constant that accelerates the expansion at late times

An early phase of cosmic inflation produced the observed fluctuations

The universe is spatially flat

An anthropic argument on multi-verses could explain the unnatural value of the cosmological constant