

Notes on Gunn-Peterson effect

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Photons bluer than the Lyman- α frequency can be absorbed by neutral hydrogen along the line of sight

We compute the opacity for photons emitted at frequency $\nu_e > \nu_{Ly\alpha}$ at redshift z_e .

These photons have frequency ν_0 at the observing time

$$d\tau_\nu = \alpha_\nu dl = N_{HI} \sigma_\nu dl \quad \text{from Ribickiy & Lightman}$$

The frequency ν here depends on redshift:

$$\nu_0(1+z) = \nu$$

The cross section is:

$$\sigma_\nu = \frac{\pi e^2}{m_e c} f g(\nu - \nu_{Ly\alpha}), \quad f = 0.616 \text{ is the oscillator strength}$$

$$g \text{ is the line profile, } \int_0^\infty g(x) dx = 1$$

Integrating along the line of sight from the emission redshift of the source to now:

$$\tau_{\nu_0} = \int_0^{z_e} N_{HI}(z) \sigma_\nu \frac{dl}{dz} dz$$

dl is the proper length interval for a redshift interval dz

From eq. 2.48 of Vittorio:

$$\frac{dl}{dz} = \frac{c}{H_0} \frac{1}{(1+z)} \frac{1}{\sqrt{\Omega_0(1+z)^3 + \Omega_\Lambda}} = \frac{c}{H_0 E(z)} \frac{1}{(1+z)} = \frac{c a}{H}$$

$$\text{or: } l = a \int \frac{cdt}{a} \Rightarrow dl = c dt \quad (\text{obvious!})$$

$$dt = \frac{dt}{da} \frac{da}{dz} dz = \frac{1}{\dot{a}} a^2 dz = \frac{a}{H} dz$$

$$dl = \frac{ac}{H} dz \Rightarrow \frac{dl}{dz} = \frac{c}{H} a = \frac{1}{(1+z)} \frac{c}{H_0} \frac{1}{E(z)}$$

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So:

$$\tau_{\nu_0} = \int_0^{z_e} n_{\text{HI}}(z) \frac{\pi e^2}{m_e c} f g(\nu_0(1+z) - \nu_{L\alpha}) \frac{c}{H_0} \frac{1}{E(z)} \frac{1}{(1+z)} dz$$

We want to integrate over the argument of g , and

$$d(\nu_0(1+z) - \nu_{L\alpha}) = \nu_0 dz$$

so we multiply and divide the integrand by ν_0 .

All z -dependent functions vary smoothly, so they can be taken out of the integral, obtaining:

$$\begin{aligned} \tau_{\nu_0} &\cong n_{\text{HI}}(z) \frac{\pi e^2}{m_e} f \frac{1}{H_0} \frac{1}{E(z)} \frac{1}{\nu_0(1+z)} \int_0^{z_e} g(\nu_0(1+z) - \nu_{L\alpha}) \nu_0 dz \\ &\cong n_{\text{HI}}(z) \frac{\pi e^2}{m_e} f \frac{1}{H_0} \frac{1}{E(z)} \frac{1}{\nu_{L\alpha}} \end{aligned}$$

where z is such that $\nu_0(1+z) = \nu_{L\alpha}$ ($= \frac{c}{\lambda_{L\alpha}} = 2.47 \times 10^{15} \text{ s}^{-1}$)

So:

$$\tau_{\nu} \cong n_{\text{HI}}(z) (4.7 \times 10^{10} \text{ h}^{-1} \text{ cm}^{-3}) \frac{1}{E(z)}$$

If $n_{\text{HI}} \cong n_b$ (no He and no ionized H)

$$n_{\text{HI}}(z) = n_{b0} (1+z)^3 \cong 2.5 \times 10^{-7} \text{ cm}^{-3} \left(\frac{\Omega_b h^2}{0.022} \right) (1+z)^3$$

$$\Rightarrow \tau_{\nu_0} \cong 10^4 \text{ h}^{-1} \left(\frac{\Omega_b h^2}{0.022} \right) \frac{(1+z)^3}{E(z)}$$

at, say, $z=2$: $\tau_{\nu_0}(z=2) \cong 1.3 \times 10^5$ ($h=0.67$)

This means that the flux bluewards of Lyman- α should NOT be observed, unless hydrogen is highly ionized